

This theory is 0

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Key words, Mandelbrot set, Dirac equation, Metric

Abstract All QM physicists know about *real* eigenvalue (Hermetian) observables (eg. eq.11 and its circles). All mathematicians know that the limit of a Cauchy sequence of rational numbers is a Cauchy real number (Cantor(7) 1872). So all we did here is show we postulated $real \neq 0$ by using it to derive a rational Cauchy sequence with limit 0. We did this because that same postulate (of $real \neq 0$) math *also* implies *the* real eigenvalues and so fundamental theoretical physics (eg., the Newpde solutions below). So the Ultimate Occam's Razor postulate implies the ultimate math-physics theory. Nothing is more (real)'Occam' than postulate 0.

So this **Theory** is 0. But we need to define the algebra first and use it to write the postulate 0. So define

1) numbers $1 \equiv 1+0$ in $0 \equiv 0 \times 0, 1 \equiv 1 \times 1$ as symbol $z = zz$: (algebraic definition of 0). So we can write
 2) Postulate *real* number 0 if $z' = 0$ and $z' = 1$ plugged into $z' = z'z' + C$ (eq.1) results in some $C = 0$ constant (ie $\delta C = 0$).

Single **Application**: Above theory says only to plug 1,0 into eq.1.

• Plug in $z=0=z_0=z'$ in **eq.1**. Note the equality sign in eq.1 demands we substitute z' on left (**eq.1**) into right $z'z'$ repeatedly and get iteration $z_{N+1} = z_N z_N + C$. So using that *other available* $C=1$ $z_1 = 0 \times 0 + 1 = (0+1) \equiv 1$ so $z_2 = 1 \times 1 + 1 \equiv (1+1) \equiv 2$ is now available etc., (both sums *defined* algebraically to be $(a+b) \equiv c$) and *define* rules of algebra (on these big numbers) like $a+b = b+a$ (eg., ring-field) with no new axioms. So postulate 0 also *generates* the big number algebra and calculus we can now use.

For example we can now define constant C with that $\delta C = 0$. When applied on iteration $z_{N+1} = z_N z_N + C, z_0 = 0$ it also requires we reject the C s for which $\delta C = \delta(z_{N+1} - z_N z_N) = \delta(\infty - \infty) \neq 0$. The C s that are left over define the fractal **Mandelbrot set** with (D3) lemniscate (so eq.11 circle observables) subsets $C_M = C = \delta z' = 10^{40N} \delta z, N = \text{integer}$ (since the derivation of eq.5 and thereby eq.11 requires we set the observer scale N anyway) also giving that required $C = 0 \in C_M$ ($C \approx 0$ from $z = 1 + \delta z'$ plugin is below). See fig1 zoom. Thus these fractal scales have their own $\delta z'$ (tiny circles) that must perturb that $z = 1$ Nth scale putting ansatz $z = 1 + \delta z$ into **eq.1** to get $\delta z + \delta z \delta z = C$ (3) Thus the (δz observables) numbers allowed are only in the Mandelbrot set.

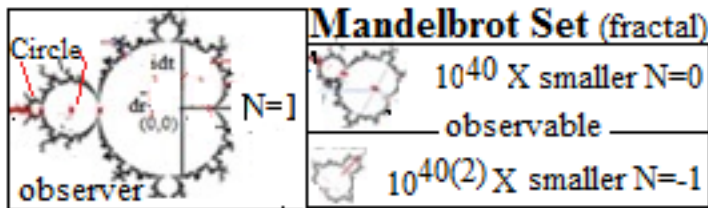


fig1 lemniscate subset. Defining the 'observers' (required by eq.11, so circle) scale as fractal scale N then $M < N$ (implied by eq.3) is the 'observables' scale M . For example we can define the fig1 'observer' fractal scales as $N=1$ implying $|\delta z| \gg 1$ since C is then huge in comparison to the $M=0$ scale. In addition to the above iteration, we must also solve eq 3 as a quadratic equation

$$\delta z = (-1 \pm \sqrt{1 + 4C}) / 2 \equiv dr + idt \text{ if } C < -1/4 \text{ (complex)} \quad (4)$$

(so also temporal observability). Note the Mandelbrot set iteration (ie., $z_{N+1}=z_N z_N - C$) for this $\delta C=0$ extremum $C=-1/4$ is a rational number Cauchy sequence $-1/4, -3/16, -55/256, \dots, 0$ thereby proving our above postulated *real#0* math. QED Note the $<$ in eq.4 implies the actual real number is the $N=0$ limit $C \approx \delta z = dr$ is not at exactly 0.

• **Plug in $z=1$** in $z'=1+\delta z$ in **eq1**, So $\delta C=0$ [eq1 implies eq3] $=\delta(\delta z + \delta z \delta z) = \delta \delta z(1) + \delta \delta z(\delta z) + (\delta z)\delta \delta z =$ (observer $|\delta z| \gg 1$ implying $M < N$) $\approx \delta(\delta z \delta z) = 0 =$ (plug in eq.4) $= \delta[(dr+idt)(dr+idt)] = \delta[(dr^2-dt^2) + i(dr dt + dt dr)] = 0$ (5)

$= 2D \delta[(\text{Minkowski metric, } c=1) + i(\text{Clifford algebra} \rightarrow \text{eq.7a})]$ (\equiv Dirac eq)

Factor **real** eq.5 $\delta(dr^2-dt^2) = \delta[(dr+dt)(dr-dt)] = 0 = [[\delta(dr+dt)](dr-dt)] + [(dr+dt)[\delta(dr-dt)]] = 0$ (6)

so $-dr+dt=ds, -dr-dt=ds \equiv ds_1 (\rightarrow \pm e)$ Squaring & eq.5 gives circle in e, v (dr, dt) 2nd, 3rd quadrants (7) & $dr+dt=ds, dr-dt=ds, dr \pm dt=0$, light cone ($\rightarrow v, \bar{v}$) in **same** (dr, dt) plane 1st, 4th quadrants (8)

& $dr+dt=0, dr-dt=0$ so $dr=dt=0$ defines vacuum (while eq.4 derives space-time) (9) Those quadrants give *positive* scalar $dr dt$ in eq.7 (if *not* vacuum) so imply the eq.5 *non* infinite extremum **imaginary** $\equiv dr dt + dt dr = 0 \equiv \gamma^i dr^i dt + \gamma^j dt^j dr = (\gamma^i \gamma^j + \gamma^j \gamma^i) dr dt$ so $(\gamma^i \gamma^j + \gamma^j \gamma^i) = 0, i \neq j$ (from **real** eq5 $\gamma^i \gamma^i = 1$) (7a) Thus from eqs 5, 7a: $ds^2 = dr^2 - dt^2 = (\gamma^i dr + i \gamma^j dt)^2$ Note how eq5 Dirac eq. and C_M Mandelbrot set just fall (pop) out of eq.1, amazing!

We square eqs.7 or 8 or 9 $ds_1^2 = (dr+dt)(dr+dt) = (-dr-dt)(-dr-dt) = [dr^2+dt^2] + (dr dt + dt dr) \equiv ds^2 + ds_3 = \text{Circle} + \text{invariant}$. **Circle** $\equiv \delta z = dse^{i\theta} = dse^{i(\Delta\theta + \theta_0)} = dse^{i((\cos\theta dr + \sin\theta dt)/(ds) + \theta_0)}$, $\theta_0 = 45^\circ$ min of $\delta ds^2 = 0$ given eq.7 constraint for $N=0$ $\delta z'$ perturbation of eq5 flat space. We define $k = dr/ds$, $\omega = dt/ds$, $\sin\theta \equiv r$, $\cos\theta \equiv t$. $dse^{i45^\circ} \equiv ds'$. Take ordinary derivative dr (since flat space) of 'Circle'.

$$\frac{\partial (dse^{i(\frac{rdr}{ds} + \frac{tdt}{ds})})}{\partial r} = i \frac{dr}{ds} \delta z \text{ so } \frac{\partial (dse^{i(rk+wt)})}{\partial r} = ik \delta z, k \delta z = -i \frac{\partial \delta z}{\partial r} \quad (11).$$

Recall from above that we proved that dr is a real number. So $k = dr/ds$ is an operator with *real* eigenvalues (So k is an observable). Also $k = 2\pi/\lambda$ (eg., in $\delta z = \cos kr$) thereby deriving the DeBroglie wavelength λ . Note the derivation of eq11 from that circle. Also from eq.11 we can write $\langle p_r \rangle^* = \int (p_r \psi)^* \psi d\tau = \int \psi^* p_r \psi d\tau = \langle p_r \rangle$. Therefore $p_r = \hbar k$ is Hermitian given dr is real.

Also need a $C \approx 0$ for $z=1$ plugin

For the $N=1$ huge observer $\delta z \gg \delta z \delta z$ from eq.3. Thus the required $N=-1, N=0$ tiny observable ($\delta z' \ll \delta z$) is a perturbation of the eq.7 $\delta z \approx dr \approx dt$ at 45° $(dr - \delta z') + (dt + \delta z') \equiv dr' + dt' = ds$ (12) But for the high energy big $\delta \delta z$ (extreme "axis" perturbations) δz is small. So finding big $\delta \delta z$ 'observables' requires we artificially stay on the circle (appendix C) implying this additional $\delta z'$ eq7 perturbation.

Define $\kappa_{rr} \equiv (dr/dr')^2 = (dr/(dr - \delta z'))^2 = 1/(1 - r_H/r)^2 = A_1/(1 - r_H/r) + A_2/(1 - r_H/r)^2$. The partial fractions A_I can be split off from RN and so

$$\kappa_{rr} \approx 1/[1 - r_H/r]. \quad (13)$$

$$\text{in } ds^2 = \kappa_{rr} dr'^2 + \kappa_{oo} dt'^2 \quad (14)$$

$$\kappa_{rr} = 1/\kappa_{oo} \quad (15)$$

From eq.7a $dr' dt' = \sqrt{\kappa_{rr}} dr' \sqrt{\kappa_{oo}} dt' = dr dt$ so

• **Both $z=0, z=1$** together (in eq1. Use 3D orthogonality to get (2D+2D curved space)). Thus $\delta z' + \delta z = (dx_1 + idx_2) + (dx_3 + idx_4) \equiv dr + idt$ given $dr^2 - dt^2 = (\gamma^i dr + i \gamma^j dt)^2$ if $dr^2 \equiv dx^2 + dy^2 + dz^2$ (3D orthogonality) so that $\gamma^i dr \equiv \gamma^x dx + \gamma^y dy + \gamma^z dz, \gamma^i \gamma^j + \gamma^j \gamma^i = 0, i \neq j, (\gamma^i)^2 = 1$. From eq.14

$(\gamma^x \sqrt{\kappa_{xx}} dx + \gamma^y \sqrt{\kappa_{yy}} dy + \gamma^z \sqrt{\kappa_{zz}} dz + \gamma^t \sqrt{\kappa_{tt}} idt)^2 = \kappa_{xx} dx^2 + \kappa_{yy} dy^2 + \kappa_{zz} dz^2 - \kappa_{tt} dt^2 = ds^2$. Multiply both sides by \hbar^2/ds^2 and $\delta z^2 \equiv \psi^2$ use eq11 inside brackets () get 4D QM $\gamma^\mu (\sqrt{\kappa_{\mu\mu}}) \partial \psi / \partial x_\mu = (\omega/c) \psi \equiv \text{Newpde}$ for $e, v, \kappa_{oo} = 1 - r_H/r = 1/\kappa_{rr}, r_H = C_M/\xi = e^2 X 10^{40N}/m$ ($N = -1, 0, 1, \dots$). **Postulate(0) \rightarrow Newpde**

Solution to Newpde e,v:

stable $2P_{3/2}$ at $r=r_H$

$N=0$ Mandelbulbs: **Free space:** τ, μ, e leptons. OnSphere $2P_{3/2}$ at $r=r_H$ $3e$ baryons (QCD not required)

$N=1$ inside zitterbewegung oscillation $r < r_C$ $N=1$ puts us in the **cosmological expansion** stage

Other results (besides Newpde) from postulate0, eg., Copenhagen stuff and 10^{82} electrons e between consecutive fractal scales such as $N=1$ cosmological $r=r_H, 2P_{3/2}$

perturbation objects B,C

Object B $N=0$ perturbations of κ_{oo} and κ_{rr} in the Newpde **E&M**, $N=-1$ gravity GR

Object C $N=0$ “ “ “ **weak**, both SM

That eq1 **iteration generates** (rel0)& possibly larger numbers $1+1 \equiv 2, 2+1 \equiv 3$.etc as *defined* symbols $a+b=c$ and **algebra rules** (eg.,ring-field def. like $a+b \equiv b+a$). *with no new axioms*.

Thus (with the math&physics) we understand *everything* (eg GR, cosmology, QM, e,v SM, baryons, rel#).

So the *simplest idea imaginable 0* implies all *fundamental math-physics*. no more, no less(eg simply 4D)

•**Conclusion:** So by merely (plugging 0,1 into eq.1) **postulating 0**, out pops the whole universe, BOOM! easily the most important discovery ever made or that will ever be made again. We finally figured it out.

Reminder:

Note also that 0 is not the same as the null set \emptyset since 0 is a real number and \emptyset isn't.

So when we postulate 0 we are also implicitly postulating the real# nature of 0.

Summary This Theory is 0 The rest is a real number 0 **definition**

We need to define the algebra first and use it to write postulate0. So define numbers $1 \equiv 1+0$, and $0 \equiv 0 \times 0, 1 \equiv 1 \times 1$ as symbol $z=zz$ the simplest algebraic definition of 0. **So**

Theory	<u>Real# 0 definition</u>
Postulate 0	if <u>$z=1$</u> and <u>$z=0$</u> (plugged) into <u>$z=zz+C$ eq.1</u> gives some $C=0$ constant (ie $\delta C=0$)

Can't be more Occam than postulate(0). All the rest is those 2 'plug in' **applications**. So must

1) **plug $z=0$** into eq.1 (given the implied $z=zz+C$ eq1 iteration) to get the 2D (fractal) Mandelbrot set to prove postulate real 0. Iteration numbers might be bigger $1+1 \equiv 2, 1+2 \equiv 3$ etc., with *defined* symbols $a+b \equiv c$ and algebra rules eg., $a+b \equiv b+a$. So the *postulate generates real# math* without extra axioms

2) **plug $z=1$** into eq.1 and get 2D Dirac equation.

Both Mandelbrot and Dirac results together give $2+2=4D$ Newpde (plus some Copenhagen stuff)

Solve the differential equation (Newpde) to get the physical universe, no more, no less.
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Conclusion: Ultimate Occam's razor postulate(0) implies ultimate math-physics.

This Theory is 0

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Section 1

Postulate *real number 0 if $z'=0$ and $z'=1$ plugged into $z'=z'z'+C$ (eq.1) results in some $C=0$ constant (ie $\delta C=0$)*

These are 2 plugin **Applications**: Plug 0 into eq.1 iteration get 2D Mandelbrot set (fractal scale N)
Plug 1 into eq.1 get 2D Dirac eq. $N=1$

Dirac with Mandelbrot together gets e, ν 4D Newpde:

$\gamma^\mu (\sqrt{\kappa_{\mu\nu}}) \partial \psi / \partial x_\mu = (\omega/c) \psi \equiv$ **Newpde** for $e, \nu, \kappa_{00}=1-r_H/r=1/\kappa_{rr}, r_H=C_M/\xi=e^2 X 10^{40N}/m$ ($N=-1, 0, 1, \dots$).

Solve Newpde e, ν ; stable state **2P_{3/2}** at $r=r_H$

$N=0$ Mandelbulbs: **Free space**: τ, μ, e leptons. On **Sphere 2P_{3/2}** at $r=r_H$ **3e baryons** (QCD not required)

$N=1$ **inside e** zitterbewegung oscillation $r < r_C$ $N=1$ puts us in the **cosmological expansion** stage.

Appendix A Other QM results (besides Newpde) from postulate0, eg., **Copenhagen** stuff. and 10^{82} electrons **e** between fractal scales such as cosmological $N=1$ **e** objects B,C inside $r=r_H, 2P_{3/2}$
Newpde perturbation of κ_{00}, κ_{rr} with these **e** objects B,C

Appendix B $N=0$ perturbations of the Newpde κ_{00} and κ_{rr} by object B **E&M**, $N=-1$ gravity GR

Appendix C $N=0$ “ “ “ object C **weak**, both **SM**

Appendix D Mathematical symbols: That eq1 **iteration generates** (rel0)& possibly larger numbers $1+1 \equiv 2, 2+1 \equiv 3$. etc as *defined* symbols $a+b=c$ and **algebra rules** (eg., ring-field def. like $a+b \equiv b+a$). *with no new axioms*.

So Ultimate Ocam's razor postulate(0) implies ultimate math-physics!!!!!!!!!!!!!!!!!!!!!!!!!!!!!!!!!!!!

sect.1 So this **Theory** is **0**

Abstract All QM physicists know about *real* eigenvalue (Hermetian) observables (eg. eq.11 and its circles). All mathematicians know that the limit of a Cauchy sequence of rational numbers is a Cauchy real number (Cantor(7) 1872). So all we did here is show we postulated *real#0* by using it to derive a rational Cauchy sequence with limit 0. We did this because that same postulate (of *real#0*) math *also* implies *the* real eigenvalues and so fundamental theoretical physics (eg., the Newpde solutions below). So the Ultimate Occam's Razor postulate implies the ultimate math-physics theory. Nothing is more (real) 'Occam' than postulate0.

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1) *numbers* $1 \equiv 1+0$ in $0 \equiv 0X0, 1 \equiv 1X1$ as *symbol* $z=zz$: (algebraic definition of **0**). So we can write
2) **Postulate** *real* number **0** if $z'=0$ and $z'=1$ plugged into $z'=z'z'+C$ (**eq.1**) results in some $C=0$ constant (ie $\delta C=0$).

So Postulate0 is our *entire* (ultimate Occam's razor postulate) **theory**.

Note: Also that "some C" implies there *might be* other C with the only other *available* number $C=1$ in the above theory beside 0 (available \equiv previously derived from postulate(0) as in a symbolic logic proof.)

So 'postulate0' as *real* (in the single 'observables' application: eq.11 for *real* eigenvalues).

Single **Application:** Above theory says only to plug 1,0 into eq.1.

Jump to “• Conclusion” below for result.

•**Plug in $z=0=z_0=z'$ in eq1.** [Must use iteration *and* quadratic eq.] Note the equality sign in eq.1 demands we substitute z' on left (**eq1**) into right $z'z'$ repeatedly and get iteration $z_{N+1}=z_N z_N + C$. So using that *other available* $C=1$ $z_1=0 \times 0 + 1 = (0+1) \equiv 1$ so $z_2=1 \times 1 + 1 \equiv (1+1) \equiv 2$ is now available etc. (both sums *defined* algebraically to be $(a+b) \equiv c$) and *define* rules of algebra (on these big numbers) like $a+b=b+a$ (eg., ring-field) with no new axioms. So postulate 0 also *generates* the big number algebra and calculus we can now use.

For example we can now define constant C with that $\delta C=0$. When applied on iteration $z_{N+1}=z_N z_N + C$, $z_0=0$ it also requires we reject the C s for which $\delta C = \delta(z_{N+1} - z_N z_N) = \delta(\infty - \infty) \neq 0$. The C s that are left over define the fractal **Mandelbrot set** with (D3) lemniscate (so eq.11 circle observables) subsets $C_M = C = \delta z' = 10^{40N} \delta z$, $N = \text{integer}$ (since the derivation of eq5 and thereby eq.11 requires we set the observer scale N anyway) also giving that required $C=0 \in C_M$ ($C \approx 0$ from $z=1+\delta z'$ plugin is below.). See fig1 zoom. Thus these fractal scales have their own $\delta z'$ (tiny circles) that must perturb that $z=1$ Nth scale putting ansatz $z=1+\delta z$ into **eq.1** to get $\delta z + \delta z \delta z = C(3)$ Thus the (δz observables) numbers allowed are only in the Mandelbrot set.

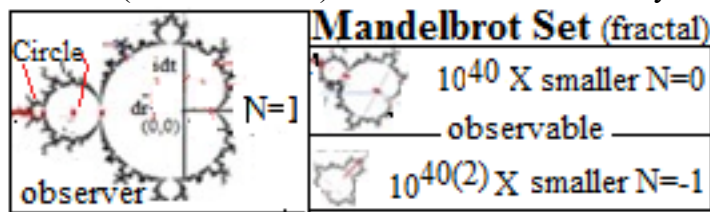


fig1 lemniscate subset. Defining

‘observers’ (required by eq.11, so circle) scale as fractal scale N then $M < N$ (implied by eq.3) is the ‘observables’ scale M . For example we can define the fig1 ‘observer’ fractal scales as $N=1$ implying $|\delta z| \gg 1$ since C is then huge in comparison to the $M=0$ scale. In addition to the above iteration, we must also solve eq 3 as a quadratic equation

$$\delta z = (-1 \pm \sqrt{1 + 4C}) / 2 \equiv dr + idt \text{ if } C < -1/4 \text{ (complex)} \quad (4)$$

(so also temporal observability). Note the Mandelbrot set iteration (ie., $z_{N+1} = z_N z_N - C$) for this $\delta C = 0$ *extremum* $C = -1/4$ is a rational number Cauchy sequence $-1/4, -3/16, -55/256, \dots, 0$ thereby proving our above postulated *real#0* math. QED Note the $<$ in eq.4 so the actual real number is the $N=0$ limit $C \approx \delta z = dr$ not at exactly 0.

•**Plug in $z=1$ in $z'=1+\delta z$ in eq1,** So $\delta C=0 = [\text{eq1 implies eq3}] = \delta(\delta z + \delta z \delta z) = \delta \delta z(1) + \delta \delta z(\delta z) + (\delta z) \delta \delta z = (\text{observer } |\delta z| \gg 1 \text{ implying } M < N) \approx \delta(\delta z \delta z) = 0 = (\text{plug in eq.4}) = \delta[(dr+idt)(dr+idt)] =$

$$\delta[(dr^2 - dt^2) + i(dr dt + dt dr)] = 0 \quad (5)$$

$$= 2D \delta[(\text{Minkowski metric, } c=1) + i(\text{Clifford algebra} \rightarrow \text{eq.7a})] \quad (\equiv \text{Dirac eq})$$

Factor **real** eq.5 $\delta(dr^2 - dt^2) = \delta[(dr+dt)(dr-dt)] = 0 = [[\delta(dr+dt)](dr-dt)] + [(dr+dt)[\delta(dr-dt)]] = 0 \quad (6)$

Here we make $ds > 0$ by putting these $dr \pm dt$ in appropriate quadrants so that ds is the radius of a circle (eg., In the 4th quadrant $dr-dt=ds$ is positive) so our rotations are for observables, as in eq11 so $-dr+dt=ds, -dr-dt=ds \equiv ds_1 (\rightarrow \pm e)$ Squaring & eq.5 gives circle in e, v (dr, dt) 2nd, 3rd quadrants (7)

& $dr+dt=ds, dr-dt=0$, light cone ($\rightarrow v$) in **same** (dr, dt) plane 1st quadrant (8)

& $dr-dt=ds, dr+dt=0$, light cone ($\rightarrow \bar{v}$) in **same** (dr, dt) plane 4th quadrant. (9)

& $dr+dt=0, dr-dt=0$ so $dr=dt=0$ defines vacuum (while eq.4 derives space-time)

Quadrants give *positive* scalar $drdt$ of eq.7 (if *not* vacuum) imply the eq.5 *non* infinite extremum **imaginary** $\equiv drdt + dt dr = 0 \equiv \gamma^i dr \gamma^j dt + \gamma^j dt \gamma^i dr = (\gamma^i \gamma^j + \gamma^j \gamma^i) dr dt$ so $(\gamma^i \gamma^j + \gamma^j \gamma^i) = 0$, $i \neq j$ (from **releq5** $\gamma^i \gamma^i = 1$) (7a)
 Thus from eqs 5, 7a: $ds^2 = dr^2 - dt^2 = (\gamma^r dr + i \gamma^t dt)^2$ (10)

We square eqs.7 or 8 or 9 $ds_1^2 = (dr+dt)(dr+dt) = (-dr-dt)(-dr-dt) = [dr^2+dt^2] + (drdt+dt dr)$
 \equiv **Circle** + invariant. **Circle** $\equiv \delta z = dse^{i\theta} = dse^{i(\Delta\theta+\theta_0)} = dse^{i((\cos\theta dr + \sin\theta dt)/(ds)+\theta_0)}$, $\theta_0 = 45^\circ$ min of $\delta ds^2 = 0$
 given eq.7 constraint for $N=0$ $\delta z'$ perturbation of eq5 flat space. We define $k \equiv dr/ds$, $\omega \equiv dt/ds$,
 $\sin\theta \equiv r$, $\cos\theta \equiv t$. $dse^{i45^\circ} \equiv ds'$. Take ordinary derivative dr (since flat space) of 'Circle'

$$\frac{\partial \left(dse^{i \left(\frac{r dr}{ds} + \frac{t dt}{ds} \right)} \right)}{\partial r} = i \frac{dr}{ds} \delta z \text{ so } \frac{\partial (dse^{i(rk+wt)})}{\partial r} = ik \delta z, \quad k \delta z = -i \frac{\partial \delta z}{\partial r} \quad (11).$$

Recall from above that we proved that dr is a real number. So $k = dr/ds$ is an operator with *real* eigenvalues (So k is an observable). Also $k = 2\pi/\lambda$ (eg., in $\delta z = \cos kr$) thereby deriving the DeBroglie wavelength λ . Note the derivation of eq11 from that circle. Also from eq.11 we can write $\langle p_r \rangle^* = \int (p_r \psi)^* \psi d\tau = \int \psi^* p_r \psi d\tau = \langle p_r \rangle$. Therefore $p_r = \hbar k$ is Hermitian given dr is real. Recall dr are real numbers and so our eigenvalues dr/ds are real observables. Thus the Mandelbrot set iteration here did double duty also as proof of the real number eigenvalues(observables) in eq.11.

Cancel that e^{i45° coefficient ($45^\circ = \pi/4$) then multiply both sides of eq.11 by \hbar and define $\delta z \equiv \psi$, $p \equiv \hbar k$. Eq.11: the familiar '**observables**' p_r in

$$p_r \psi = -i \hbar \frac{\partial \psi}{\partial r} \quad (11)$$

Repeat eq.3 for the τ , μ respective δz Mandelbrot set lobes in fig.6 so they each have their own neutrino ν : Lepton generations.

That means the **mathematics and the physics** come from (**postulate 0**): *everything*. Recall from eq.7 that $dr+dt=ds$. So combining in quadrature eqs 7&11 $SNR \times \delta z = (dr/ds+dt/ds)\delta z = ((dr+dt)/ds)\delta z = (1)\delta z$ (11a) and so having come *full circle* back to sect.1 postulate 0 as a real#

Also need a $C \approx 0$ for $z=1$ plugin

For the $N=1$ huge observer $\delta z \gg \delta z \delta z$ from eq.3. Thus the required $N=-1, N=0$ tiny observable ($\delta z' \ll \delta z$) is a perturbation of the eq.7 $\delta z \approx dr \approx dt$ at 45° $(dr - \delta z') + (dt + \delta z') \equiv dr' + dt' = ds$ (12)
 But for the high energy big $\delta \delta z$ (extreme "axis" perturbations) δz is small. So finding big $\delta \delta z$ 'observables' requires we artificially stay on the circle (appendix C) allowing this additional $\delta z'$ eq7 perturbation.

so big ξ . $C_M = e^2 10^{40N}$ defines charge, $\xi = \gamma$ defines mass.

That figure 1 Mandelbrot set structure can be pulled out of the zoom clutter because of the above 4X circle observability sequence in fig1

We can pull out the above 4X circle observability sequence in fig1 from the zoom clutter

$$\text{Recall } C \text{ is a function on the complex } (dr, idt) \text{ plane so } \delta C = \left(\frac{\partial C}{\partial r} \right)_t dr + \left(\frac{\partial C}{\partial t} \right)_r idt = 0 \quad (12)$$

implying there are several $\delta C=0$ (dr, idt) extreme possible here. The first 1D extremum is provided by eq.4 and is that dr axis extremum $C_M = -1/4$ which incidently is the only rational number extremum on our C_M . Another extremum clearly is that $\partial C / \partial t = 0$, $dr = \text{constant}$, The last 1D extremum is $\partial C / \partial r = 0$, $dt = \text{constant}$ $N=2$ (observable internal QMS jumps in fig1 in partIII) with the rest unobservable.

To get the only 2D dr, idt extremum we divide eq.12 by dt so that above fig.1 4X sequence of those *observable* circles $drdt = d\text{area}_M \neq 0$ (so eq.11 observables) the real $\delta C=0$ extremum given the decreasing observable *real* circle radius sequence $\lim_{m \rightarrow \infty} \frac{\partial C}{\partial (drdt)_m} dr_m = \lim_{m \rightarrow \infty} \frac{\partial C}{\partial \text{area}_m} dr_m =$

$\lim_{m \rightarrow \infty} \frac{\partial C}{\partial \text{Circle}_m} dr_m = KX0 = 0$ (since $dr_\infty \approx 0$) = Fiegenbaum point = $f^\alpha = (-1.40115, i0) = C_M = \text{end}$ (12a)

our final realization of $\delta C = 0$. So random circles in the zoom don't do $\delta C = 0$. Note if a circle (or many circles) is rotated (U), translated (D), shrunk (S) equally in both dimensions (i.e.,

$(\partial X^i / \partial X'^k) f^i = f'^k \equiv \begin{bmatrix} f_{1N} \\ f_{2N} \end{bmatrix} = S_N \begin{bmatrix} U_{11} & U_{12} \\ U_{21} & U_{22} \end{bmatrix} \begin{bmatrix} f \\ 0 \end{bmatrix} + \begin{bmatrix} D_{1N} \\ D_{2N} \end{bmatrix}$) it is still a circle, eq.11 still holds, so *it's still an observable* as seen in the N fractal scale zoom. Thus you can pick out from that zoom these fig.1 Mandelbrot set extremum 4X diameter circles as the only **observables** and $\delta C = 0$ extremum geometry in all that clutter. Reset the zoom, restart at such $S_N C_M = 10^{40N} C_M$ in eq.17.

C=C_M Source Small C definition and two Mandelbulb masses

Our postulate 0: (z=zz+C)≈(z=zz+0)) requires a small C

But for N=0 eq.3 $\delta z + \delta z \delta z = C$ reads $C \approx \delta z$. So that postulated small $C \approx 0$ implies an eq.5 Lorentz (Fitzgerald) contraction (9) $1/\gamma$ boosted frame of reference (fig.6) **small $C \approx \delta z / \gamma = C_M / \xi = \delta z'$** (10)

Therefore $\delta C = 0$ and eq.10 implies we take variation of $C = C_M = \xi \delta z$

So this same ξ is merely large in eq.17 with this N=0 $\delta z'$ the curved space perturbation $\delta z'$ in eqs.11,12. Also in *sect.1* $z' = 1 + \delta z$ z is called the perturbation z' . So on N=0 **$\delta C = 0 = \delta(\delta z) = \delta(z' - 1) = \delta z = 0$** so even perturbation z is the extreme of $\delta z' = -1$ or $\delta z = 0$ corresponding to fundamental $z = 0, 1$.

So take variation $\delta C = \delta C_M = (\delta \xi) \delta z + \xi \delta \delta z = 0$. Also recall above plugin $z = 1 + \delta z$. So if

$$\delta z \text{ is small so } \delta \xi \text{ and } \xi \text{ can be large (unstable large mass } \tau + \mu, \text{ sect.D4).} \quad (14)$$

And extremum perturbation $z = 1$ is the reduced mass $\tau + \mu = 2m_p$.

For large $|\delta z|$ in the above variation then

$$\delta \xi \text{ and } \xi \text{ can be small (stable small mass: electron ground state } \delta z \text{)} \quad (15)$$

with perturbation $\delta z = -1$.

From here on look only at what we are **allowed to observe**: eq.11 circles: so **$\delta(ds^2) = 0$** , proper frame. **Nothing else matters but these observables.** (Which are also $N < 1$ for $N = 1$ observer except for observer $N = 2$ seeing what we see: 'observables' can thereby be $N = 1$ cosmology objects (eq.B3a))

For N=1 At high energy Lorentz boost $1/\gamma$ of $\lambda = \delta z = dr$ then gets small relative to 1 and so $\delta \delta z$ gets bigger since we start approaching $N = 0$ instead (of $N = 1$) and so eq.5 fails except for **observables** if for them we still keep (circle) **$dr^2 - dt^2 = ds^2 = \text{radius}^2$** constant by expressing 'large $\delta \delta z'$ as a rotation at 45° in a slightly modified eq.7: $(dr - \delta z') + (dt + \delta z') \equiv dr' + dt' = ds$ (16)

For N=0 $\theta_0 = 45^\circ$ min of $\delta ds^2 = 0$ given eq.7 constraint $\delta z'$ perturbation of eq5 flat space and so $\delta z'$ in eq.16 is large relative to dr, dt . So given the max extremum for ds^2 is on the axis' each extreme can now be $\Delta \theta = \pm 45^\circ$. So in eq.16 the 4 rotations $45^\circ + 45^\circ = 90^\circ$ define 4 Bosons (see Ch.6). But

For N=-1 $45^\circ - 45^\circ$ $N < 0$ then contributes (appendix A2) so you also have other (smaller and **infinitesimal** $N = -1$) fractal scale extreme $\delta z'$ (eg., tiny Fiegenbaum pts so $N = 1$ $dr = r$, for $N_{ob} = -1$) so metric coefficient $\kappa_{rr} \equiv (dr/dr')^2 = (dr/(dr - (C_M/\xi_1)))^2 = 1/(1 - r_H/r)^2 = A_1/(1 - r_H/r) + A_2/(1 - r_H/r)^2$.

The partial fractions A_i can be split off from RN and so $\kappa_{rr} \approx 1/[1 - ((C_M/\xi_1)r)]$ (17)

$$(C_M \text{ defined to be } e^2 \text{ charge, } \gamma \equiv \xi_1 \text{ mass). So: } ds^2 = \kappa_{rr} dr'^2 + \kappa_{oo} dt'^2 \quad (18)$$

$$\text{From eq.7a } dr' dt' = \sqrt{\kappa_{rr}} dr' \sqrt{\kappa_{oo}} dt' = dr dt \text{ so } \kappa_{rr} = 1/\kappa_{oo} \quad (19)$$

We can then do a rotational dyadic coordinate transformation of $\kappa_{\mu\nu}$ to get the Kerr metric which is all we need for our applications(9). Recall also from eqs5,7a that $dr^2-dt^2=(\gamma^r dr+i\gamma^t dt)^2$

1.5 Both $z=0, z=1$ together using orthogonality get (2D+2D curved space) . So $(z=1)+(z=0)=(dx_1+idx_2)+(dx_3+idx_4) \equiv dr+idt$ given $dr^2-dt^2=(\gamma^r dr+i\gamma^t dt)^2$ if $dr^2 \equiv dx^2+dy^2+dz^2$ (3D orthogonality) so that $n\gamma^r dr \equiv \gamma^x dx + \gamma^y dy + \gamma^z dz$, $\gamma^i \gamma^j + \gamma^j \gamma^i = 0$, $i \neq j$, $(\gamma^i)^2 = 1$, rewritten (with invariant (8) $\kappa_{\mu\nu}$ eq.17-19) $(\gamma^x \sqrt{\kappa_{xx}} dx + \gamma^y \sqrt{\kappa_{yy}} dy + \gamma^z \sqrt{\kappa_{zz}} dz + \gamma^t \sqrt{\kappa_{tt}} idt)^2 = \kappa_{xx} dx^2 + \kappa_{yy} dy^2 + \kappa_{zz} dz^2 - \kappa_{tt} dt^2 = ds^2$. Multiply both sides by $1/ds^2$ and $\delta z^2 \equiv \psi^2$ use use operator equation 11 inside brackets () get curved space 4D

$$\gamma^\mu (\sqrt{\kappa_{\mu\mu}}) \partial \psi / \partial x_\mu = (\omega/c) \psi \quad (20)$$

\equiv Newpde for $e, \nu, \kappa_{00} = 1 - r_H/r = 1/\kappa_{rr}$, $r_H = e^2 X 10^{40N}/m$ ($N = -1, 0, 1, \dots$). Also $C_M/\xi = r_H =$

*small C so big $\xi = \gamma$ boost so $z = \xi z$ so **postulate 0**. So we really did just postulate 0. So

Postulate 0 \rightarrow Newpde

Solutions to New pde (given $e^2 10^{40N}$ X fractal scales N)

$N = -1$ is GR. $e^2 X 10^{40(-1)} = e^2/10^{40} \equiv Gm_e^2$, solve for G. So given eq,17-19 $\kappa_{\mu\nu} \equiv g_{\mu\nu}$ Schwarzschild metric

$N = 0$ At $r \gg r_H$ becomes usual Dirac equation e solutions with *nonrelativistic* limit the Schrodinger equation and appendix C Standard electroweak Model (SM) perturbations.

$N = 0$ e perturbations. At $r \approx r_H$ $1S_{1/2} \mu$ $2S_{1/2} \tau$, $2P_{3/2}$ $2P_{1/2}$ $3e$ is baryon core(QCD not required). (Part2)

1.6 Newpde $2P_{3/2}$ at $r=r_H$ state Contrast with QCD

In part II we find the only stable 3 body solution(state) to the Newpde is the $2P_{3/2}$ at $r=r_H$.

The electron (solution to that new pde) spends 1/3 of its time in each $2P_{3/2}$ (at $r=r_H$) lobe, *explaining the lobe multiples of 1/3e fractional charge* (The ‘lobes’ can be named ‘quarks’ or George if you want). The lobes are locked into the center of mass, can’t leave, *giving asymptotic freedom* (otherwise yet another ad hoc postulate of qcd). As derived in PartII the two positrons are ultrarelativistic ($\gamma=917$, sect.7.5, $3e=(\gamma m_e + \gamma m_e) = m_{p\delta\delta}$) so the field line separation is narrowed into plates explaining the strong force (otherwise postulated by qcd). Also there are 6 $2P$ states *explaining the 6 quark flavors*. P wave scattering *gives the jets*. We have *stability* ($dt'^2 = (1 - r_H/r) dt^2$) since the dt' clocks stop at $r=r_H$. That 2 γ ray scattering off the 3rd mass (in $2P_{3/2}$) diagonal metric(eq.17) time reversal invariance also reverses the γ ray pair annihilation with the subsequent e^\pm pair creation inside the r_H volume given $\sigma = \pi r_H^2 \approx (1/20)$ barn making it merely a virtual creation-annihilation event. So our $2P_{3/2}$ composite $3e$ (proton) at $r=r_H$ is the *only* stable multi e composite So quarks don’t exist, it’s all just 2 Newpde positrons and electron in $2P_{3/2}$ at $r=r_H$ states.

Newpde $1S_{1/2} 2S_{1/2}$ at $r \leq r_H$ States: Recall that $C = \delta z/\gamma = C_M/\gamma \equiv C_M/\xi$. $\xi = e + \mu + \tau = 2P$. Given only stable $2P_{3/2}$ at $r=r_H$: then there are only (Hund’s rule) $2m_p = G + 1S_{1/2} + 2S_{1/2} = e + \mu + \tau$. Here we use this relation and the Schrodinger equation for the observer comoving with the P COM to derive the ratios between muon to taoun to electron masses. Recall from sect.1:

$$(\gamma^x \sqrt{\kappa_{xx}} dx + \gamma^y \sqrt{\kappa_{yy}} dy + \gamma^z \sqrt{\kappa_{zz}} dz + \gamma^t \sqrt{\kappa_{tt}} idt)^2 = \kappa_{xx} dx^2 + \kappa_{yy} dy^2 + \kappa_{zz} dz^2 - \kappa_{tt} dt^2 = ds^2. \quad (23)$$

Recall the nonrelativistic limit of the Dirac equation is the Schrodinger equation where our energies are close to rest mass energy.

In partII that $3e 2P_{3/2}$ at $r=r_H$ was the only multibody stable state(i.e.,proton) with that $2P = m_\tau + m_\mu + m_e$ free space from $G + 1S_{1/2}, 2S_{1/2} = 3k$. Hund rule where this energy is the same as

that reduced mass two electron motion (those two positrons in orbit around the central electron) energy. It is an analog state of the group 2 (alkaline earth) electronic configuration in the periodic table of elements. G is the electron, the 'ground state' for them all, just as in chemistry. Here though we differ from chemistry in that we are at $r=r_H$ much smaller than the Bohr radius.

Schrodinger equation for Newpde for this state

Recall the Schrodinger equation is the nonrelativistic limit of the Newpde (modified Dirac) equation.

1) Recall associated 2 body *energy eigenvalues of Newpde Schrodinger equation* hydrogen atom $r > r_H$ Rydberg formula

$$E = R_b / N^2$$

2) The resulting *energy eigenvalues of the Newpde Schrodinger eq.* at $r=r_H$ is the Kioda formula:

$$\frac{m_\tau + m_\mu + m_e}{(\sqrt{m_\tau} + \sqrt{m_\mu} + \sqrt{m_e})^2} = \frac{2}{3}$$

Kioda eq. derivation from Newpde Schodinger equation at $r=r_H$.

Nonrelativistic reduced COM $r > r_H$ observer model For 2P=D Deuterium

Also recall Schrodinger equation (nonrelativistic): $H\psi = -\frac{\hbar}{2m} \frac{\partial^2}{\partial r^2} \psi$, $P\psi = -\frac{\hbar}{2P} \frac{\partial^2}{\partial r^2} \psi = -\frac{\hbar}{D} \frac{\partial^2}{\partial r^2} \psi$ or with eq.11 $\hbar(dr/ds)\psi = -i\hbar d\psi/dr$ with \hbar canceling out:

$$k\psi = -\frac{1}{2m} \frac{\partial^2}{\partial r^2} \psi \quad (k=dr/ds) = -\frac{1}{2m} \left(\frac{dr}{ds}\right)^2 \psi = \left(\sqrt{\frac{1}{2m} \frac{dr}{ds}}\right)^2 \psi$$

This D is 2Xproton mass singlet here (Not the actual ortho.) so regard this as a Boson allowing us to exactly drop the Pauli term.

Associated with the $2P_{3/2}$ state is the usual Hund's rule G, $1S_{1/2}$, $2S_{1/2}$ $m_\tau + m_\mu + m_e = 2P$ free space particles wrapped around the $2P_{3/2}$ spherical shell at $r=r_H$ interior mass giving the two ultrarelativistic positron energies of each $2P_{3/2}$ which is the only stable 3e composite state. Thus the reduced mass P is composed of these 2 relativistic particles which for the outside observer (outside of r_H) have a *nonrelativistic* COM mass D in the comoving system allowing us to still use the Schrodinger equation. Recall also (sect.1.5) that the linear dx_i s ($= dr' = \gamma^r \sqrt{\kappa_{rr}} dr \equiv dr'$, $\gamma^r \gamma^r = 1$) observables perturbations add in the complex plane so the Dirac equation for lepton multiplets G, $1S_{1/2}$, $2S_{1/2}$ can be summed under the square(brackets) in eq.23

$$3k\psi = \left(\sum_1^3 \sqrt{\frac{1}{2m} \frac{dr}{ds}}\right)^2 \psi$$

So all the relativistic effects are thrown into the $P=m$ mass black box allowing us to still use the exact nonrelativistic Schrodinger equation outside r_H for the COM proton P. Recall from the above that $m = (m_\tau + m_\mu + m_e) / 2 = 2\text{Proton} = (2P) / 2 = D / 2$ reduced mass of the two positron motion so

$$\frac{D}{2} 3\psi = P 3\psi = \left(\sqrt{\frac{1}{2m_\tau} \frac{dr'}{ds}} + \sqrt{\frac{1}{2m_\mu} \frac{dr'}{ds}} + \sqrt{\frac{1}{2m_e} \frac{dr'}{ds}}\right)^2 \psi \text{ stable solution: Newpde } 2P_{3/2} \text{ state at}$$

$r=r_H$.

Replace black box mass D with its interior ultrarelativistic values

Replace the mass D black box terms using Newpde $\gamma^r \sqrt{\kappa_{rr}} dr \equiv dr'$. Use $\gamma^r \gamma^r = 1$

But from eq.23 (and note τ, μ, e are Dirac equation-Newpde particles so) we can define the black box mass relativistic part: $\gamma^r \sqrt{\kappa_{rr}} dr \equiv dr'$. Use $\gamma^r \gamma^r = 1$ so that

$$\frac{D}{2} 3\psi = \frac{(m_\tau + m_\mu + m_e)}{2} 3\psi = \left(\frac{\gamma^r}{1} \sqrt{\frac{\kappa_{rr\tau}}{D_\tau/2}} \frac{dr}{ds} + \frac{\gamma^r}{1} \sqrt{\frac{\kappa_{ee\mu}}{D_\mu/2}} \frac{dr}{ds} + \frac{\gamma^r}{1} \sqrt{\frac{\kappa_{rre}}{D_e/2}} \frac{dr}{ds} \right)^2 \psi$$

Given the black box interior positron ultrarelativistic (so at 45°: $\sqrt{2}dr=ds$), $\kappa_{rr}=m^2$ for 0 speed from B10, eq.15) motion inside r_H :

$$3 \frac{(m_\tau + m_\mu + m_e)}{2} = \left(\frac{\gamma^r}{1} \sqrt{\frac{\kappa_{rr\tau}}{m_\tau/2}} \frac{dr}{\sqrt{2}dr} + \frac{\gamma^r}{1} \sqrt{\frac{\kappa_{rr\mu}}{m_\mu/2}} \frac{dr}{\sqrt{2}dr} + \frac{\gamma^r}{1} \sqrt{\frac{\kappa_{rre}}{m_e/2}} \frac{dr}{\sqrt{2}dr} \right)^2$$

so that (again $\sqrt{\kappa_{rr}} = m$):

$$3(m_\tau + m_\mu + m_e) = 2(\sqrt{m_\tau} + \sqrt{m_\mu} + \sqrt{m_e})^2 \quad \text{so}$$

$$\frac{m_\tau + m_\mu + m_e}{(\sqrt{m_\tau} + \sqrt{m_\mu} + \sqrt{m_e})^2} = \frac{2}{3} \quad \text{Koide}$$

Turns out that m_τ , m_μ , m_e move up and down together with the motion of object A zitterbewegung keeping the Koide 2/3 constant. Note these are unique solutions for $2m_p = G+1S_{1/2}+2S_{1/2}=m_e+m_\mu+m_\tau$. Also this equation is really a quartic with 3 other complex solutions. We could also use this relation to derive the value of m_t out to 7 sig.fig.(to muon mass accuracy.)

Ratios of the real valued masses that solve Koide are $m_\tau/m_\mu/m_e = 1/.05946/.0002826$, good to at least 4 significant figures.

Masses proportional to charge in $e/2m_e = g_e$, $e/(2(m_e(1+m_\mu))) = g_\mu$. Note m_μ and m_e are both changing together (as in the Mercuron equation) but the gyromagnetic ratio of the muon $g_\mu=e/2m_e(1+m_\mu)$ will change and gyromagnetic ratio of the electron $g_e=e/2m_e$ will not.

1.7 Value of C_M source term in New pde

Origin of Mass is 3 extreme Mandelbulbs

Recall postulate of 1 requires that at the end of all these derivations that $C \approx 0$. Thus we require a Fitzgerald contracted C provided by a eq.5 Minkowski metric frame of reference γ of moving the eq.7 object. From equation 3 for $N=0$ $C \approx \delta z$. So $C = \delta z / \gamma = C_M / \gamma \equiv C_M / \xi$. So that $\xi = m_e \gamma$ ($= \tau + \mu = 2m_p$ in Mandelbrot set fig.6 for *smallest* stable (so most *observable*) λ_C) in $C = C_M / \gamma = C_M / \text{mass} = r_H$ which also thereby *requires* us to define both mass $\alpha \gamma$ and charge $C_M = e^2$

Again $N=0$ equation 17 $\delta z = C_M / \xi$ satisfies extreme condition equation 3 (that is straight from the postulate) and for a (eq11) circular C $N=0$ nonflat perturbation makes an ξ observable mass (energy operator H). So that 45° extreme δz small (circle) Mandelbulb μ and the tiny antenna Mandelbulb circle then are both observable Newpde masses, so leptons. But μ is not a constant in time because of $N=1$ eq.12 angle Newpde zitterbewegung variable time t in $\delta z = e^{i\omega t}$ contribution (eq.20) to the δz chord in the small Mandelbulb of the 45° (fig6 below). In contrast the next higher energy antenna is from eq.4 quadratic equation solution at the Feigenbaum point (so it gives our 2 *fundamental extreme* excited state Mandelbulb) mass τ that does not change over cosmological time in $N=1$ allowing us to normalize it to 1). Note these are Mandelbulb radii just as eq.7-9 are in fig6, fig4, fig3 of the section 1 eq.3 application for the τ , μ respective Mandelbulb radii δz lobes in fig.6 so they *each* have their own neutrino ν . eq.7,8,9 with its electron' and neutrino still the core equations even for the muon and taunon thereby deriving the 3 generations of leptons.

Object B mass Effects (see appendix B also)

Note in appendix B the (fractally) selfsimilar to electron (ignoring zitterbewegung for the moment) Kerr metric here is rotating at near c at the equator but inertially frame drags (eg., ergosphere) to the point we see it internally (almost) only as a Schwarzschild metric. Due to the drop in inertial frame dragging caused by object B however the eq.B9 Kerr term $(a/r)^2$ is not zero anymore which (N=1 observer, N=0 observable) is equal to the $C_M/(\delta z \delta z)$ (with $r^2=|\delta z|^2$, define $a^2=C_M$) =mass= $1+\epsilon+\Delta\epsilon$ (fig6) whose Newpde fractal mass-energy- zitterbewegung frequency ω is also in the zitterbewegung exponent. We call the charge C_M which in other units and off the light cone is e^2 . Note also δz (in $C_M/(\delta z \delta z)$) is also determined by the frame of reference so by the magnitude of the Lorentz transformation γ boost of δz creating (small C) ξ input into eq.13 in $r_H = C_M/\xi$.

Note these 2D τ, μ Mandelbulbs can be on a flat 2D ($z=1$) plane or this spherical 2D shell ($z=0$)

Note the above 3e composite spherical $2P_{3/2}$ shell at $r=r_H$ is the only other stable 2D space (in addition to these $z=1$ flat 2D) Newpde ground state to define these Mandelbulbs on. Thus high energy 2D $\tau+\mu$ Mandelbulbs provide 3e stability in μ and 3e in τ so $\mu+\tau=3e+3e=(\gamma m_e + \gamma m_e)_\tau + (\gamma m_e + \gamma m_e)_\mu$ as 2 $2P_{3/2}$ orbitals with S and L inside the horizon r_H so unobserved so all that is seen from the outside is (no longer the inside 2P) net $J=S'=1/2$.

For N=0 observable For N=0 observable $z'=1+\delta z$ so z' is perturbation z .

$z'=0$, 11b, Spherical shell: $\delta z = \text{Compton wavelength } \lambda_C$ on the high energy $2P_{3/2}$ $r=r_H$ 2D spherical shell then is a domain of these same 2D Mandelbulbs μ, τ giving on the 2D shell: $\mu+\tau=3e+3e=(\gamma m_e + \gamma m_e)_\tau + (\gamma m_e + \gamma m_e)_\mu = 3e+3e = m_p + m_p$. two body motion equipartition of energy of the interacting positrons in each of two **baryons** each with $J=S'=1/2$. Eq 11b so for each positron $\delta z' = r_H = C_M/\xi_0 = C_M/m_e$ in eq.12. with C_M and ξ real numbers.

$z'=1$, 11a, Free Space: $r'_H \ll r_H$ (so not that shell) because for $z=1$ $\xi_1 \gg \xi_0$

$\delta z = \lambda = h/mc = \text{Compton wavelength}$, $2\pi r'_H = \lambda$, $m = \xi_1$. Again 3e for each of 2D free space domain high energy quasi stable μ, τ ,: $\tau+\mu=3e+3e = 2$ free space **leptons** each with $J=S'=1/2$. 11a so $\delta z = r'_H = C_M/\xi_1 = C_M/(\tau+\mu)$ (21)

For N=1 observer eq.3 implies $C = \delta z \delta z / \xi$ so that $\xi = C / \delta z \delta z = C / (\text{Mandelbulb radius})^2 = \text{mass}$ (from fig.6). or as a fraction of τ , with $2m_p = \tau + \mu + e = \xi_1$ electron $\Delta\epsilon = .00058$ (22)

Recall eq.3 $\delta z + \delta z \delta z = C$. So for N=1 observer $|\delta z| \gg 1$ so $\delta z \delta z = C$. Given eq.3 for N=0 $|\delta z| \gg |\delta z \delta z|$, ($C \approx \delta z$ sect.1 for N=0, eq10).

1.8 Postulate 0 implied finally

But γ (observer) = γ (observable) so for the N=0 observable we got the γ from the N=1 observer case in $r_H = C_M/\gamma = C_M/\xi = C$ for small C and so postulate 0. Thus we really did just **postulate 0**.

Conclusion

So the **small C** at the end was required. So we really did just **postulate 0**

So we just do *what is simplest* (let Occam be your guide), just **postulate 0**: the physics (Newpde) will then follow, top down:

Ultimate Occam's Razor postulate 0

It means here *ultimate* simplicity, the *simplest* idea imaginable. So for example $z=zz$ is *simpler* than $z=zzzz$. Therefore **0** in this context (uniquely algebraically defined by $z=zz$) is this ultimate Occam's razor **postulate**. How could you not be more 'Occam' than postulating 0? Recall that the null set \emptyset postulates absolutely nothing

Intuitive Notion (of postulate **0 \Leftrightarrow Newpde+Copenhagen stuff)**

The Mandelbrot set introduces that $r_H = C_M/\xi_1$ horizon in $\kappa_{oo}=1-r_H/r$ in the Newpde, where C_M is fractal by 10^{40} Xscale change(fig.2) So we have found (davidmaker.com) that: Given that fractal selfsimilarity astronomers are observing from the inside of what particle physicists are studying from the outside, that **ONE** *New pde* **e** electron r_H , **one** thing (fig.1). *Everything* we observe big (cosmological) and small (subatomic) is then that (*New pde*) r_H , even baryons are composite **3e**. So we understand, *everything*. This is the only Occam's razor first principles theory.

Summary: So instead of doing the usual powers of 10 simulation we do a single power of 10^{40} simulation and we are immediately back to where we started!

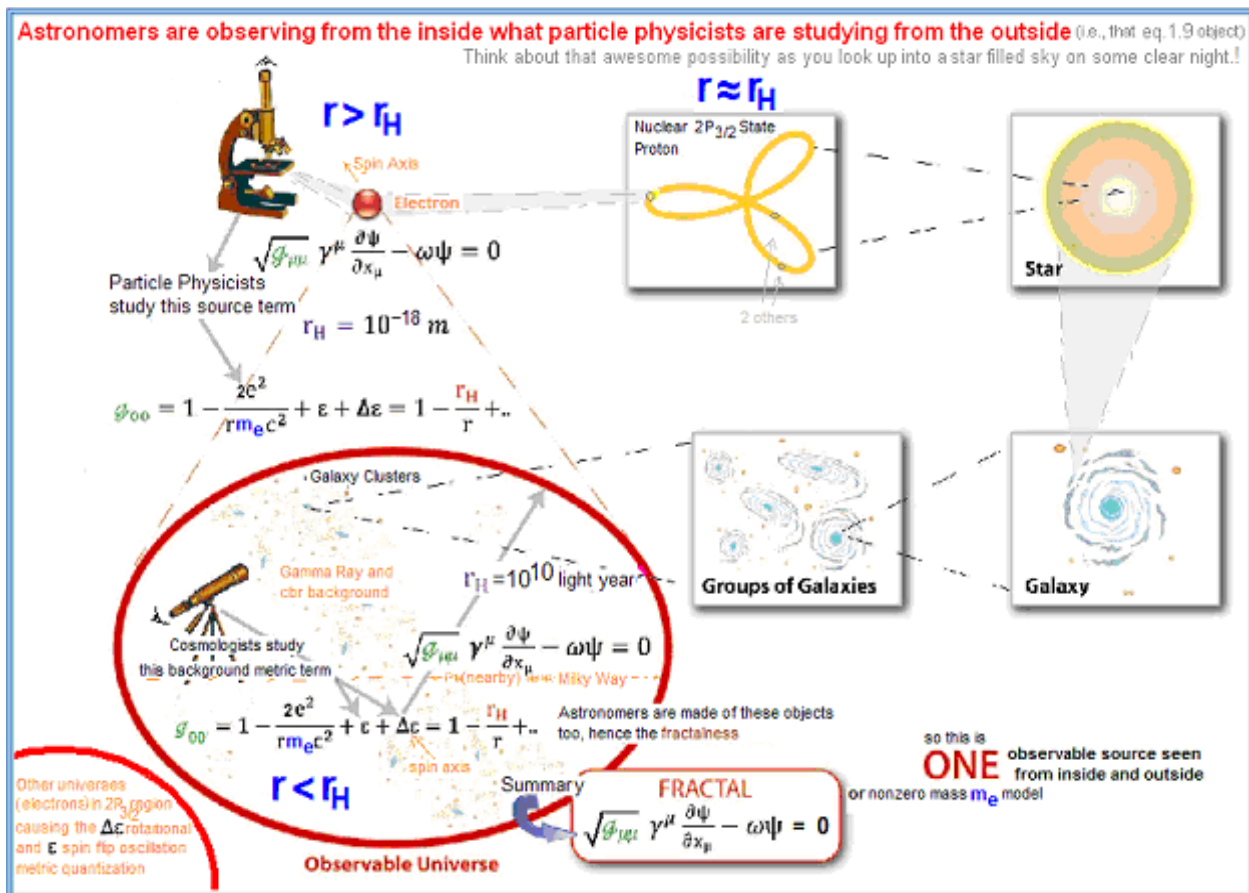


fig2.

(↑lowest left corner) Object C $g_{00} = \kappa_{00}$ caused caused metric quantization jumps: Reyleigh Taylor->galaxy->globular->protostar nebula,etc. X100 scale change metric quantization jumps (PartIII)

References

- (6) Penrose in a utube video implied that the Mandelbrot set might contain physics. Here we merely showed how to find it. For example the (fractal Mandelbulb neighborhood area $|drdt| > 0$ of the) Feigenbaum point is a subset (containing that 10^{40} Xselfsimiilar scale jump: Fig1)
- (7) Cantor: Ueber die Ausdehnung eines Satzes aus der Theorie der trigonometrischen Reihen, "Ueber eine elementare Frage der Mannigfaltigkeitslehre" Jahresbericht der Deutschen Mathematiker-Vereinigung.". Cantor proved the real# were dense with a binary # (1,0) argument (Our $z=zz$ solutions also implying 11c and appendix F). Thus we capture all the core real# properties with postulate0
- (8) Tensor Analysis, Sokolnikoff, John Wiley
- (9) The Principle of Relativity, A Einstein, Dover
- (10) Quantum Mechanics, Merzbacher, John Wiley
- (11) lemniscate circle sequence (Wolfram, Weisstein, Eric)

Direct Results From Postulate 0 Besides the Newpde

Extra Copenhagen Interpretation & 10^{82} between fractal scales stuff besides Newpde also directly from postulate 0

Quantum Mechanics core *Is* The Newpde $\psi \equiv \delta z$ (for each N fractal scale) but other stuff comes out of postulate 0 as well (as the Newpde) i.e., the Copenhagen stuff. For example recall from eq.3 for observable fractal scale $N=0$ we have $C \approx \delta z$ (A1) with C the Mandelbrot set. The interior of the inner boundary (fig3) of the electron, muon and tauon Mandelbulbs for small angle $\delta z/ds$ rotations is filled with C points so we can impose a given C^2 continuous envelope function over these points such as $\delta z^* \delta z$ and its integral over a volume V_o given by $(\int [(\delta z^* \delta z)/V_o] dV)/V_o = (\int [C^* C/V_o] dV)/V_o$ (from eq.A1) which gives a measure of the number of C s in V_o thereby implying $\delta z^* \delta z/V_o^2$ is a probability density (**in Copenhagen**). So if the number $\int [C^* C/V_o] dV/V_o$ is equal to 1 then the total probability is 1 that the electron is in V_o . So we did not have to postulate noise C for the purpose of introducing probabilities, we derived it instead given that the Mandelbrot set is plenty noisy with all those C points especially on the edges.. Also recall the solution to (postulate 1) $z=zz$ is **1,0**. Recall eq.11b that the electron is $\delta z=-1$. In $z=1-\delta z$, $\delta z^* \delta z$ is $-1^* -1=1$ and so from eq. A1 can then be interpreted as probability density, the probability of z being **0**. Recall $z=0$ is the $\xi_o=m_e$ electron solution(11b) to the new pde so $\delta z^* \delta z=1$ is the probability we have just an electron (11b). So $z=zz$ even thereby conveniently provides us with an automatic normalization of δz . Note also that $(\delta z^* \delta z)/dr$ is also then a one dimensional probability ‘density’. So Bohr’s probability density “postulate” for $\psi^* \psi$ ($\equiv (\delta z^* \delta z)$) is derived here and even contains the normalization to 1 here. So it is not a postulate anymore. (Thus Bohr was very close to the postulate of 0, and so using $z=zz$ here.). Note this result came directly out of the postulate of 0, not the Newpde.

Note also that the electron-positron eq.7 has *two* components(i.e., $dr+dt$ & $dr-dt$) that *both* solve eq.5 (and therefore eq.3) *together* as analogous to creating a $\delta z \equiv \psi = \frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle)$ singlet state relation with spin S of two opposite spin electrons $(S_1+S_2)^2 = S^2$. This singlet ψ can be used as a paradigm-model of the iconic idler-signal (Alice and Bob) singlet QM $\delta(p_A-p_B)$ conservation law state, in the Bell’s inequality functions of the idler-signal correlations.. We could then label these two parts of eq.7 *observer* and *object* with associated eq.7 wavefunctions ψ_1, ψ_2 and singlet ψ . Thus if we observe ψ_1 (idler) we must infer that there is a ψ_2 (signal from eq.7) *and* so our singlet wavefunction ψ . So we ‘collapsed’ our wavefunction to our singlet wave function ψ by observing ψ_1 since *we knew the singlet wave function* existed at the beginning (ala Bertlmann’s socks). Then apply the same mathematical reasoning to every other such analog of $\delta z \equiv \psi = \frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle)$ singlet cases (eg.,H,V polarized photon emission) and we will also have thereby derived the correlation functions in Bell’s inequalities This is then a derivation of the wave function collapse part of the **Copenhagen interpretation** of Quantum Mechanics from eq.7 and so from the first principles **postulate 0**.

But this (Copenhagen interpretation) wave function collapse is actually a trivial principle (i.e.,so it could be the wave function ψ is trivially just what you measure) except, as EPR pointed out, in this kind of conservation law singlet case laboratory initialization paradigm ψ . To (actually) know the initial S_1+S_2 in this $\delta z \equiv \psi = \frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle)$ QM singlet state is actually a **rare (laboratory setting) case** and so its spooky superluminal collapse is not a universal attribute (that being the new fad taking theoretical physics by storm) of all observed particles. So even the core Bertlmann’s socks situation is rare and without it Bell’s inequalities don’t even apply and so in that case there is no such spookiness. For the trivial single particle case we can say that measurement caused decoherence was the cause of that type of wave function collapse.

Also recall from appendix C dr^2+dt^2 is a second derivative *operator* wave equation (A1,eq.11) that holds all the way around the circle and gives the wave equation, waves. In eq.16, $N=1$ error magnitude $C \approx \delta z$ (sect.2.3) is also a $\delta z'$ angle measure on the dr, dt plane. One extremum ds ($z=0$) is at 45° so the largest C is on the diagonals (45°) where we have eq.5 extremum holding: particles. So a wide slit has high uncertainty, so large C (rotation angle) so we are at 45° (eg., particles, Newpde photoelectric effect). For a *small slit* we have less uncertainty in position so smaller C , not large enough for 45° , so only the *wave equation* $C1$ holds (then small slit diffraction). Thus we derived “wave particle duality” here. So complementarity is derived here, not postulated thereby completing the derivation of the Copenhagen interpretation.

We can count electrons and light quanta here also

Also recall wave equation eq.C1 iteration of the New pde with eq.11 operator formalism. So $dr/ds=k$ in the sect.1 circle $\delta z=dse^{i\theta}$ θ exponent kx with $k=2\pi/\lambda \equiv p/\hbar$. Multiplying both sides by \hbar with $\hbar k \equiv mv$ as before we then have the DeBroglie equation that relates particle momentum to wavelength in quantum mechanics as we already mentioned in section 1. For all the rotations in fig.4 (except the eq.11 IVth to Ist quadrants: in eq.B1 each quadrant rotation provides one derivative for each v) $\left[\left(\frac{dr+dt}{ds} \right) + \left(\frac{dr+dt}{ds} \right) \right] \delta z = 2 \frac{ds}{ds} \delta z = 2(1)\delta z$ Equation 11 (sect.1) then counts units N of each 2 half integer $S=1/2$ angular momentums= 1 unit oelectrons (spin1 for W and Z) off the light cone. For the rotation in the eq.11 IVth to Ist quadrants (each quadrant rotation provides one derivative for each v) at 45° $dr=dt$ (on the light cone in fig.4) so for Hamiltonian H : $2H\delta z=2(dt/ds)\delta z=2(1/2)\delta z=(1)\hbar\omega\delta z=\hbar k\delta z$ on the diagonal so that $E=pv=\hbar\omega$ for the two v energy components, universally. Thus we can state the most beautiful result in physics that $E=Nhf$ for the energy of light with N equal N monochromatic photons. Thus this eq.11c merely counts the number of electrons. It is not list of energy levels (states) as in the (well known) quantization of the energy levels N of the E&M field with SHM.

By the way the Casimir force is simply then the relativistic component of the Van der Waals force, has nothing to do with zero point energy vacuum fluctuations. See Robert Jaffe paper from 2005.

Appendix A Object A, the cosmological positron we are inside of

A Extra fractal zoom stuff besides Newpde directly from postulate0 such as 10^{82} objects (so including objects A,B,C) between fractal scales

The Fiegenbaum point (11a) is the only part of the Mandelbrot set we zoom from. At the Fiegenbaum point (imaginary) time $X10^{-40}=\Delta$ and real -1.40115 (sect.1.2). At the very beginning (top) C was defined to be constant *only* at $C \approx 0$ ($\|C\| \ll 1$). So at the end of all these derivations we still have to have a small C . This implies a boosted SR Lorentz transformation universal reference frame to random (since this transformation drops (cancels) noise C in eq.2, fig6), small C_M subset $C \approx \delta z'$ (from eq.3) =real distance =real $\delta z/\gamma = 1.4011/\gamma = C_M/\gamma \equiv C_M/\xi_1$ using large ξ_1 . Note at the Fiegenbaum point distance $1.4011/\gamma$ shrinks a lot but time $X10^{-40}\gamma$ doesn't get much bigger since it was so small to begin with at the Fiegenbaum point. Eq.1 then means we have Ockam's razor optimized **postulated 0**. Given the New pde r_H we only see the $r_H=e^2 10^{40N}/m$ with 10^{82} sources from our $N=0$ observer baseline. We never see the $r < r_H$

<http://www.youtube.com/watch?v=0jGai087u3A> which explores the Mandelbrot set interior near the Fiegenbaum point. Reset the zoom start at such extremum $S_N C_M = 10^{40N} C_M$ in eq.13. The splits are in 3 directions from the orbs. There appear to be about 2.5 splits going by each second (given my PC baud rate) and the next Mandelbrot set comes up in about 62 seconds. So

$3^{2.7 \times 62} = 10^N$ so $172 \log 3 = N = 82$. So there are 10^{82} splits. So there are about 10^{82} splits per initial split. But each of these Mandelbrot set Feigenbaum points is a $C_M/\xi = r_H$ in electron (eq.10 above). So for each larger electron there are **10^{82} constituent electrons**. Also the scale difference between Mandelbrot sets as seen in the zoom is about **10^{40} , the scale change** between the classical electron radius and 10^{11} ly with the C noising giving us our fractal universe.

Recall again we got from eq.3 $\delta z + \delta z \delta z = C$ with quadratic equation result:

$\delta z = \frac{-1 \pm \sqrt{1-4C}}{2}$. is real for noise $C < 1/4$ creating our noise on the $N=0$ th fractal scale. So $1/4 = (3/2)kT/(m_p c^2)$. So T is 20MK. So here we have *derived the average temperature of the universe* (stellar average). That $z' = 1 + \delta z$ substitution also introduces Lorentz transformation rotational and translation noise that does not effect the number of splits, analogous to how a homeomorphism does not change the number of holes (which is a Topological invariant). So the excess C noise (due to that small C' boost) causes the Feigenbaum point neighborhood internal structure to become randomized (as our present universe is) but the number of electrons (10^{82}) remains invariant. See appendix D mixed state case2 for further organizational effects. $N = r^D$. So the **fractal dimension** = $D = \log N / \log r = \log(\text{splits}) / \log(\#r_H \text{ in scale jump})$
 $= \log 10^{80} / \log 10^{40} = \log(10^{40})^2 / \log(10^{40}) = 2$. (See appendix E for Hausdorff dimension & measure) which is the same as the 2D of eq.4 and the Mandelbrot set. The next smaller (subatomic) fractal scale $r_1 = r_H = 2e^2/m_e c^2$, $N=0$ th, $r_2 = r_H = 2GM/c^2$ is defined as the $N=1$ th where $M = 10^{82} m_e$ with $r_2 = 10^{40} r_1$ So the Feigenbaum pt. gave us a lot of physics:

eg. **#of electrons in the universe, the universe size, temp.** With 10^{82} electrons between any two fractal scales we are also *certainly allowed objects B&C* in the Newpde $2P_{2/3}$ state at $r = r_H$.

time independent perturbation

N=1 observer (eq.17,18,19 gives our Newpde metric $\kappa_{\mu\nu}$ at $r < r_H, r > r_H$)
 Found General Relativity (GR) GR from eq.17- eq.19 so Schwarchild metric and so can do a dyadic coordinate transformation on it to get the Kerr metric and all these free space metrics to get all the solutions to $R_{ij} = 0$. $N = -1$, $e^2 10^{40(-1)} = e^2 / 10^{40} = G m_e^2$, solve for G, get GR. So we can now write the Ricci tensor R_{uv} (and fractally self similar perturbation Kerr metric since frame dragging decreased by external object B, sect.B2). Also for fractal scale $N=0$, $r_H = 2e^2/m_e c^2$, and for $N=-1$ $r'_H = 2G m_e / c^2 = 10^{-40} r_H$. **D=5 if using N=-1, and N=0, N=1 contributions in same $R_{ij} = 0$**
 Note the $N=-1$ (GR) is yet another δz perturbation of $N=0$ $\delta z'$ perturbation of $N=1$ observer thereby adding, if these scales share the same time coordinate, at least 1 independent parameter dimension to our $\delta z + (dx_1 + idx_2) + (dx_3 + idx_4)$ (4+1) *explaining why Kaluza Klein 5D $R_{ij} = 0$ works so well*: GR is really 5D if E&M

Appendix object A Fractal mass and cosmology

From Newpde (eg., eq.1.13 Bjorken and Drell special case) $i\hbar \frac{\partial \psi}{\partial t} = \frac{\hbar c}{i} \left(\alpha_1 \frac{\partial \psi}{\partial x^1} + \alpha_2 \frac{\partial \psi}{\partial x^2} + \alpha_3 \frac{\partial \psi}{\partial x^3} \right) + \beta m c^2 \psi = H \psi$. For electron at rest: $i\hbar \frac{\partial \psi}{\partial t} = \beta m c^2 \psi$ so: $\delta z = \psi_r = w^r(0) e^{-i \epsilon_r \frac{m c^2}{\hbar} t}$
 $\epsilon_r = +1, r=1,2; \epsilon_r = -1, r=3,4$): This implies an oscillation frequency of $\omega = m c^2 / \hbar$. which is fractal here ($\omega = \omega_0 10^{-40N}$). So the eq.16 the 45° line has this ω oscillation as a (that eq.7-9 δz variation) rotation at radius ds. On our own fractal cosmological scale $N=1$ we are in the expansion stage of one such oscillation. Thus the fractalness of the Newpde explains cosmology. The next higher cosmological scale is independent (but still connected by relativistic superposition of speeds

implying a inverse separation of variables result: $i\hbar \frac{\partial \psi}{\partial t} = \beta \sum_N (10^{-40N} (\omega t)_{\varepsilon+\Delta\varepsilon}) \psi = \beta \sum_N (10^{-40N} m_{\varepsilon+\Delta\varepsilon} c^2 / \hbar) \psi$. Note this means that fractal scale $N=1$ the 45° small Mandelbulb chord ε (Fig6) is now, given this ω , getting larger with time so $1-t \propto \varepsilon$. (See Mercuron equation) B3a. But the fig6 Mandelbulb antenna taun is stationary so its mass can be set to 1. So at this time (relative to the taun) the muon $\varepsilon=.05946$, electron $\Delta\varepsilon=.0005899/2=.0002826$. (B1)

Set average $e^{(-\varepsilon+\Delta\varepsilon)t} = \delta |e^{i\tau tz}|$ Newpde zitterbewegung oscillation but τ constant (fig6), doesn't vary in cosmological time t_c . So cosmologically (eq. B11) outside r_H of object B for $N=0$ use t_z . For $N=1$ use t_c for cosmologically relevant time dependence.

Define average $(e^{i(\tau+\varepsilon+\Delta\varepsilon)t_z}) \equiv \delta \bar{z}_0$, So $|\delta z| = |e^{-i\varepsilon_r \frac{mc^2}{\hbar} t} \delta \bar{z}_0| = \delta \bar{z}_0 e^{i\omega t} = e^{i(\tau+\varepsilon+\Delta\varepsilon)t_z + i(-\varepsilon+\Delta\varepsilon)t_c} = \delta \hat{z}_0 e^{i(\varepsilon+\Delta\varepsilon)t} = \delta \hat{z}_0 \sqrt{\kappa_{rr}}$ in $(\sqrt{\kappa_{rr}} = e^{i(\varepsilon+\Delta\varepsilon)t})$ $dr^2 = \kappa_{rr} dr^2 = e^C \kappa_{00} dr^2 = e^{i(-\varepsilon+\Delta\varepsilon)t} \kappa_{00} dr^2$ (A2)

But seen from inside at $N=1$ $E=1/\sqrt{\kappa_{00}}=1/\sqrt{(1-r_H/r)}$ (B20) then $r < r_H$ & E becomes imaginary in $e^{iEt/\hbar} = \delta z = \sqrt{\kappa_{00}} dt = e^{-i\varepsilon_r \frac{mc^2}{\hbar} t} dt \rightarrow e^{(-\varepsilon+\Delta\varepsilon)t}$ (A2)

The negative sign from equation B2a below. The reduced mass ground state rotater ($\Delta\varepsilon$) for ε for this κ_{00} part of derivation). This $e^{i2\Delta\varepsilon/(1-2\varepsilon)} = \kappa_{00}$ asymptotic value must be equal to g_{00} in galaxy halos in the plane of the galaxy (sect.11.4). Ricci tensor is given by oscillating source.

‘Observer’ scale $N > M$ ‘observables’ scale.

Recall from sect.1 if our scale $N > M$ for some object then N is the observer scale and M is the ‘observable’ scale. Note the scale difference can be very small. Since we are all electrons that means a slightly smaller scale electron is the observable. But this seems to eliminate astronomy as observation of ‘observables’ since those objects exist at a *larger* scale $N=1$. But not to the $N=2$ scale (the ‘gid’ scale as I call it) since to him ($N=2$) the $N=1$ astronomy scale is an ‘observable’ scale since $N=2 > N=1$.

A2 Two perturbations of the $N=1$ scale as seen by $N=2$

We also have two perturbations of the $N=1$ scale here. The first perturbation is due to the Dirac equation object A zitterbewegung harmonic oscillation (which equivalently could be the source or the manifold). Recall in that regard Weinberg (eg., eq 10.1.9 “Gravitation & Cosmology”) calls it a “harmonic coordinate system” (here as eq.1.13 Bjorken and Drell) thereby also providing our manifold in that 2nd case. The second much smaller perturbation is due to the drop in inertial frame dragging due to nearby object B.

Ricci tensor source term for interior to object A

In that regard the Ricci tensor $= R_{ij} = -1/2 \Delta(g_{ij})$ (where Δ is the Laplace-Beltrami second derivative operator) is not zero and the right side is the metric source. Recall limit R_{ij} as $r \rightarrow 0$ is the source, where alternatively gravity creates gravity feedback loop in the Einstein equations which becomes the modulation of the DeSitter ball implied by the zitterbewegung oscillation of object A. Geometrically, the Ricci curvature is the mathematical object that controls the (comoving *observer*) *growth rate of the volume* of metric balls in a manifold in this case given by the New pde source zitterbewegung. Thus the above Laplace Beltrami source eq. A2 $-\sin\omega t = -\sin\mu \approx -\sin\varepsilon$ here comes out of the Newpde zitterbewegung A2.

$N=2$ ‘observer’ sees what we see if $i \rightarrow 1$ in $\sin\mu \rightarrow \sinh\mu$ in $R_{22} = -\sinh\mu$: which makes our $N=1$ ‘observables’.

But $R_{22} = e^{-\lambda} [1 + 1/2 r(\mu' - \nu')] - 1$ with $\mu = \nu$ (spherical symmetry) and $\mu' = -\nu'$. So as $r \rightarrow 0$, $\text{Im} R_{22} = \text{Im}(e^\mu - 1) = \mu + .. = \sin\mu = \mu + ..$ for outside r_H imaginary μ for small r (at the source) so

zitterbewegung $\sin\mu$ becomes a gravitational source (alternatively gravity itself can create gravity in a feedback mechanism). The $N=2$ observer then multiplies by iR_{22} , $-i\sin\mu$ and μ to get $R_{22}=-\sinh\mu$ (A2A)

to see what the $N=2$ observer sees that we see inside r_H so:

$$R_{22}=e^{-\nu}[1+\frac{1}{2}r(\mu'-\nu')]-1=-\sinh\nu=(-(e^\nu-e^{-\nu})/2), \quad \nu'=-\mu' \text{ so}$$

$$(e^\mu-1)=-\sinh\mu \text{ for positive } \mu \text{ in } \sinh\mu \text{ then the } \mu=\varepsilon \text{ in the } e^\mu \text{ on the left is negative} \quad (\text{A2B}).$$

Object B mostly contributes to μ' in $-r\mu\omega$, with object C providing a tiny perturbation of μ' , implying there is no such positive $\sinh\mu$ constraint for object C. Thus the object C perturbation μ_c in e^{μ_c} coefficient can be positive or negative

$$e^{-\mu}[-r(\mu')]=-\sinh\mu-e^{-\mu}+1=(-(e^{-\mu}+e^\mu)/2)-e^{-\mu}+1=(-(e^{-\mu}+e^\mu)/2)+1=-\cosh\mu+1. \text{ So given } \nu'=-\mu'$$

$$e^{-\nu}[-r(\mu')]=1-\cosh\mu. \text{ Thus}$$

$$e^{-\mu}r(d\mu/dr)=1-\cosh\mu$$

$$\text{This can be rewritten as:} \quad e^\mu d\mu/(1-\cosh\mu)=dr/r$$

We set the phase μ so that when $t=0$ then $r=0$ so use $r=\sin\omega t$ in eq.A1. Given the fractal universe a temporarily comoving proper frame at minimum radius lowest γ must imply a μ Mandelbulb chord 45° intersection that implies minimally the Newpde ground state (Which can't go away analogously as for a hydrogen atom orbital electron.) $\Delta\varepsilon$ electron for comoving outside observer where then at time=0, in B1,B2 $\tau-\varepsilon\approx\omega t=\Delta\varepsilon \approx 1-1=0$ so that $\omega t=\Delta\varepsilon$ when $\sin\omega t\approx 0$. So the integration of B3 is from $\xi_1=\mu=\varepsilon=1$ to the present day mass of the $\mu=\mu_{\text{muon}}=.05946$ (X tauon mass) giving us:

$$\ln(r_{M+1}/r_{bb})+2=[1/(e^\mu-1)-\ln[e^\mu-1]]^2 \quad (\text{A3C})$$

implying $gr=e/2me(1+\mu)$ gyromagnetic ratio ($\mu=m$) is changing with time as was discovered recently at Fermi lab 2023 (Ch.7) with CERN 1974 gr muon data for comparison.

A2 Writing The feedback mechanism two different ways

Introduction to $\Delta\varepsilon$ contribution to what $N=2$ sees

We have two perturbations, one due to the zitterbewegung and a smaller one due to the drop in inertial frame dragging due to nearby objectB.

So inside object A we can include the zitterbewegung oscillation $\delta z=\sqrt{\kappa_{oo}}dt=e^{-i\varepsilon r\frac{mc^2}{\hbar}t}dt \rightarrow e^{(-\varepsilon+\Delta\varepsilon)^2}dt$ in the source as $-\sinh\mu=R_{22}$

Alternatively zitterbewegung oscillation $\delta z=\sqrt{\kappa_{oo}}dt=e^{-i\varepsilon r\frac{mc^2}{\hbar}t}dt \rightarrow e^{(-\varepsilon+\Delta\varepsilon)^2}dt$, with $r\rightarrow\infty$, $g_{\alpha\alpha}\rightarrow\text{constant}\neq 1$, can be the manifold itself, so relative to this manifold the motion is flat space so sourceless. Thereby we set $R_{22}=-\sinh\mu=0$ with $R_{\alpha\alpha}=0$.

So these 2 perturbations then give the $N=1$ contribution to what $N=2$ sees.

$N=2$ sees local interior contribution of object A

Object B $N=1$ ambient metric $C=\text{constant}$ (nonrotating)

From eqs17-18 but with ambient metric ansatz: $ds^2=-e^\lambda(dr)^2-r^2d\theta^2-r^2\sin\theta d\phi^2+e^\mu dt^2$ (A3)

so that $g_{oo}=e^\mu$, $g_{rr}=e^\lambda$. From eq. $R_{ij}=0$ for spherical symmetry in free space and $N=0$

$$R_{11}=\frac{1}{2}\mu''-\frac{1}{4}\lambda'\mu'+\frac{1}{4}(\mu')^2-\lambda'/r=0 \quad (\text{A4})$$

$$R_{22}=e^{-\lambda}[1+\frac{1}{2}r(\mu'-\lambda')]-1=0 \quad (\text{A5})$$

$$R_{33}=\sin^2\theta\{e^{-\lambda}[1+\frac{1}{2}r(\mu'-\lambda')]-1\}=0 \quad (\text{A6})$$

$$R_{oo}=e^{\mu-\lambda}[-\frac{1}{2}\mu''+\frac{1}{4}\lambda'\mu'-\frac{1}{4}(\mu')^2-\mu'/r]=0 \quad (\text{A7})$$

$$R_{ij}=0 \text{ if } i\neq j$$

(eq. A4-A7 from pp.303 Sokolnikof(8)): Equation B4 is a mere repetition of equation A6. We thus have only three equations on λ and μ to consider. From equations A4, A7 we deduce that $\lambda' = -\mu'$ so that radial $\lambda = -\mu + \text{constant} = -\mu + C$ where C represents a possible \sim constant ambient metric contribution which (allowing us to set $\sinh\mu=0$) could be imaginary in the case of the slowly oscillating ambient metric of nearby object B from B2. So $e^{-\mu+C} = e^\lambda$. Then A3-A7 can be written as:

$$e^{-C} e^\mu (1+r\mu') = 1. \quad (\text{A9})$$

Set $e^\mu = \gamma$. So $e^{-\lambda} = \gamma e^{-C}$ ε and $\Delta\varepsilon$ are time dependent. So integrating this first order equation (equation A9) we get:

$\gamma = -2m/r + e^C \equiv e^\mu = g_{00}$ and $e^{-\lambda} = (-2m/r + e^C) e^{-C} = 1/g_{rr}$
 or $e^{-\lambda} = 1/\kappa_{rr} = 1/(1-2m'/r)$, $2m/r + e^C = \kappa_{00}$. With (reduced mass ground state rotater ($\Delta\varepsilon$) for charged if $-\varepsilon$) dr zitterbewegung from B1 $\kappa_{rr} dr^2 = e^C \kappa_{00} dr'^2 = e^{i(-\varepsilon+\Delta\varepsilon)^2} \kappa_{00} dr^2$ from A2. We found

$$\kappa_{00} = e^C - 2m/r = e^{i(-\varepsilon+\Delta\varepsilon)^2} - 2m/r \quad (\text{A10})$$

$\Delta\varepsilon$ here is reduced ground state mass $\Delta\varepsilon$ as in Schrodinger eq $E = \Delta\varepsilon = 1/\sqrt{\kappa_{00}}$. (A10a)

does not add anything to r_H/r in κ_{rr} since e^C is not added to r_H/r there.

Appendix B Object B

Add Perturbative Kerr rotation $(a/r)^2$ to r_H/r in κ_{rr} Here nothing gets added to r_H/r in κ_{00}

Our new pde has spin S and so the self similar ambient metric on the N=0 th fractal scale is the Kerr metric which contains those ambient metric **perturbation rotations** ($d\theta/dt$ T violation so (given CPT) then **CP violation**) due to object B

$$ds^2 = \rho^2 \left(\frac{dr^2}{\Delta} + d\theta^2 \right) + (r^2 + a^2) \sin^2 \theta d\phi^2 - c^2 dt^2 + \frac{2mr}{\rho^2} (a \sin^2 \theta d\theta - c dt)^2, \quad (\text{B1})$$

where $\rho^2(r, \theta) \equiv r^2 + a^2 \cos^2 \theta$; $\Delta(r) \equiv r^2 - 2mr + a^2$, In our 2D $d\phi=0$, $d\theta=0$ Define:

$$\left(\frac{r^2 + a^2 \cos^2 \theta}{r^2 - 2mr + a^2} \right) dr^2 + \left(1 - \frac{2m}{r^2 + a^2 \cos^2 \theta} \right) dt^2 \quad \theta \neq 0$$

$\rho^2(r, \theta) \equiv r^2 + a^2 \cos^2 \theta$; $\Delta(r) \equiv r^2 - 2mr + a^2$, $r^{\wedge 2} \equiv r^2 + a^2 \cos^2 \theta$, $r'^2 \equiv r^2 + a^2$. Inside r_H $a \ll r$, $r \gg 2m$

$$\left(\frac{(r^{\wedge})^2}{(r')^2 - 2mr} \right) dr^2 + \left(1 - \frac{2mr}{(r^{\wedge})^2} \right) dt^2 + .. = \left(\frac{1}{(r')^2 - \frac{2mr}{(r^{\wedge})^2}} \right) dr^2 + \left(1 - \frac{2mr}{(r^{\wedge})^2} \right) dt^2.$$

The (r^{\wedge}/r') ² term is $\frac{(r')^2}{(r^{\wedge})^2} = \frac{r^2 + a^2}{r^2 + a^2 \cos^2 \theta} = \frac{1 + \frac{a^2}{r^2}}{1 + \frac{a^2}{r^2} \cos^2 \theta} \approx 1/g_{rr} (\approx g_{00})$

$$= \left(1 + \frac{a^2}{r^2} \right) \left(1 - \frac{a^2}{r^2} \cos^2 \theta \right) + .. = 1 - \frac{a^4}{r^4} \cos^2 \theta - \frac{a^2}{r^2} \cos^2 \theta + \frac{a^2}{r^2} + .. = 1 + \frac{a^2}{r^2} (1 - \cos^2 \theta) + .. = 1 + \frac{a^2}{r^2} \sin^2 \theta + .. \equiv 1 + \left(\frac{a}{r} \right)^2 u^2 = \left(\text{from fig. 6 mass} = \frac{C_M}{\delta z \delta z} \right) = 1 + 2(\varepsilon + \Delta\varepsilon) + .. \quad (\text{B2})$$

since $\varepsilon + \Delta\varepsilon$ are time dependent, and add $2m/r$ to this $1 + \varepsilon + \Delta\varepsilon$ at the end. $\Delta\varepsilon$ is total

(Mandelbulb) mass as in $C_M/(\delta z \delta z) = (a/r)^2$ in fig6 contributing to inertial frame dragging drop

We can normalize out $1 + \varepsilon$ over a region we know it is (at least appromately) a constant. That in turn makes the metric coefficients at $r \gg 0$ flat which is what they should be. In contrast rotation adds to κ_{rr} (B2) and only oblates $2m/r$ in κ_{00} .

Summary: Our Newpde metric including the effect of object B (with $\tau + \mu = 2m_p = \xi_1$.) is for the $\tau + \mu + e$ Mandelbulbs in Fig6

$$\tau + \mu \text{ in free space } r_H = e^2 10^{40(0)} / 2m_p c^2, \kappa_{00} = e^{i(2\Delta\varepsilon/(1-2\varepsilon))} \cdot r_H/r, \kappa_{rr} = 1 + 2\Delta\varepsilon/(1+\varepsilon) - r_H/r \text{ Leptons} \quad (\text{B3})$$

$\tau+\mu$ on $2P_{3/2}$ sphere at $r_H=r$, $r_H=e^2 10^{40(0)}/2m_e c^2$, comoving with $\gamma=m_p/m_e$. Baryons, part2 (B4)
Imaginary $i\Delta\varepsilon$ in this cosmological background metric $\kappa_{00}=e^{i\Delta\varepsilon}$ B13 makes no contribution to the Lamb shift but is the core of part III cosmological application $g_{00}=\kappa_{00}$ of eq B13 of this paper.

B1 N=0 eq.B3 Application example: anomalous gyromagnetic ratio Separation Of Variables On New Pde.

After separation of variables the “r” component of Newpde can be written as:

$$\left[\left(\frac{dt}{ds} \sqrt{\kappa_{00}} m_p \right) + m_p \right] F - \hbar c \left(\sqrt{\kappa_{rr}} \frac{d}{dr} + \frac{j+3/2}{r} \right) f = 0 \quad B5$$

$$\left[\left(\frac{dt}{ds} \sqrt{\kappa_{00}} m_p \right) - m_p \right] f + \hbar c \left(\sqrt{\kappa_{rr}} \frac{d}{dr} - \frac{j-1/2}{r} \right) F = 0. \quad B6$$

Using the above Dirac equation component we find the anomalous gyromagnetic ratio Δgy for the spin polarized $F=0$ case. Recall the usual calculation of rate of the change of spin S gives $dS/dt \propto m \propto gyJ$ from the Heisenberg equations of motion. We note that $1/\sqrt{\kappa_{rr}}$ rescales dr in $\left(\sqrt{\kappa_{rr}} \frac{d}{dr} + \frac{j+3/2}{r} \right) f$ in equation B5 with κ_{rr} from B3. Thus to have the same rescaling of r in the second term we must multiply the second term denominator (i.e., r) and numerator (i.e., $J+3/2$) each by $1/\sqrt{\kappa_{rr}}$ and set the numerator ansatz equal to $(j+3/2)/\sqrt{\kappa_{rr}} = 3/2 + J(gy)$, where gy is now the gyromagnetic ratio. This makes our equation B5, B6 compatible with the standard Dirac equation allowing us to substitute the gy into the Heisenberg equations of motion for spin S : $dS/dt \propto m \propto gyJ$ to find the correction to dS/dt . Thus again:

$$\begin{aligned} [1/\sqrt{\kappa_{rr}}](3/2 + J) &= 3/2 + Jgy, \text{ Therefore for } J = 1/2 \text{ we have:} \\ [1/\sqrt{\kappa_{rr}}](3/2 + 1/2) &= 3/2 + 1/2 gy = 3/2 + 1/2(1 + \Delta gy) \end{aligned} \quad B7$$

Then we solve for Δgy and substitute it into the above dS/dt equation.

Thus solve eq. B7 with Eq.B1 values in $\sqrt{\kappa_{rr}} = 1/\sqrt{(1+2\Delta\varepsilon/(1+\varepsilon))} = 1/\sqrt{(1+2\Delta\varepsilon/(1+0))} = 1/\sqrt{(1+2 \times 0.0002826/1)}$. Thus from equation B1:

$[\sqrt{(1+2 \times 0.0002826)}](3/2 + 1/2) = 3/2 + 1/2(1 + \Delta gy)$. Solving for Δgy gives anomalous **gyromagnetic ratio correction of the electron** $\Delta gy = .00116$.

If we set $\varepsilon \neq 0$ (so $\Delta\varepsilon/(1+\varepsilon)$) instead of $\Delta\varepsilon$ in the same κ_{00} in Newpde we get the anomalous **gyromagnetic ratio correction of the muon** in the same way.

Composite 3e: Meisner effect For B just outside r_H . (where the zero point energy particle eq. 9.22 is $.08 = \pi^\pm$) See A4

Composite 3e CASE 1: Plus $+r_H$, therefore is the proton + charge component. Eq.A4: $1/\kappa_{rr} = 1 + r_H/r_H + \varepsilon = 2 + \varepsilon$. $\varepsilon = .08$ (eq.9.22). Thus from eq.B17 $\sqrt{2 + \varepsilon}(1.5 + .5) = 1.5 + .5(gy)$, $gy = 2.8$

The gyromagnetic ratio of the proton

Composite 3e CASE 2: negative r_H , thus charge cancels, zero charge:

$$\begin{aligned} 1/\kappa_{rr} = 1 - r_H/r_H + \varepsilon = \varepsilon \quad \text{Therefore from equation B7 and case 1 eq.A3 } 1/\kappa_{rr} = 1 - r_H/r_H + \varepsilon \\ \sqrt{\varepsilon}(1.5 + .5) = 1.5 + .5(gy), \quad gy = -1.9. \end{aligned}$$

the **gyromagnetic ratio of the neutron** with the other charged and those ortho neutral hyperon magnetic moments scaled using their masses by these values respectively.

A4 eq.B3 κ_{00} application example: Lamb shift

After separation of variables the “r” component of Newpde can be written as

$$\left[\left(\frac{dt}{ds} \sqrt{\kappa_{00}} m_p \right) + m_p \right] F - \hbar c \left(\sqrt{\kappa_{rr}} \frac{d}{dr} + \frac{j+3/2}{r} \right) f = 0 \quad B8$$

$$\left[\left(\frac{dt}{ds} \sqrt{\kappa_{00}} m_p \right) - m_p \right] f + \hbar c \left(\sqrt{\kappa_{rr}} \frac{d}{dr} - \frac{j-1/2}{r} \right) F = 0. \quad B9$$

Comparing the flat space-time Dirac equation to the left side terms of equations B8 and B9:

$$(dt/ds)\sqrt{\kappa_{00}}=(1/\kappa_{00})\sqrt{\kappa_{00}}=(1/\sqrt{\kappa_{00}})=\text{Energy}=E \quad \text{B10}$$

We have normalized out the e^c in equation B10 to get the pure measured r_H/r coupling relative to a laboratory flat background given thereby in that case by κ_{00} under the square root in equation B10.

Note for electron motion around hydrogen proton $mv^2/r=ke^2/r^2$ so $KE=1/2mv^2=(1/2)ke^2/r=PE$ potential energy in $PE+KE=E$. So for the electron (but not the tauon or muon that are not in this orbit) $PE_e=1/2e^2/r$. Write the hydrogen energy and pull out the electron contribution B10a. So in eq.B2 and B8 $r_H=(1+1+.5)e^2/(m_\tau+m_\mu+m_e)/2=2.5e^2/(2m_p c^2)$. B11

Variation $\delta(\psi^*\psi)=0$ At $r=n^2a_0$

Next note for the variation in $\psi^*\psi$ is equal to zero at maximum $\psi^*\psi$ probability density where for the hydrogen atom is at $r=n^2a_0=4a_0$ for $n=2$ and the $\psi_{2,0,0}$ eigenfunction. Also recall eqA4 eq.11 $a\xi_1=m_L c^2=(m_\tau+m_\mu+m_e)c^2=2m_p c^2$ normalizes $1/2ke^2$ (Thus divide $\tau+\mu$ by 2 and then multiply the whole line by 2 to normalize the $m_e/2$.result. $\epsilon=0$ since no muon ϵ here.): Recall in eq.11a ξ_0 has to be pulled in a Taylor expansion as an operator since it a separate observable. So substituting eqs.B1 for κ_{00} , values in eq.B10:

$$E_e = \frac{(tauon+muon)(\frac{1}{2})}{\sqrt{1-\frac{r_H}{r}}} - (tauon + muon + PE_\tau + PE_\mu - m_e c^2) \frac{1}{2} =$$

$$2(m_\tau c^2 + m_\mu c^2) \frac{1}{2} + 2 \frac{m_e c^2}{2} + 2 \frac{2.5e^2}{2r(m_L c^2)} m_L c^2 - 2 \frac{2e^2}{2r(m_L c^2)} m_L c^2 - 2 \frac{3}{8} \left(\frac{2.5e^2}{r m_L c^2} \right)^2 m_L c^2$$

$$- 2(m_\tau c^2 + m_\mu c^2) \frac{1}{2}$$

$$= \frac{2m_e c^2}{2} + 2 \frac{e^2}{4r} - 2 \frac{3}{8} \left(\frac{2.5}{r m_L c^2} \right)^2 m_L c^2 = m_e c^2 + \frac{e^2}{2r} - 2 \frac{3}{8} \left(\frac{2.5e^2}{r m_L c^2} \right)^2 m_L c^2$$

So: $\Delta E_e = 2 \frac{3}{8} \left(\frac{2.5}{r m_L c^2} \right)^2 m_L c^2 =$ (Third order $\sqrt{\kappa_{\mu\mu}}$ Taylor expansion term)=

$$\Delta E = 2 \frac{3}{8} \left[\frac{2.5(8.89 \times 10^9)(1.602 \times 10^{-19})^2}{(4(.53 \times 10^{-10}))^2((1.67 \times 10^{-27})(3 \times 10^8))^2} \right]^2 (2(1.67 \times 10^{-27})(3 \times 10^8))^2$$

$$= hf = 6.626 \times 10^{-34} \times 27,360,000 \text{ so that } f = 27 \text{ MHz Lamb shift.}$$

The other 1050Mhz comes from the zitterbewegung cloud.

Note: Need infinities if **flat space Dirac 1928 equation**. For flat space $\partial g_{ik}/\partial x^j = 0$ as a limit. Then must take field $g^{km} = 1/0 = \infty$ to get finite Christoffel symbol $\Gamma^{m}_{ij} \equiv (g^{km}/2)(\partial g_{ik}/\partial x^j + \partial g_{jk}/\partial x^i - \partial g_{ij}/\partial x^k) = (1/0)(0) = \text{undefined}$ but still implying *nonzero* acceleration on the left side of the

geodesic equation: $\frac{d^2 x^\mu}{ds^2} = -\Gamma^{\mu}_{\nu\lambda} \frac{dx^\nu}{ds} \frac{dx^\lambda}{ds}$ So we need infinite fields for flat space. Thus QED

requires (many such) infinities. But we have in general curved space $g_{ij} = \kappa_{ij}$ in the New pde so do not require that anything be infinite and yet we still obtain for the third order Taylor expansion term of $\sqrt{\kappa_{\mu\nu}}$ the Lamb shift and anomalous gyromagnetic ratio correction (see above sections B3,B4).

So renormalization is a perturbative way (given it's flat space Dirac equation and minimal interaction gauge origins) of calculating these (above) same, *NON*perturbative results, it's a perturbative GR theory. But renormalization gives lots of wrong answers too, eg., 10^{96} grams/cm³ vacuum density for starters. (So we drop it here since we don't need it any longer for the high

precision QED results.) In contrast note near the end of reference 5 our $G_{00}=0$ for a 2D MS. Thus a vacuum really is a vacuum. Also that large $\xi_1=\tau(1+\epsilon')$ in r_H in eq.B13,11a is the reason leptons appear point particles (in contrast to the small ξ_0 in the composite 3e baryons).

B5 eq.B3 κ_{00} application example: metric quantization from $g_{00}=\kappa_{00}$

Given the subatomic fractal scale is dominated by quantum mechanics phenomena in a fractal universe the next higher $N=1$ fractal scale should bring the QM back: In galaxy halos $g_{00}=\kappa_{00}$ (eq.4.13) with resulting Metric Quantization $N=1$ result $g_{00}=\kappa_{00}$, in galaxy halos (eg.,replacing need for dark matter Note we have yet to use the $e^{i(2\Delta\epsilon/(1-2\epsilon))}$ in $\kappa_{00}=e^{i(\Delta\epsilon/(1-2\epsilon))}-r_H/r$ of equation B13. $mv^2/r=GMm/r^2$ is always true (eg.,globulars orbiting out of plane) but so is $g_{00}=\kappa_{00}$ in the plane of a flattened galaxy (rotating central black hole planar effect partIII). That $g_{00}=\kappa_{00}$ in the halo of the Milky Way galaxy is the fundamental equation of metric quantization. So again $mv^2/r=GMm/r^2$ so $GM/r=v^2$ COM in the galaxy halo(circular orbits) ($1/(1-2\epsilon)$ term from κ_{00} in B13) so

Pure state $\Delta\epsilon$ (ϵ excited $1S_{1/2}$ state of ground state $\Delta\epsilon$, so not same state as $\Delta\epsilon$)

$Rel\kappa_{00}=\cos\mu$ from B13 κ_{00}

$$\text{Case 1 } 1-2GM/(c^2r)=1-2(v/c)^2=1-(2\Delta\epsilon/(1-2\epsilon))^2/2 \quad (B12)$$

So $1-2(v/c)^2=1-(2\Delta\epsilon/(1-2\epsilon))^2/2$ so $=(2\Delta\epsilon/(1-2\epsilon))c/2=2X.0002826/(1-(.05946)^2)(3X10^8)/2=99km/sec \approx 100km/sec$ (Mixed $\Delta\epsilon,\epsilon$, states classically here are grand canonical ensembles with nonzero chemical potential.). For ringed (not hub) galaxies the radial value becomes $100/2=50km/sec$.

Also $v=(2\Delta\epsilon/(1-2\epsilon))c/2$ so $v/c=\text{constant}$.

Mixed state $\epsilon\Delta\epsilon$ (Again $GM/r=v^2$ so $2GM/(c^2r)=2(v/c)^2$.)

$$\text{Case 2 } g_{00}=1-2GM/(c^2r)=Rel\kappa_{00}=\cos[2\Delta\epsilon+\epsilon]=1-[\Delta\epsilon+\epsilon]^2/2=1-[(2\Delta\epsilon+\epsilon)^2/(\Delta\epsilon+\epsilon)]^2/2=1-[(2\Delta\epsilon^2+\epsilon^2+2\epsilon\Delta\epsilon)/(2\Delta\epsilon+\epsilon)]^2$$

The $2\Delta\epsilon^2$ is just the above first case (Case 1) so just take the mixed state cross term $[\epsilon 2\Delta\epsilon/(\epsilon+2\Delta\epsilon)]=c[2\Delta\epsilon/(1+\Delta\epsilon/\epsilon)]/2=c[2\Delta\epsilon+\Delta\epsilon^2/\epsilon+\dots 2\Delta\epsilon^{N+1}/\epsilon^{N+1}]/2=\Sigma v_N$. Note each term in this expansion is itself a (mixed state) operator. So there can't be a single v in the large gradient 2nd case so in the equation just above we can take $v_N=[2\Delta\epsilon^{N+1}/(2\epsilon^N)]c$. (B13)

From eq. B23 for example $v=m100^N km/sec$. $m=2,N=1$ here (Local arm). In part III we list hundreds of examples of B13: (sun1,2km/sec, galaxy halos m100km/sec). The linear mixed state subdivision by this ubiquitous ~ 100 scale change factor in r_{bb} (due to above object B zitterbewegung spherical Bessel function resonance boundary conditions resulting in nodes) created the voids. Same process for $N-1$ (so 100X smaller) antinodes get galaxies, 100Xsmaller: globular clusters, 100Xsmaller solar systems, etc., So these smaller objects were also created by mixed state metric quantization (eq.B13) resonance oscillation inside initial radius r_{bb} .

We include the effects of that object B drop in inertial frame dragging on the inertial term m in the Gamow factor and so lower Z nuclear synthesis at earlier epochs ($t>18by$)BCE. (see partIII)

Appendix C. Object C with spinor ansatz for eq.12(gives ordinary field theory SM)

For the $N=1$ huge observer $\delta z \gg \delta z \delta z$ from eq.3. Thus the required $N=-1,N=0$ tiny observable ($\delta z' \ll \delta z$) is a perturbation of the eq.7 $\delta z \approx dr \approx dt$ at 45° $(dr-\delta z')+(dt+\delta z') \equiv dr'+dt'=ds$ (12)

But for the high energy big $\delta \delta z$ (extreme "axis" perturbations) δz is small. So finding big $\delta \delta z$

'observables' requires we artificially stay on the circle (appendix C) implying this additional $\delta z'$ eq7 perturbation. These large rotations can then be done as spinor rotations \rightarrow Pauli matrices \rightarrow isomorphic to quaternions

So we finally include the orthogonal axis' to orthogonal axis extreme ds rotations in eq. 12. So on the circle. Recall from sect.1 eq.3 that $\delta C = \delta(\delta z + \delta z \delta z) = \delta \delta z(1) + \delta \delta z(\delta z) + (\delta z)\delta \delta z = \delta C = 0$ so C is split between $\delta \delta z$ noise and $\delta z \delta z$ and classical ds^2 proper time. Note for $N=1$ $|\delta z| \gg 1$ and $C_M \gg 1$. So eq.5 holds then. But for high energies (like those provided by an accelerator) as γ is boosted observer $\delta z/\gamma$, C/γ gets smaller than the huge $N=1$ scale (so higher energy, smaller wavelength beam probes) $\delta \delta z(1)/ds$ noise angle gets relatively larger (relative to $\delta(\delta z \delta z)/ds$, sect.1) until finally the next smaller (and next smaller one after that at $N=-1$) is the $N=0$ fractal scale becomes relatively large.

Large rotation angle $\delta \delta z/ds$ can then be those large axis' ds extreme $\pm 45^\circ$ min ds and so two possible 45° rotations so through a total of two quadrants for $\pm \delta z'$ in eq.12. (a single δz just gives e,v back) One such rotation around a axis (SM) and the other around a diagonal (SC).

These rotations are

I \rightarrow II, II \rightarrow III, III \rightarrow IV, IV \rightarrow I required extremum to eq.16 extremum rotations in eq.7-9 plane Give SM Bosons at high interaction COM energies (where $\delta \delta z$ gets big). $N_{ob} = 0$

Note in fig.3 dr, dt is also a rotation. and so has an eq.11 rotation operator observable θ . Thus from equation 11 for (θ) angle rotations $\theta \delta z = (dr/ds)\delta z = -i\partial(\delta z)/\partial r$ for the first 45° rotation. So we got through one Newpde derivative for each 45° rotation. For the next 45° rotation in fig.4 it is then a second derivative $\theta \theta \delta z' = e^{i\theta p} e^{i\theta'} \delta z = e^{i(\theta p + \theta')} \delta z = (dr/ds)((dr/ds)dr') = -i\partial(-i\partial(dr'))/\partial r = -\partial^2(dr')/\partial r^2$ large angle rotation in figure 3. In contrast for $z=1$, $\delta z'$ small so $45^\circ-45^\circ$ small angle rotation in figure 3 (so then $N=-1$). Do the same with the time t and get for $z=0$ rotation of $45^\circ+45^\circ$ (fig.4) then $\theta \theta \delta z' = (d^2/dr^2)z' + (d^2/dt^2)\delta z'$ (C1)

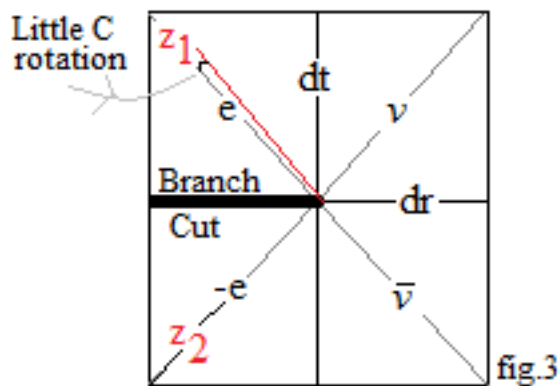


fig.3. for $45^\circ-45^\circ$ So two body (e,v) singlet $\Delta S = 1/2 - 1/2 = 0$ component so pairing interaction (sect.4.5). Also ortho $\Delta S = 1/2 + 1/2 = 1$ making 2 body (at $r=r_H$) $S=1$ Bosons and so a field theory. Note we also get these Laplacians characteristic of the Boson field equations by those $45^\circ+45^\circ$ rotations so eq.16 implies Bosons accompany our leptons (given the $\delta z'$), **so these leptons exhibit "force"**.

Newpde $r=r_H$, $z=0$, $45^\circ+45^\circ$ rotation of composites e,v implied by Equation 12

So $z=0$ allows a large C z rotation application from the 4 different axis' max extremum (of eq.16) branch cuts gives the 4 results: Z,+W, photon bosons of the Standard Model. So we have derived the Standard Model of particle physics in this very elegant way (from the four axis'). You are physically at $r=r_H$ if you rotate through the electron quadrants (I, IV).of eq.7-9. So we

have large C_M dichotomic 90° rotation to the next Reimann surface of eq.12, eq.A1 $(dr^2+dt^2)z''$ from some initial extremum angle(s) θ . Eq.16 solutions imply complex 2D plane Stern Gerlach dichotomic rotations using eq.A1 thereby using Pauli matrices σ_i algebra, which maps one-to-one to the quaternion algebra. Using eq.12 we start at some initial angle θ and rotate by 90° the noise rotations are: $C=\delta z'' = [e_L, v_L]^T \equiv \delta z'(\uparrow) + \delta z'(\downarrow) \equiv \psi(\uparrow) + \psi(\downarrow)$ has a eq.12 infinitesimal unitary generator $\delta z'' \equiv U = 1 - (i/2)\epsilon n^* \sigma$, $n = \theta/\epsilon$ in $ds^2 = U^t U$. But in the limit $n \rightarrow \infty$ we find, using elementary calculus, the result $\exp(-(i/2)\theta^* \sigma) = \delta z''$. We can use any axis as a branch cut since all 4 are eq.20 large extremum so for the 2nd rotation we move the branch cut 90° and measure the angle off the next diagonal since Pauli matrix dichotomic rotations are actually axis rotations, leaving our e and v directions the same. In any case $(dr+dt)z''$ in eq.16 can then be replaced by eq.A1 $(dr^2+dt^2+..) \delta z'' = (dr^2+dt^2+..) e^{\text{quaternionA}}$ Bosons because of eq.C1. C2 Then use eq. 12 and quaternions to rotate $\delta z''$ since the quaternion formulation is isomorphic to the Pauli matrices. $dr' = \delta z_r = \kappa_r dr$ for **Quaternion A** $\kappa_{ii} = e^{iA_i}$.

Possibly large $\delta\delta z$ in eq.3 $\delta(\delta z + \delta z \delta z) = 0$ so large rotations in eq12 i.e., high energy, tiny $\sqrt{\kappa_{00}}$ since τ normalized to 1 allows formalism for object C

C1 for the eq.12: large $\theta = 45^\circ + 45^\circ$ rotation (for $N=0$ so large $\delta z' = \theta r_H$). Instead of the equation 13,15 formulation of κ_{ij} for small $\delta z'$ ($z=1$) and large $\theta = 45^\circ + 45^\circ$ we use A_r in dr direction with $dr^2 = x^2 + dy^2 + dz^2$. So we can again use 2D (dr, dt) $E = 1/\sqrt{\kappa_{00}} = 1/\sqrt{e^{iA_i}} = e^{-iA/2}$. The 1 is mass energy and the first real component after that in the Taylor expansion is field energy A^2 . For 2 particles together the other particle ϵ negative means r_H is also negative. Since it is $e_1^* e_2 = r_H$. So $1/\kappa_r = 1 + (-\epsilon + r_H/r)$ is \pm and $1 - (-\epsilon + r_H/r)$ 0 charge. (C0)

For baryons with a 3 particle r_H/r may change sign without third particle ϵ changing sign so that at $r=r_H$. Can normalize out the background ϵ in the denominator of $E = (\tau + \epsilon)/\sqrt{(1 + \epsilon + \Delta\epsilon - r_H/r)}$ for Can normalize out the background ϵ in the denominator of $E = (\tau + \epsilon)/\sqrt{(1 + \epsilon + \Delta\epsilon - r_H/r)}$ for small conserved (constant) energies $1/\sqrt{(1 + \epsilon)}$ and (so $E = (1/\sqrt{(1 + x)}) = 1 - x/2 +$) large r (so large λ so not on r_H) implies the normalization is:

$E = (\epsilon + \tau)/\sqrt{((1 - \epsilon/2 - \epsilon/2)/(1 \pm \epsilon/2))}$, $J=0$ para e, v eq.9.23 π^\pm, π^0 . For large $1/\sqrt{\Delta\epsilon}$ energies given small $r=r_H$, Here $1 + \epsilon$ is locally constant so can be normalized out as in

$$E = (\epsilon + \tau)/\sqrt{(1 - (\Delta\epsilon/(1 \pm \epsilon)) - r_H/r)}, \text{ for charged if -, ortho } e, v \text{ } J=1, W^\pm, Z_0 \text{ (11d)}$$

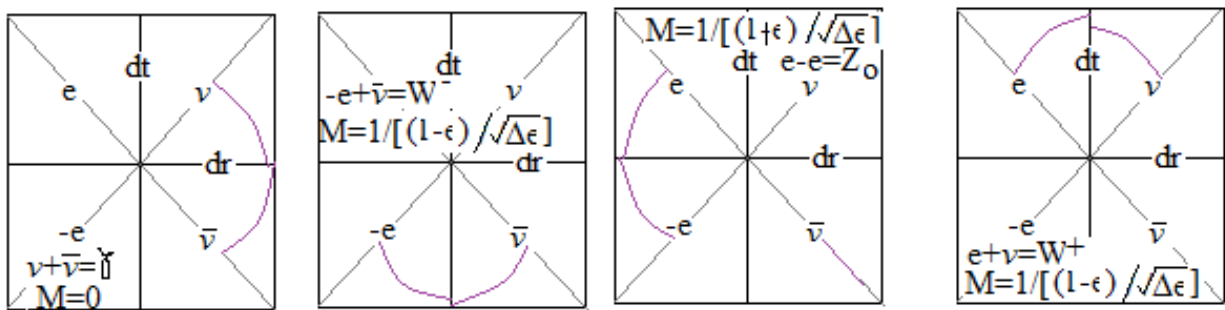


fig4

Fig.4 applies to eq.9 $45^\circ + 45^\circ = 90^\circ$ case: **Bosons**.

C2 These quadrants were defined in eq.7-9 and used in eq.12. The Appendix C4 derivation applies to the far right side figure. Recall from eq.16 $z=0$ result $C_M = 45^\circ + 45^\circ = 90^\circ$, gets Bosons. $45^\circ - 45^\circ =$ leptons. The v in quadrants II (eq.5) and III (eq.9). e in quadrants I (eq.7) and IV (eq.7).

Locally normalize out $1+\varepsilon$ (appendix D). For the **composite e,v** on those required large $z=0$ eq.9 rotations for $C\rightarrow 0$, and for stability $r=r_H$ (eg., for $2P_{1/2}$, $I\rightarrow II$, $III\rightarrow IV$, $IV\rightarrow I$) unless $r_H=0$ ($II\rightarrow III$)

Example:

C4 Quadrants IV→I rotation eq.C2 $(dr^2+dt^2+..)_e^{\text{quaternion } A}$ =rotated through C_M in eq.16.

example C_M in eq.C1 is a 90° CCW rotation from 45° through v and anti v

A is the 4 potential. From eq.17 we find after taking logs of both sides that $A_o=1/A_r$ (A2)

Pretending we have a only two i,j quaternions but still use the quaternion rules we first do the r

derivative: From eq. C1 $dr^2\delta z=(\partial^2/\partial r^2)(\exp(iA_r+jA_o))=(\partial/\partial r[(i\partial A_r/\partial r+\partial A_o/\partial r)(\exp(iA_r+jA_o))]$

$$=\partial/\partial r[(\partial/\partial r)iA_r+(\partial/\partial r)jA_o](\exp(iA_r+jA_o))+[i\partial A_r/\partial r+j\partial A_o/\partial r]\partial/\partial r(iA_r+jA_o)(\exp(iA_r+jA_o))+$$

$$(i\partial^2 A_r/\partial r^2+j\partial^2 A_o/\partial r^2)(\exp(iA_r+jA_o))+[i\partial A_r/\partial r+j\partial A_o/\partial r][i\partial A_r/\partial r+j\partial/\partial r(A_o)] \exp(iA_r+jA_o) \quad (A3)$$

Then do the time derivative second derivative $\partial^2/\partial t^2(\exp(iA_r+jA_o))=(\partial/\partial t[(i\partial A_r/\partial t+\partial A_o/\partial t)$

$(\exp(iA_r+jA_o))]=\partial/\partial t[(\partial/\partial t)iA_r+(\partial/\partial t)jA_o](\exp(iA_r+jA_o))+$

$[i\partial A_r/\partial r+j\partial A_o/\partial t]\partial/\partial r(iA_r+jA_o)(\exp(iA_r+jA_o))+(i\partial^2 A_r/\partial t^2+j\partial^2 A_o/\partial t^2)(\exp(iA_r+jA_o))$

$$+[i\partial A_r/\partial t+j\partial A_o/\partial t][i\partial A_r/\partial t+j\partial/\partial t(A_o)]\exp(iA_r+jA_o) \quad (C4)$$

Adding eq. C2 to eq. C4 to obtain the total D'Alambertian $C3+C4=$

$[i\partial^2 A_r/\partial r^2+i\partial^2 A_r/\partial t^2]+[j\partial^2 A_o/\partial r^2+j\partial^2 A_o/\partial t^2]+ii(\partial A_r/\partial r)^2+ij(\partial A_r/\partial r)(\partial A_o/\partial r)$

$+ji(\partial A_o/\partial r)(\partial A_r/\partial r)+jj(\partial A_o/\partial r)^2++ii(\partial A_r/\partial t)^2+ij(\partial A_r/\partial t)(\partial A_o/\partial t)+ji(\partial A_o/\partial t)(\partial A_r/\partial t)+jj(\partial A_o/\partial t)^2$.

Since $ii=-1, jj=-1, ij=-ji$ the middle terms cancel leaving $[i\partial^2 A_r/\partial r^2+i\partial^2 A_r/\partial t^2]+$

$[j\partial^2 A_o/\partial r^2+j\partial^2 A_o/\partial t^2]+ii(\partial A_r/\partial r)^2+jj(\partial A_o/\partial r)^2+ii(\partial A_r/\partial t)^2+jj(\partial A_o/\partial t)^2$

Plugging in C2 and C4 gives us cross terms $jj(\partial A_o/\partial r)^2+ii(\partial A_r/\partial t)^2=jj(\partial(-A_r)/\partial r)^2+ii(\partial A_r/\partial t)^2$

$$=0. \text{ So } jj(\partial A_r/\partial r)^2=-jj(\partial A_o/\partial t)^2 \text{ or taking the square root: } \partial A_r/\partial r+\partial A_o/\partial t=0 \quad (C5)$$

$$i[\partial^2 A_r/\partial r^2+i\partial^2 A_r/\partial t^2]=0, \quad j[\partial^2 A_o/\partial r^2+i\partial^2 A_o/\partial t^2]=0 \text{ or } \partial^2 A_\mu/\partial r^2+\partial^2 A_\mu/\partial t^2+..=1 \quad (C6)$$

A4 and A5 are **Maxwell's equations** (Lorentz gauge formulation) in free space, if $\mu=1,2,3,4$.

$$\square^2 A_\mu=1, \quad \square \bullet A_\mu=0 \quad (C7)$$

This looks like the Lorentz gauge formalism but it is actually a fundamental field equation (not interchangeable with some other as in gauge theories) hence it is *no gauge at all* and we have also avoided the Maxwell overdeterminism problem (8eq, 6 unknowns E_i, B_i). Must use Newpde 4D orthogonalization here. Amplitudes of physical processes in QED in the noncovariant Coulomb gauge coincide with those in the covariant Lorenz gauge. The Aharonov–Bohm effect depends on a line integral of A around a closed loop, and this integral is not changed by $A\rightarrow A+\nabla\psi$ which doesn't change $B=\nabla\times A$ either. So formulation in the Lorentz gauge mathematics works (but again 6.7 is no longer a gauge).

C5 Other $45^\circ+45^\circ$ Rotations (Besides above quadrants IV→I)

Proca eq

In the 1st to 2nd, 3rd to 4th quadrants the A_u is already there as a single v in the rotation the mass is in both quadrants and in the end we multiply by the A_u so get the $m^2 A_u^2$ term in the Proca eq.for the W^+, W^- . The mass still gets squared for the 2nd to 3rd quadrant rotation $Z_o..$

For the **composite e,v** on those required large $z=0$ eq.16 rotations for $C\approx 0$, and for stability $r=r_H$ for $2P_{1/2}$ ($I\rightarrow II$, $III\rightarrow IV$, $II\rightarrow III$) unless $r_H=0$ ($IV\rightarrow I$) are:

Ist→IInd quadrant rotation is the W^+ at $r=r_H$. Do similar math to C2-C7 math and get instead a Proca equation The limit $\varepsilon \rightarrow 1 = \tau$ (D13) in ξ_1 at $r=r_H$. since Hund's rule implies $\mu = \varepsilon = 1S_{1/2} \leq 2S_{1/2} = \tau = 1$. So the ε is negative in $\Delta\varepsilon/(1-\varepsilon)$ as in case 1 charged as in appendix A1 case 2.

$E = 1/\sqrt{(\kappa_{00})} - 1 = [1/\sqrt{(1-\Delta\varepsilon/(1-\varepsilon)-r_H/r)}] - 1 = [1/\sqrt{(\Delta\varepsilon/(1-\varepsilon))}] - 1$. $E_t = E + E = 2/\sqrt{(\Delta\varepsilon/(1-\varepsilon))} = W^+$ mass.
 $E_t = E - E$ gives E&M that also interacts weakly with weak force.

IIIrd →IV quadrant rotation is the W^- . Do the math and get a Proca equation again.

$E = 1/\sqrt{(\kappa_{00})} - 1 = [1/\sqrt{(1-\Delta\varepsilon/(1-\varepsilon)-r_H/r)}] - 1 = [1/\sqrt{(\Delta\varepsilon/(1-\varepsilon))}] - 1$. $E_t = E + E = 2/\sqrt{(\Delta\varepsilon/(1-\varepsilon))} = W^-$ mass.
 $E_t = E - E$ gives E&M that also interacts weakly with weak force.

II → III quadrant rotation is the Z_0 . Do the math and get a Proca equation. C_M charge cancelation. B14 gives $1/(1+\varepsilon)$ gives 0 charge since $\varepsilon \rightarrow 1$ to case 1 in appendix C2.

$E = 1/\sqrt{(\kappa_{00})} - 1 = [1/\sqrt{(1-\Delta\varepsilon/(1+\varepsilon)-r_H/r)}] - 1 = [1/\sqrt{(\Delta\varepsilon/(1+\varepsilon))}] - 1$. $E_t = E + E = 2/\sqrt{(\Delta\varepsilon/(1+\varepsilon))} - 1 = Z_0$ mass.
 $E_t = E - E$ gives E&M that also interacts weakly with weak force. Seen in small left handed polarization rotation of light.

IV→I quadrant rotation through those 2 neutrinos gives 2 objects. $r_H = 0$

$E = 1/\sqrt{\kappa_{00}} - 1 = [1/\sqrt{(1-\Delta\varepsilon/(1+\varepsilon))}] - 1 = \Delta\varepsilon/(1+\varepsilon)$. Because of the +- square root $E = E + -E$ so E rest mass is 0 or $\Delta\varepsilon = (2\Delta\varepsilon)/2$ reduced mass.

$E_t = E + E = 2E = 2\Delta\varepsilon$ is the pairing interaction of SC. The $E_t = E - E = 0$ is the 0 rest mass photon Boson. Do the math (eq.C2-C7) and get Maxwell's equations. Note there was no charge C_M on the two ν s. Note we get SM particles out of composite e, ν using required eq.9 rotations for

C6 Object B Effect On Inertial Frame Dragging (from appendix B)

The fractal implications are that we are inside a cosmological positron inside a proton $2P_{3/2}$ at $r=r_H$ state. The cosmological object (electron) we are inside of is a positron and call it object A which orbits electron object B with a given distant 3rd object C. Object B is responsible for the mass of the electron since it's frame dragging creates that Kerr metric $(a/r)^2 = m_e c^2$ (B9) result used in eq.D9. So Newpde ground state $m_e c^2 \equiv \langle H_e \rangle$ is the fundamental Hamiltonian eigenvalue defining idea for composite e, ν , $r=r_H$ implying Fermi 4 point $E = \int \psi^\dagger H \psi dV = \int \psi^\dagger \psi H dV = \int \psi^\dagger \psi G$

Recall for composite e, ν all interactions occur inside r_H $(4\pi/3)\lambda^3 = V_{rH}$. $\frac{1}{v^{1/2}} = \psi_e = \psi_3 \frac{1}{v^{1/2}} =$

$$\psi_\nu = \psi_4 \text{ so } 4\text{pt} \iiint_0^{r_H} \psi_1 \psi_2 \psi_3 \psi_4 dV = 2G \iiint_0^{r_H} \psi_1 \psi_2 \frac{1}{v^{1/2}} \frac{1}{v^{1/2}} V$$

$$\equiv \iiint_0^{r_H} \psi_1 \psi_2 G \equiv \iiint_0^{r_H} \psi_1 \psi_2 (2m_e c^2) dV_{rH} = \iiint_0^{r_H} \psi_1 (2m_e c^2) \psi_2 dV_{rH} \quad (A8)$$

Object C adds its own spin (eg., as in 2nd derivative eq.A1) to the electron spin (1,IV quadrants) and the W associated with the $2P_{3/2}$ state at $r=r_H$ thereby adds a derivative in a neutrino quadrant (fig.4) thereby including neutrinos in the Fermi 4pt. So 2nd derivative

$$\Sigma((\gamma^\mu \sqrt{\kappa_{\mu\mu}} dx_\mu) - i\kappa)(\gamma^\nu \sqrt{\kappa_{\nu\nu}} dx_\nu + i\kappa)\chi = \Sigma((\gamma^\mu \sqrt{\kappa_{\mu\mu}} dx_\mu) - i\kappa)\psi \text{ so } \frac{1}{2}(1 \pm \gamma^5)\psi = \chi. \quad (A9)$$

In that regard the expectation value of γ^5 is speed and varies with $e^{i3\phi/2}$ in the trifolium. The spin $^{1/2}$ decay proton $S_{1/2} \propto e^{i\phi/2} \equiv \psi_1$, the original ortho $2P_{1/2}$ particle is chiral $\chi = \psi_2 \equiv \frac{1}{2}(1 - \gamma^5)\psi = \frac{1}{2}(1 - \gamma^5 e^{i3\phi/2})\psi$. Initial $2P_{1/2}$ electron ψ is constant. Start with initial ortho state χ . These γ^5 terms then

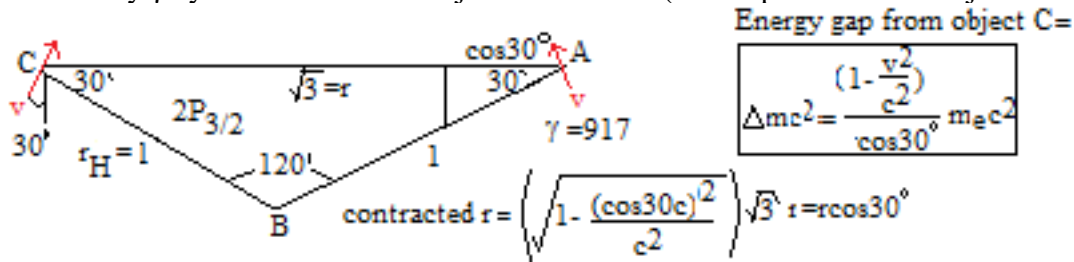
modify equation A8 to read $= \iiint_0^{r_H} \psi_1 \psi_2 (2m_e c^2) dV_{rH} = \iint \psi_{S1/2}^* (2m_e c^2 V_{rH}) \chi dV_\phi =$

$$K \int \langle e^{i\frac{\phi}{2}} [\Delta\varepsilon V_{rH}] \left(1 - \gamma^5 e^{i\frac{3\phi}{2}}\right) \psi \rangle d\phi = KG_F \int \langle e^{i\phi/2} - \gamma^5 e^{i(4/2)\phi} \rangle d\phi = KG_F \left(\frac{2e^{i\phi}}{i} \Big|_0^{2\pi+C} - \right.$$

$\frac{2\gamma^5 e^{i4\phi}}{i^4} |0^{2\pi+c}) = k1(1/4+i\gamma^5) = k(.225+i\gamma^5 0.974) = k(\cos 13^\circ + i\gamma^5 \sin 13^\circ)$ **deriving the 13° Cabbibo angle.** With previously mentioned CP result(direct evidence of fractal universe) get CKM matrix

C7 Object C Effect on Inertial Frame Dragging and G_F found by using eq.C8 again (N=1 ambient cosmological metric)

Review of $2P_{3/2}$ Next higher fractal scale ($X10^{40}$), cosmological scale. Recall from B9 $m_e c^2 = \Delta \epsilon$ is the energy gap for object B vibrational stable iegenstates of composite $3e$ (vibrational perturbation r is the only variable in Frobenius solution, partII Ch.8,9,10) proton. Observer in objectA. $\Delta m_e c^2$ gap=object C scissors eigenstates. is what we see at object A but $\Delta m_e c^2$ gets boosted by γ by rotation into the object B direction.(to compare with the object B $m_e c^2$ gap).



From fig 7 $r^2 = 1^2 + 1^2 + 2(1)(1)\cos 120^\circ = 3$, so $r = \sqrt{3}$. Recall for the positron motion $\gamma = \frac{1}{\sqrt{1 - \frac{v^2}{c^2}}} = 917$.

So start with the distances we observe which are the Fitzgerald contracted $AC =$

$r_{CA} = 1 \sqrt{1 - \frac{\cos^2 30^\circ c^2}{c^2}} \sqrt{3} = .866 = \cos 30^\circ = CA$ and Fitzgerald contracted $AB = r_{BA} = x/\gamma = 1/\gamma$ so for Fitzgerald contracted $x=1$ for AB (fig7). We can start at $t=0$ with the usual Lorentz transformation for the time component.

$$t' = \gamma(ct - \beta x) = kmc^2.$$

since time components are Lorentz contracted proportionally also to mc^2 , both with the γ multiplication.

In the object A frame of reference we see $\Delta m_e c^2$ which is the average of left and right object C motion effect. We go into the AB frame of reference to compare the object B $m_e c^2$ with this $\Delta m_e c^2$. Going into the AB frame automatically boosts $\Delta m_e c^2$ to $\gamma \Delta m_e c^2$. So start from a already Fitzgerald contracted x/γ . Next do the time contraction γ to that frame:

$$t'' = k\gamma \Delta m_e c^2 = \gamma \beta r_{AB} = \gamma \beta \left(\frac{x}{\gamma}\right) = \frac{1}{\sqrt{1 - \frac{v^2}{c^2}}} \beta \left(\sqrt{1 - \frac{v^2}{c^2}} \sqrt{1}\right) = \beta$$

with k defining the projection of tiny $\Delta m_e c^2$ "time" CA onto BA = $\cos \theta$ = projection of BA onto CA. But $m_e c^2$ is the result of object B of both of the motion and inertial frame dragging reduction (D9) so its γ is large. To make a comparison of ΔE to AB mass $m_e c^2$ CA is rotated and translated to the high speed AB diection and distance with its large γ so thereby *object C becomes mathematically object B* with the same k because of these projection properties of: CA onto BA. So we define projection k from projection of $m_e c^2$: So again

$$t' = \gamma(ct - \beta x) = kmc^2 = t'' = km_e c^2 = \gamma \beta r_{CA} = \left(\frac{1}{\sqrt{1 - \frac{v^2}{c^2}}}\right) \beta \left(\sqrt{1 - \frac{\cos^2 30^\circ c^2}{c^2}} \sqrt{3}\right) = \gamma \beta \cos 30^\circ$$

Take the ratio of $\frac{k\gamma\Delta m_e c^2}{km_e c^2}$ to eliminate k: thus

$$\frac{k\gamma\Delta m_e c^2}{km_e c^2} = \frac{\gamma\beta\left(\frac{x}{\gamma}\right)}{\gamma\beta r_{CA}} = \frac{1\beta 1}{\gamma\beta\cos 30^\circ} = \frac{1}{\gamma\cos 30^\circ} \quad \text{so}$$

$$\Delta m_e c^2 = \frac{\beta m_e c^2}{\beta\cos 30^\circ \gamma^2} = \frac{\left(1-\frac{v^2}{c^2}\right)m_e c^2}{\cos 30^\circ} \quad (\text{A10})$$

allowing us to finally compare the energy gap caused by object C ($\Delta m_e c^2$) to the energy gap caused by object B ($m_e c^2$, C8). So to summarize $\Delta E = (m_e c^2 / ((\cos 30^\circ)^2 \gamma^2)) = m_e c^2 / 728000$. So the energy gap caused by object C is $\Delta E = (m_e c^2 / ((\cos 30^\circ)^2 \gamma^2)) = m_e c^2 / 728000$. The weak interaction thereby provides the ΔE perturbation ($\int \psi^* \Delta E \psi dV$) inside of r_H creating those Frobenius series (partII) $r \neq 0$ states, for example in the unstable equilibrium $2P_{1/2}$ electrons m_e . so in the context of those e, v rotations giving W and Z_0 . The G can be written for E&M decay as $(2m_e c^2) X V_{r_H} = 2m_e c^2 [(4/3)\pi r_H^3]$. But because this added object C rotational motion is eq.A9 Fermi 4 point it is entirely different than a mere 'weak' E&M. So for weak decay from equation A8 it is $G_F = (2m_e c^2 / 728,000) V_{r_H} = G_F = 1.4 \times 10^{-62} \text{ J}\cdot\text{m}^3 = .9 \times 10^{-4} \text{ MeV}\cdot\text{F}^3$ **the strength of the Fermi 4pt weak interaction constant** which is the coupling constant for the Fermi 4 point weak interaction integral. Note $2m_e c^2 / 729,000 = 1.19 \times 10^{-19} \text{ J}$. So $\Delta E = 1.19 \times 10^{-19} / 1.6 \times 10^{-19} = .7 \text{ eV}$ which \pm that r perturbation (instability) states in the Frobenius solution (partII) and so weak decay. interaction integral. Note $2m_e c^2 / 729,000 = 1.19 \times 10^{-19} \text{ J}$. So $\Delta E = 1.19 \times 10^{-19} / 1.6 \times 10^{-19} = .7 \text{ eV}$ which is our ΔE gap for the weak interaction (from operator H) inside the Fermi 4pt. integral for G_F .

The perturbation r in the Frobenius solution is caused by this ΔH in ($\int \psi^* \Delta H \psi dV$) with available phase space $\psi^* = \psi_p \psi_e \psi_\nu$ for $\psi = \psi_N$ decay where ψ_e and ψ_ν are from the factorization. The neutrino adds a $e^2(0)$ to the set of $e^2 10^{40N}$ electron solutions to Newpde r_H with electron charge $\pm e$ and intrinsic angular momentum conservation laws $S = 1/2$ holding for both e and ν .

The neutrino mass increases with nonisotropic homogenous space-time (sect.3.1 and our direction of motion here) whereas that Kerr metric $(a/r)^2$ term (B9) in general is isotropic and homogenous and so only effects the electron mass.

C8 NONhomogeneous and NONisotropic Space-Time

Recall 2D $N=1$ and that 2D $N=0$ (perturbation) orientations are not correlatable so we have $2D+2D=4D$ degrees of freedom. But this is all still embedded in the same complex (2D) plane. So this theory is still geometrically complex 2D Z then. Recall the $\kappa_{\mu\nu} = g_{\mu\nu}$ metrics (and so R_{ij} and R) were generated in section 1.

In that regard for 2D for a homogenous and isotropic g_{ij} we have identically $R_{\mu\mu} - 1/2 g_{\mu\mu} R = 0$ (3.1.1) \equiv source $= G_{00}$ since in 2D $R_{\mu\mu} = 1/2 g_{\mu\mu} R$ identically (Weinberg, pp.394) with $\mu=0, 1, \dots$. Note the 0 ($=E_{\text{total}}$ the energy density source) and we have thereby proven the existence of a net zero energy density vacuum. Thus our 2D **theory implies the vacuum is really a vacuum!** It is then the result of the fractal and 2D nature of space time!

A ultrarelativistic electron is essentially a transverse wave 2D object (eg., the $2P_{1/2}$ electron in the neutron). In an isotropic homogenous space time $G_{00}=0$. Also from sect.2 eqs. 7 and 8 (9) occupy the same complex 2D plane. So eqs. 7+8 is $G_{00} = E_e + \sigma \cdot p_r = 0$ so $E_e = -\sigma \cdot p_r$. So given the negative sign in the above relation the **neutrino chirality is left handed**. But if the space time is not isotropic and homogenous then G_{00} is not zero and the **neutrino gains mass**.

C9 Derivation of the Standard Electroweak Model from Newpde but with No Free parameters

Since we have now derived M_W , M_Z and their associated Proca equations, and Dirac equations for m_τ, m_μ, m_e etc., and G, G_F, ke^2 Maxwell's equations, etc. we can now write down the usual Lagrangian densities that implies these results. In the formulation $M_Z = M_W / \cos\theta_W$ you can find the Weinberg angle θ_W , $g \sin\theta_W = e$, $g' \cos\theta_W = e$; solve for g and g' , etc., We will have thereby derived the standard model from first principles (i.e., postulate 0). **It no longer contains free parameters.**

Note $C_M =$ Figenbaum pt really is the U(1) charge and equation 16 rotation is on the complex plane so it really implies SU(2) (C1) with the sect.1.2 2D eqs. $7+8 = G_{00} = E_e + \sigma \cdot p_r = 0$ gets the left handedness. Recall the genius of the SM is getting all those properties (of χ, Z_0, W^+, W^-) from SU(2)XU(1)_L so we really have completely derived the electroweak standard model from eq.16 which comes out of the Newpde given we even found the magnitude of its input parameters (eg., G_F (appendix C7), Cabbibo angle C6).

Appendix D Counting actual quanta numbers N (instead of just n energy level 2nd quantization states |n>)

D1 Recal from equation 11 $\left[\left(\frac{dr+dt}{ds} \right) \right] \delta z = \frac{ds}{ds} \delta z = (1) \delta z$ In that "implied iteration of the first application $\left[\left(\frac{dr+dt}{ds} \right) + \left(\frac{dr+dt}{ds} \right) \right] \delta z = 2 \frac{ds}{ds} \delta z = 2(1) \delta z$ For all the rotations in fig.4 (except the eq.11 IVth to Ist quadrants: in eq.6.1 each quadrant rotation provides one derivative for each ν) $\left[\left(\frac{dr+dt}{ds} \right) + \left(\frac{dr+dt}{ds} \right) \right] \delta z = 2 \frac{ds}{ds} \delta z = 2(1) \delta z$. Equation 11 (sect.1) then counts units N of each 2 half integer $S = 1/2$ angular momentums = 1 unit oelectrons (spin1 for W and Z) off the light cone. For the rotation in the eq.11 IVth to Ist quadrants (each quadrant rotation provides one derivative for each ν) at 45° $dr = dt$ (on the light cone in fig.4) so for Hamiltonian H: $2H \delta z = 2(dt/ds) \delta z = 2(1/2) \delta z = (1) \hbar \omega \delta z = \hbar c k \delta z$ on the diagonal so that $E = p_r = \hbar \omega$ for the two ν energy components, universally. Thus we can state the most beautiful result in physics that $E = N h f$ for the energy of light with N equal N monochromatic photons. Thus this eq.11c counting N does not require the (well known) quantization of the E&M field with SHM (sect.6.10 below). Which seemed to me at least a adhoc process on the face of it since the Maxwell equations have nothing to do with SHM.

Given this comes from equation 11, these numbers are thereby "observables". We have come full circle, getting eq.11 'observables' and using equation 11 to define our inputs into the '1' in $1 = 1+0, 1 = 1X1, 0 = 0X0$ as an observable (Newpde electrons), thereby starting our entire derivation all over again..

All defined numbers, and resulting symbols and rules, that are larger than 1 ($N > 1$) we define as "applications" given our ultimate Occam's Razor attribute of the postulate of 0. Note applications can be arbitrarily complicated.

D2 Postulate 0 also implies the underlying 1,0 relationship and $n > 1$ "applications"

Review Postulate 0: No need for a complicated definition because there is nothing there to define! The null set would be simpler and ultimate occam's razor but you don't postulate it, since it is subset of every set anyway.

So by the process of elimination we arrive at the ultimate Occam's razor *postulate* real#0, the very next level up.

But we need to define the algebra first and use it to write the postulate 0. So define

1) numbers $1 \equiv 1+0$ and $0 \equiv 0 \times 0, 1 \equiv 1 \times 1$ as symbol $z=zz$: the simplest algebraic definition of 0. So
 2) Postulate real number 0 if $z'=0$ and $z'=1$ plugged into $z'=z'z'+C$ (eq.1) results in some $C=0$ constant (ie $\delta C=0$).

This is our *entire* (Ultimate Occam's Razor postulate(0)) theory

Application: (i.e., plug $z=1,0$ into eq.1 as required by above theory.)

Plug in $z=0=z'$ in eq.1. The equality sign in eq.1 demands we substitute z' on left (eq.1) into right $z'z'$ repeatedly and get iteration $z_{N+1}=z_N z_N + C$. If $C=1$ and $z_N=1$ then $z_{N+1}=2$. If $C=2$ and $z_N=1$ then $z_{N+1}=3$, etc., . So the numbers z_N possibly are larger than 1 so the larger $1+1 \equiv 2, 1+2 \equiv 3$, etc (defined to be $a+b=c$) and define rules of algebra on these numbers like $a+b=b+a$ (eg., ring-field) with no new axioms. So postulate 0 also generates the big numbers and thereby the algebra we can now use:

If we state different rules than the standard ring-field algebra rules we still get the same physics but using these different math rules in the physics laws.

Postulate 1 also gets us set theory. For example $1 \cup C \equiv 1+C$ (If $A \cap B = \emptyset$). with algebraic definition of $1 = zz$ having both 1,0 as solutions so defining negation \sim with $0=1-1$ Thus we can define intersection \cap with $\sim((A \cup B) \sim B \sim A) \equiv A \cap B$. So we have defined both union \cup and intersection \cap so we have derived set theory.

So in postulate 1 $z=zz$ why did 0 come along for the ride? The deeper reason in set theory is that \emptyset is an element of every set. Note \emptyset and 0 aren't really new postulates since they postulate literally "nothing". So we just derived set theory from the postulate of 1.

Relationship between 0 and \emptyset

The null set \emptyset is the subset of every set. In the more fundamental set theory formulation. \emptyset is not a real number so \emptyset and 0 are not the same. But {some of the properties overlap such as $\emptyset \subset \{\text{all sets}\} \Leftrightarrow \{0\} \subset \{1\}$ since $\emptyset = \emptyset \cup \emptyset \Leftrightarrow 0+0=0, \{\{1\} \cup \emptyset\} = \{1\} \Leftrightarrow 1+0=1$.

So list $1 \cup 1 \equiv 1+1 \equiv 2, 2 \cup 1 \equiv 1+2 \equiv 3, \dots$ all the way up to 10^{82} (as an "application" so we haven't violated Occam's razor. See Feigenbaum point) and define all this list as $a+b=c$, etc., to create our algebra and numbers (rings^fields) which we use to write equation 1 $z=zz+C, \delta C=0$ for example.

D2 Alternative ways of adding 2D+2D→4D

Recall from section 1 that adding the $N=0$ fractal scale 2D δz perturbation to $N=1$ eq.7 2D gives curved space 4D. So $(dx_1+id x_2)+(dx_3+id x_4) \equiv dr+idt$ given (eqs5,7a) $dr^2-dt^2=(\gamma^r dr+i\gamma^t dt)^2$ if $dr^2 \equiv dx^2+dy^2+dz^2$ (3D orthogonality) so that $\gamma^r dr \equiv \gamma^x dx + \gamma^y dy + \gamma^z dz, \gamma^i \gamma^i + \gamma^j \gamma^j = 0, i \neq j, (\gamma^i)^2 = 1$, rewritten (with curved space $\kappa_{\mu\nu}$ eq.14-17)

$$(\gamma^x \sqrt{\kappa_{xx}} dx + \gamma^y \sqrt{\kappa_{yy}} dy + \gamma^z \sqrt{\kappa_{zz}} dz + \gamma^t \sqrt{\kappa_{tt}} idt)^2 = \kappa_{xx} dx^2 + \kappa_{yy} dy^2 + \kappa_{zz} dz^2 - \kappa_{tt} dt^2 = ds^2.$$

But there are alternatives to this 3D orthogonalization method. For example satisfying this 4D Clifford algebra and complex orthogonalization requirement is a special case of any 2 $x_i x_j$ in eq.3 (directly from postulate1): Imposing orthogonality thereby creates 6 pairs of eqs.3&5. So each particle carries around it's own $dr+idt$ complex coordinates with them on their world lines.

Alternatively this 2D $dr+idt$ is a 'hologram' 'illuminated' by a modulated $dr^2+dt^2=ds^2$ 'circle' wave (as 2nd derivative wave equation operators from eq.11 circle) since 4Degrees of freedom are imbedded on a 2D (dr,dt) surface here, with observed coherent superposition output as eq.16 solutions. A more direct way is to simply write the 4Degrees of freedom on the 2D surface as

$dr+idt = (dr_1+idt_1)+(dr_2+idt_2) = (dr_1, \omega dt_2), (dr_2, idt_2) = (x, z, y, idt) = (x, y, z, idt)$, where $\omega dt \equiv dz$ is the z direction spin $\frac{1}{2}$ component ω (angular velocity) axial vector of the Newpde lepton (eqs.7-9); which we get anyway from lepton equation Newpde.

N=-1 and dimensionality

Note the N=-1 (GR) is yet another δz perturbation of N=0 $\delta z'$ perturbation of N=1 observer thereby adding at least 1 independent parameter dimension to our $dx_1+(dx_2+idx_3)+ (dx_4+idx_5)$ (4+1) explaining why Kaluza Klein 5D $R_{ij}=0$ works so well: so GR is really 5D if E&M (N=0) included. Note these N=-1 fractal scale wound up balls at $r_H=10^{-58}m$ are a lot smaller than the Planck length. But if only N=1 observer and N=-1 are used (no N=0) we still have the usual 4D GR Einstein equations. Recall the dx_1 (N=-1) is gravity.

D3 We can isolate lemniscate Mandelbrot Set implied by the perfect circle (eq.11) observability if also 4X circles included.

In section 1 we got the Circle $dr^2+dt^2=ds^2$ and so *observability* of eq.11. So including observability *only* we could have instead postulated $1^2=1^2 1^2$ or $C_{N+1}=C_N C_N + C$. $C=C_1=dr^2+dt^2$, $C_0=0$ instead of the more general $z=zz$ ($1=1X1$) implying $z_{N+1}=z_N z_N + C$. This gets the lemniscate sequence and so just the bare bones Mandelbrot set without all the flourishes of the smaller scale versions of $z_{N+1}=z_N z_N + C$

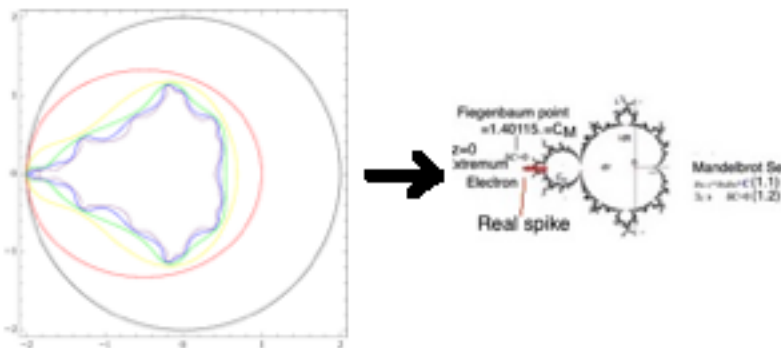


fig7 Lemniscate sequence (Wolfram, Weisstein, Eric) $C_{N+1}=C_N C_N + C$. $C=C_1=dr^2+dt^2$, $C_0=0$. After an infinite number of successive approximations $C''=C'C'+C = C_M^2$

Mandelbrot calls C_M the ER, Escape Radius (see Muency).

Note then *observability* thereby implies *only* the basic Mandelbrot set structure and so not all the other parts, the flourishes, of that zoom.

But the $\delta C=0$ extreme additionally imply states whose life times are long enough to be observable and those are at the $\delta C=0$ extreme of the (observably) 4X circles Feigenbaum point,

Δ Modification of Usual Elementary Calculus ϵ, δ ‘tiny’ definition of the limit.

Recall that: given a number $\epsilon > 0$ there exists a number $\delta > 0$ such that for all x in S satisfying

$$|x-x_0| < \delta$$

we have

$$|f(x)-L| < \epsilon$$

Then write $\lim_{x \rightarrow x_0} f(x) = L$

Thus you can take a smaller and smaller ϵ here, so then $f(x)$ gets closer and closer to L even if x never really reaches x_0 . ‘Tiny’ for $h \rightarrow L_1$ and $f(x+h)-f(x) \rightarrow L_2$ then means that $L=0 = L_1$ and L_2 . ‘Tiny’ is this difference limit.

Hausdorf (Fractal) s dimensional measure using ε, δ

Diameter of U is defined as $|U| = \sup\{|x - y| : x, y \in U\}$. $E \subset \cup_i U_i$ and $0 < |U_i| \leq \delta$

$$H_\delta^s(E) = \inf \sum_{i=1}^{\infty} |U_i|^s$$

analogous to the elementary $V=U^s$ where of $s=3$, $U=L$ then V is the volume of a cube $Volume=L^3$. Here however 's' may be noninteger (eg., fractional). The volume here would be the respective Hausdorf outer measure.

The infimum is over all countable δ covers $\{U_i\}$ of E .

To get the Hausdorf outer measure of E we let $\delta \rightarrow 0$ $H^s(E) = \lim_{\delta \rightarrow 0} H_\delta^s(E)$

The restriction of H^s to the σ field of H^s measurable sets is called a Hausdorf s-dimensional measure. $\dim E$ is called the Hausdorf dimension such that

$$H^s(E) = \infty \text{ if } 0 \leq s < \dim E, \quad H^s(E) = 0 \text{ if } \dim E < s < \infty$$

So if s implies a zero H or infinite H it is not the correct dimension. This rule is analogous to the definition of the (fractal) Mandelbrot set itself in which a C that gave infinity is rejected by the definition $\delta C = 0$ we can model as a binary pulse ($z = z z$ solution is binary $z = 1, 0$) with

$z z = z (1)$ is the algebraic definition of 1 and can add real constant C (so $z' = z' z' - C$, $\delta C = 0$ (2)), $z \in \{z'\}$

Plug $z' = 1 + \delta z$ into eq.2 and get $\delta z + \delta z \delta z = C$ (3)

so $\delta z = (-1 \pm \sqrt{1 + 4C}) / 2 = dr + idt$ (4)

for $C < -1/4$ so real line $r = C$ is immersed in the complex plane.

$z = z_0 = 0$ To find C itself substitute z' on left (eq.2) into right $z' z'$ repeatedly & get $z_{N+1} = z_N z_N - C$. $\delta C = 0$ requires us to reject the Cs for which

$-\delta C = \delta(z_{N+1} - z_N z_N) = \delta(\infty - \infty) \neq 0$. $z = z z$ solution is **1, 0** so initial

gets the **Mandelbrot set** C_M (fig2) out to some $\|\Delta\|$ distance from $C = 0$. Δ found from $\partial C / \partial t = 0$, $\delta C \equiv \delta C_r = (\partial C_M / \partial (dr dt)) dr = 0$ extreme giving the Feigenbaum point $\|C_M\| = \|-1.400115..\|$ global max given this $\|C_M\|$ is biggest of all.

If s is not an integer then the dimensionality it is has a fractal dimension.

But because the Feigenbaum point Δ uncertainty limit is the r_H horizon, which is impenetrable (sect.2.5, partI), ϵ, δ are not dr/ds eq.11a observables for $0 < \epsilon, \delta < r_H$. Instead $\epsilon, \delta > \Delta = r_H$ = the next $10^{40} \times$ smaller fractal scale Mandelbrot set at the Feigenbaum point.

Review Recall from eq.7 that $dr + dt = ds$. So combining in quadrature eqs 7&11

$SNR \delta z = (dr/ds + dt/ds) \delta z = ((dr + dt)/ds) \delta z = (1) \delta z$ (11c, append) and so having come *full circle* back to postulate 1 as a real eigenvalue ($1 \equiv$ Newpde electron). For all the rotations in fig.4 (except the eq.11 IVth to Ist quadrants: in eq.B1 each quadrant rotation provides one derivative for each ν

$\nu \left[\left(\frac{dr + dt}{ds} \right) + \left(\frac{dr + dt}{ds} \right) \right] \delta z = 2 \frac{ds}{ds} \delta z = 2(1) \delta z$ Equation 11 (sect.1) then counts units N of each 2 half integer $S = 1/2$ angular momentums = 1 = 2 units of electrons (spin1 for W and Z) off the light

cone. Alternatively diagonal $ds = \sqrt{2} dr$ in $\int \left(\frac{dr}{\sqrt{2} dr} + \frac{dr}{\sqrt{2} dr} \right)^2 dV = 1$ For the rotation in the eq.11

IVth to Ist quadrants (each quadrant rotation provides one derivative for each ν) at 45° $dr = dt$ (on the light cone in fig.4) so for Hamiltonian H: $2H \delta z = 2(dt/ds) \delta z = 2(1/2) \delta z = (1) \hbar \omega \delta z = \hbar c k \delta z$ on the diagonal so that $E = p c = \hbar \omega$ for the two ν energy components, universally. Thus we can state the most beautiful result in physics that $E = N h f$ for the energy of light with N equal N

monochromatic photons. Replaces 2nd quantization of 2 given allowed Newpde 10^{82} electrons(appendix A2) So we really do have a binary physics signal. So, having come *full circle* then: (**postulate 0** \Leftrightarrow Newpde)

Digital communication analogy: Binary ($z=zz$) 1,0 signal with white noise $\delta C=0$ in $z'+C=z'z'$. Recall the algebraic definition of 1 is $z=zz$ which has solutions 1,0.(11c). Boolean algebra. Also you could say white noise C has a variation of zero ($\delta C=0$) making it easy to filter out (eg., with a Fourier cutoff filter). So you could easily make the simple digital communication analogy of this being a binary ($z=zz$) 1,0 signal with white noise $\delta C=0$ in $z'+C=z'z'$. (However the noise is added a little differently here ($z+C=zz$) than in statistical mechanics signal theory (eg., There you might use deconvolved signal=convolution integral [(transfer function)signal]dA)). where the 'signal' actually would equal $z+C$, not the usual $(2J_1(r)/r)^2$ psf So this is not quite the same math as in signal theory statistics statistical mechanics.)

The Whole Shebang:

This theory is 0

Postulate real number **0** if $z'=0$ and $z'=1$ plugged into $z'=z'z'+C$ (**eq.1**) results in some $C=0$ constant(ie $\delta C=0$)

Plug 0 into **eq.1** and get the Mandelbrot set

Plug 1 into **eq.1** and get the Dirac eq.

Dirac plus Mandelbrot gets the Newpde

So Ultimate Occam's razor postulate(**0**) implies ultimate math-physics

So this theory is **0**. Hold that thought.