

It's Broken, fix it

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Key words, Mandelbrot set, Dirac equation, Metric

Abstract In that regard Dirac in 1928 made his equation(1) flat space(2). But space is not in general flat, there are forces.

So over the past 100 years people have had to try to make up for that mistake by adding ad hoc convoluted gauge force after gauge force until fundamental theoretical physics became a mass of confusion, a train wreck, a junk pile. So all they can do for ever and ever is to rearrange that junk pile with *zero actual progress* in the most fundamental theoretical physics* ,.. forever. We died.

By the way note that Newpde(3) $\gamma^\mu \sqrt{\kappa_{\mu\mu}} \partial\psi/\partial x_\mu = (\omega/c) \psi$ is NOT flat space (4) so it cures this problem (5).

References

(1) $\gamma^\mu \partial\psi/\partial x_\mu = (\omega/c) \psi$

(2) Spherical symmetry: $(\gamma^x \sqrt{\kappa_{xx}} dx + \gamma^y \sqrt{\kappa_{yy}} dy + \gamma^z \sqrt{\kappa_{zz}} dz + \gamma^t \sqrt{\kappa_{tt}} dt)^2 = \kappa_{xx} dx^2 + \kappa_{yy} dy^2 + \kappa_{zz} dz^2 - \kappa_{tt} dt^2 = ds^2$
 $\kappa_{xx} = \kappa_{yy} = \kappa_{zz} = \kappa_{tt} = 1$ is flat space, Minkowski, as in his Dirac equation(1).

(3) Newpde: $\gamma^\mu \sqrt{\kappa_{\mu\mu}} \partial\psi/\partial x_\mu = (\omega/c) \psi$ for e,v. So we didn't just drop the $\kappa_{\mu\nu}$ (as is done in ref.1)

(4) Here $\kappa_{oo} = 1 - r_H/r = 1/\kappa_{rr}$, $r_H = (2e^2)(10^{40N})/(mc^2)$. The $N = \dots -1, 0, 1, \dots$ fractal scales (next page)

(5) This Newpde κ_{ij} contains a Mandelbrot set(6) $e^2 10^{40N}$ Nth fractal scale source(fig1) term (from eq.13) that also successfully **unifies theoretical physics**. For example:

For $N = -1$ (i.e., $e^2 \times 10^{-40} \equiv G m_e^2$) κ_{ij} is then by inspection(4) the Schwarzschild metric g_{ij} ; so we just derived General Relativity and the gravity constant G from Quantum Mechanics in one line Wow
For $N = 1$ (so $r < r_c$) Newpde zitterbewegung expansion stage explains the universe expansion (For $r > r_c$ it's not observed, per Schrodinger's 1932 paper.).

For $N = 1$ zitterbewegung harmonic coordinates and Minkowski metric submanifold (after long time expansion) gets the De Sitter ambient metric we observe (D16, 6.2).

For $N = 0$ Newpde $r = r_H$ $2P_{3/2}$ state composite $3e$ is the baryons (sect.2, partII) and Newpde $r = r_H$ composite e,v is the 4 Standard electroweak Model Bosons (4 eq.12 rotations \rightarrow appendix A)

for $N = 0$ the higher order Taylor expansion(terms) of $\sqrt{\kappa_{ij}}$ gives the anomalous gyromagnetic ratio and Lamb shift *without* the renormalization and infinities (appendix D3): This is very important

So κ_{uv} provides the general covariance of the Newpde. Eq. 4 even provides us space-time r,t.

So we got all physics here by *mere inspection* of this (curved space) Newpde with no gauges!

We fixed it.

So where does that Newpde come from that fixed it? We just do **what is simplest** (let Occam be your guide), just **postulate 1**: the physics (Newpde) will then follow, top down:

So this Occam's razor *simplicity* requirement motivates every step. We later *derive* from our postulate 1 the eq.11 real eigenvalues (Hermitian) operator formalism so we can thereby define the **limit 1** as $1 \pm \text{uncertainty}$ and so (an observable) indistinguishable from 1 since it is inside the uncertainty bounds. Note also an 'observable' implies an 'observer'.

Postulate 1 and define limit 1

abstract Let Occam be your guide to get physics top down, So startout simplest with **Postulate 1**
The simplest algebraic definition of **1** is $z=zz$ so **define limit 1** as $z=zz+C$ for small constant C

‘constant’ defined as $\delta C=0$;

‘small’ defined as $\|C\| < 1$

solution of $z=zz$ is $z=\mathbf{1}, 0$;

Postulate 1 (i.e., $z=zz$) requires *some* $C=0$ (so $z=\mathbf{1}, 0=z'$) in the solution set of $z'=z'z'+C$ (eq1)
davidmaker

so $z=\mathbf{0}=z_0=z'$ To find *all* C substitute z' on left (eq.1) into right $z'z'$ repeatedly and get iteration
 $z_{N+1}=z_N z_N - C$. Constraint $\delta C=\mathbf{0}$ requires we reject the C s for $-\delta C=\delta(z_{N+1}-z_N z_N)=\delta(\infty-\infty)\neq 0$. The C s
that are left over define the **Mandelbrot set** C_M . eg., $\delta z'=10^{40N}\delta z$, $N=1$ fractal scale(\equiv observer)
 $z=\mathbf{1}$ in $z'=\mathbf{1}+\delta z$ in eq.1 get $\delta z+\delta z\delta z=C$ (3) so $\delta z=(-1\pm\sqrt{1+4C})/2=dr+idt$ if $C<-1/4$ (complex) (4)

$z=\mathbf{1}$ in $z'=\mathbf{1}+\delta z$ in eq1 gives for *required* $N=1$ (so $|\delta z| >> 1$) and $\delta C=\mathbf{0}=\delta(\delta z+\delta z\delta z)=\delta\delta z(1$
 $+ \delta\delta z(\delta z)+(\delta z)\delta\delta z\approx\delta(\delta z\delta z)=0$ =(plug in eq.4) $=\delta[(dr+idt)(dr+idt)]=\delta[(dr^2-dt^2)+i(dr dt+dt dr)]=0$ (5)
 $=2D$ (Minkowski metric, $c=1$) $+i$ (Clifford algebra \rightarrow eq.7a) (\equiv Dirac eq)

Factor eq.5 real $\delta(dr^2-dt^2)=\delta[(dr+dt)(dr-dt)]=0=[\delta(dr+dt)](dr-dt)+[(dr+dt)](\delta(dr-dt))=0$ (6)

so $-dr+dt=ds$, $-dr-dt=ds\equiv ds_1(\rightarrow\pm e)$ Squaring&eq.5 gives circle.in e, v (dr, dt) 2nd, 3rd quadrants (7)

& $dr+dt=ds$, $dr-dt=ds$, $dr\pm dt=0$, light cone ($\rightarrow v, v$) in *same* e, v (dr, dt) plane 1st, 4th quadrants (8)

& $dr+dt=0, dr-dt=0$ so $dr=dt=0$ defines vacuum (9)

Quadrants give *positive* scalar $drdt$ of eq.7 (if *not* vacuum) imply the eq.5 *non* infinite extremum
imaginary $\equiv drdt+dt dr=0\equiv\gamma^i dr\gamma^j dt+\gamma^j dt\gamma^i dr=(\gamma^i\gamma^j+\gamma^j\gamma^i)drdt$ so $(\gamma^i\gamma^j+\gamma^j\gamma^i)=0$, $i\neq j$ (from *releq4* $\gamma^i\gamma^i=1$)(7a)

Thus from eqs5,7a: $dr^2-dt^2=(\gamma^r dr+i\gamma^t dt)^2$

Both $z=\mathbf{0}, z=\mathbf{1}$ together using orthogonality to get $(2D+2D)$ curved space). Thus $(z=1)+(z=0)=(dx_1+idx_2)+(dx_3+idx_4)\equiv dr+idt$ given $dr^2-dt^2=(\gamma^r dr+i\gamma^t dt)^2$ if $dr^2\equiv dx^2+dy^2+dz^2$ (orthogonality)so that
 $\gamma^r dr\equiv\gamma^x dx+\gamma^y dy+\gamma^z dz$, $\gamma^i\gamma^j+\gamma^j\gamma^i=0$, $i\neq j$, $(\gamma^i)^2=1$, rewritten (κ_{ii} from $N=0$ C_M perturbation of $N=1$ eq.7)
 $(\gamma^x\sqrt{\kappa_{xx}}dx+\gamma^y\sqrt{\kappa_{yy}}dy+\gamma^z\sqrt{\kappa_{zz}}dz+\gamma^t\sqrt{\kappa_{tt}}idt)^2=\kappa_{xx}dx^2+\kappa_{yy}dy^2+\kappa_{zz}dz^2-\kappa_{tt}dt^2=ds^2$. Multiply both sides by
 $1/ds^2$ and $\delta z^2\equiv\psi^2$ use circle $-i\partial\delta z/\partial r=(dr/ds)\delta z$ inside brackets() get $4D$ $\gamma^\mu(\sqrt{\kappa_{\mu\mu}})\partial\psi/\partial x_\mu=(\omega/c)\psi$
 \equiv **Newpde** for e, v , $\kappa_{00}=1-r_H/r=1/\kappa_{rr}$, $r_H=e^2\times 10^{40N}/m$ ($N=-1, 0, 1, \dots$). Also $C_M/\xi=r_H=\text{small } C$ sobig $\xi=\gamma$
boost so $z=zz$ so **postulate 1**. So we really did just postulate 1. So **Postulate 1** \rightarrow **Newpde**

Applications of $N=0$ $r=r_H$: **Newpde** composite $3e$ $2P_{3/2}$ state = baryons and the 4 **Newpde** e, v
(quadrant)rotations are the 4 W^+, γ, W^-, Z_0 , SM Bosons. Also $N=-1$ is GR & big $\gamma=\xi=\tau+\mu$ from C_M
So “postulate 1” gives **eq1** (i.e., all of physics) and real#math, everything. See davidmaker.com

postulate 1 and define limit 1 ($\equiv 1 \pm \text{uncertainty}$)

define limit 1 Algebraic definition of 1 is $z=zz$ so **limit 1** is **defined** as $z=zz+C$ for small constant C
 'small' defined as $\|C\| \ll 1$
 'constant' defined as $\delta C=0$
 solution to $z=zz$: $z=1,0$

Postulate 1 ($z=zz$) requires some $C=0$ (so $z=1,0=z'$) in the solution set of $z'=z'z'+C$ (**eq1**)
 So plug $z=1,0$ into **eq1** (use $\delta C=0$)
 into **eq1** iteration (to get all C) get 2D Mandelbrot (fractal scale N)
 into **eq1** get 2D Dirac ($(N=1) \equiv \text{observer}$)
 4D Newpde using orthogonalization (with small C implied)

Summary **eq1 so** (it is just eq1)

z=0
z=1
both

therefore

postulate 1 \rightarrow Newpde

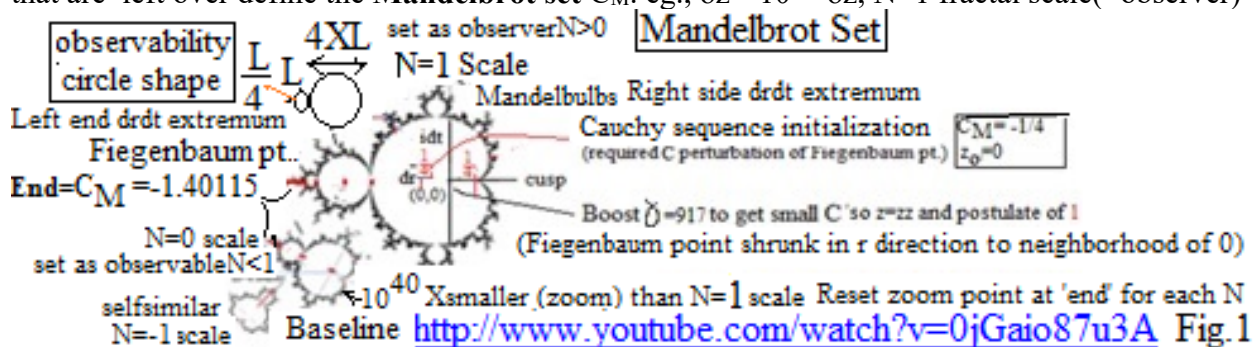
(So we get all of physics (davidmaker.com) and 1U algebra and Real#math (1 iteration is Cauchy) everything) that is physical and nothing else.

I Math Details of Postulate 1 and define limit 1

abstract Let Occam be your guide to get physics top down, So startout simplest with **Postulate 1**
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 so $z=0=z_0=z'$ To find *all* C substitute z' on left (**eq.1**) into right $z'z'$ repeatedly and get iteration
 $z_{N+1}=z_N z_N - C$. Constraint $\delta C=0$ requires we reject the Cs for $-\delta C=\delta(z_{N+1}-z_N z_N)=\delta(\infty-\infty) \neq 0$. The Cs
 that are left over define the **Mandelbrot set** C_M . eg., $\delta z'=10^{40N} \delta z$, $N=1$ fractal scale ($\equiv \text{observer}$)



$z=1$ in $z'=1+\delta z$ in **eq.1** get $\delta z+\delta z \delta z=C$ (3) so $\delta z=(-1 \pm \sqrt{1+4C})/2 = dr+idt$ if $C < -1/4$ (complex) (4)

$z=1$ in $z'=1+\delta z$ in **eq1** gives for *required* $N=1$ (so $|\delta z'| \gg 1$) and $\delta C=0=\delta(\delta z+\delta z \delta z)=\delta \delta z(1+\delta \delta z(\delta z)+(\delta z)\delta \delta z \approx \delta(\delta z \delta z)=0$ (plug in **eq.4**) $=\delta[(dr+idt)(dr+idt)]=\delta[(dr^2-dt^2)+i(dr dt+dt dr)]=0$ (5)
 $=(\text{Minkowski metric, } c=1)+i(\text{Clifford algebra} \rightarrow \text{eq.7a})$ ($\equiv \text{Dirac eq}$)

Factor **eq.5 real** $\delta(dr^2-dt^2)=\delta[(dr+dt)(dr-dt)]=0=[\delta(dr+dt)](dr-dt)+[(dr+dt)[\delta(dr-dt)]=0$ (6)

so $-dr+dt=ds, -dr-dt=ds \Rightarrow ds_1 (\rightarrow \pm e)$ Squaring **eq.5** gives circle in e, v (dr, dt) 2nd, 3rd quadrants (7)

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Quadrants give *positive* scalar $drdt$ of eq.7 (not eq.9 vacuum) imply eq.5 *non* infinite extremum *imaginary* $\equiv drdt+dtdr=0 \equiv \gamma^i dr^j dt + \gamma^j dt^i dr = (\gamma^i \gamma^j + \gamma^j \gamma^i) drdt$ so $(\gamma^i \gamma^j + \gamma^j \gamma^i)=0, i \neq j$ (from $\text{releq4 } \gamma^i \gamma^i=1$)(7a)

We square eqs.7 or 8 or 9 $ds_1^2=(dr+dt)(dr+dt)=(-dr-dt)(-dr-dt)=[dr^2+dt^2]+(drdt+dtdr) \equiv ds^2+ds_3=ds_1^2$. **Circle** $\equiv \delta z=dse^{i\theta}=dse^{i(\Delta\theta+\theta_0)}=dse^{i((\cos\theta dr+\sin\theta dt)/(ds)+\theta_0)}$, $\theta_0=45^\circ$ (δz in fig.7). We define $k \equiv dr/ds$, $\omega \equiv dt/ds$, $\sin\theta \equiv r$, $\cos\theta \equiv t$. $dse^{i45^\circ} \equiv ds'$. Take ordinary derivative dr (since (flat space)

$$\text{of 'Circle' } \frac{\partial \left(dse^{i\left(\frac{rdr}{ds} + \frac{tdt}{ds}\right)} \right)}{\partial r} = i \frac{dr}{ds} \delta z \text{ so } \frac{\partial (dse^{i(rk+wt)})}{\partial r} = ik\delta z, \quad k\delta z = -i \frac{\partial \delta z}{\partial r} \quad (11).$$

(So given $\delta z \equiv \psi$, $F \equiv k$ then from eq.11 $\langle F \rangle^* = \int (F\psi)^* \psi d\tau = \int \psi^* F \psi d\tau = \langle F \rangle$. Therefore k is Hermitian). Also from right side real# Cauchy seq. starting at $-1/4$ iteration, is the same as the the Mandelbrot set iteration(7), Ch.2,sect.2, with small $C=0$ limit making *real eigenvalues* (eg.,noise) likely. Thus the Mandelbrot set iteration here did double duty also as proof of the real number eigenvalues in eq.11. The observables $dr \rightarrow k \rightarrow p_r$ condition gotten from eq.11 **operator formalism**(10) thereby converts eq.7-9 into Dirac eq. pdes (4XCircle solution in left side fig.1 also implies observability). Thus our title 'postulate 1 (observable)' implicitly contains the eq.11 QM definition of observability. Cancel that e^{i45° coefficient ($45^\circ=\pi/4$) then multiply both sides of eq.11 by \hbar and define $\delta z \equiv \psi$, $p_r \equiv \hbar k$. Eq.11: the familiar $p_r \psi = i\hbar \frac{\partial \psi}{\partial r}$ (11). Repeat eq.3 for the τ, μ respective δz lobes in fig.6 so they each have their own neutrino ν .

$\delta C=0$ Extremum on Circle 4X sequence shapes (fig1) In Mandelbrot set pulls it out of zoom clutter because of the above 4X circle observability sequence

$\delta C=0$ gives that 45° extreme but it also applies to *local* constants (extremum peaks and valleys)

because $\delta C = \left(\frac{\partial C}{\partial r} \right)_t dr + \left(\frac{\partial C}{\partial t} \right)_r idt = 0$. So for that fig.1 4X sequence of circles $drdt=$

$darea_M \neq 0$ (so eq.11a observables) the real $\delta C=0$ extremum given the decreasing circle radius

sequence $\lim_{m \rightarrow \infty} \frac{\partial C}{\partial area_m} dr_m = KX0 = 0$ (since $dr_\infty \approx 0$) at Feigenbaum point $= f^\alpha = (-1.40115, i0) =$

$C_M \equiv \text{end}$. So random circles in the zoom don't do $\delta C=0$. Note if a circle (or many circles) is

rotated (U), translated (D), shrunk (S) equally in both dimensions (i.e., $(\partial x^j / \partial x'^k) f^j = f'^k \equiv \begin{bmatrix} f_{1N} \\ f_{2N} \end{bmatrix} =$

$S_N \begin{bmatrix} U_{11} & U_{12} \\ U_{21} & U_{22} \end{bmatrix} \begin{bmatrix} f \\ 0 \end{bmatrix} + \begin{bmatrix} D_{1N} \\ D_{2N} \end{bmatrix}$) it is still a circle, eq.11 still holds, so *it's still an observable* as seen

in the N fractal scale zoom. Thus you can pick out from that zoom these fig.1 Mandelbrot set extremum 4Xdiameter circles as the only **observables** and $\delta C=0$ extremum geometry in all that clutter. Reset the zoom, restart at such $S_N C_M = 10^{40N} C_M$ in eq.13.

1b $z=1, z=0$ steps combined (on Circle with small C boost) using orthogonality(ortho):

Postulate 1 also implies a small C in eq.1 which thereby implies a (Minkowski metric Lorentz contraction(9)) $1/\gamma$ boosted frame of reference(fig.6) in the eq.3 $C=C_M/\gamma \equiv C_M/\xi_1 = \delta z' = \Delta$ for next small smaller fractal scale $N_{ob} < 0$ so $\delta z' \ll 1$ (composite 3e: sect.2 and PartII). For $N=0$ eq.5 (which is true only for $N>0$) and so eq.7 is not quite true (and δz in eq.11 perturbed). But we keep $\delta ds^2=0$ (circle) in eq.5, on the 4X circles so we must have an angle perturbation of big $N=1$ dr, dt for $\theta_0=45^\circ$ on above **ds Circle** and so a slightly modified eq.7:

$$(dr-\delta z') + (dt+\delta z') \equiv dr' + dt' = ds \quad (12)$$

$N_{ob}=0$ extremum eq.12 rotations (observer at $N=1$, eq.7 $dr+dt=ds$ constraint)

Recall for $N_{ob}=0$ (observer at $N=1$) and eq. 7 $dr+dt=ds$ the r,t axis' are the max extremum for ds^2 , and the ds^2 at 45° is the min extremum ds^2 so each $\Delta\theta=\pm 45^\circ$ is pinned to an axis' so extreme $\Delta\theta\approx\pm 45^\circ=\delta z'$. So in eq.12 the 4 rotations $45^\circ+45^\circ=90^\circ$ define 4 Bosons (appendix A).

But for $45^\circ-45^\circ$ $N_{ob}<0$ then contributes so you also have other (smaller) fractal scale extreme $\delta z'$ (eg.,tiny Fiegenbaum pts so $N=1$ $dr=r$, for $N_{ob}<0$) so metric coefficient $\kappa_{rr}=(dr/dr')^2=(dr/(dr-(C_M/\xi_1)))^2=1/(1-r_H/r)^2=A_1/(1-r_H/r)+A_2/(1-r_H/r)^2$. The partial fractions A_i can be split off from RN and so

$$\kappa_{rr}\approx 1/[1-((C_M/\xi_1)r)] \quad (13)$$

$$(C_M \text{ defined to be } e^2 \text{ charge, } \gamma=\xi_1 \text{ mass}). \text{ So: } ds^2=\kappa_{rr}dr'^2+\kappa_{oo}dt'^2 \quad (14)$$

$$\text{From eq.7a } dr'dt'=\sqrt{\kappa_{rr}}dr'\sqrt{\kappa_{oo}}dt'=drdt \text{ so } \kappa_{rr}=1/\kappa_{oo} \quad (15)$$

We can then do a rotational dyadic coordinate transformation of $\kappa_{\mu\nu}$ to get the Kerr metric which is all we need for our applications(9).

Recall from eqs5,7a that $dr^2-dt^2=(\gamma^r dr+i\gamma^t dt)^2$

Both $z=0, z=1$ together using orthogonality to get $(2D+2D_{curved} \text{ space})$. Thus $(z=1)+(z=0)=(dx_1+id x_2)+(dx_3+id x_4)\equiv dr+idt$ given $dr^2-dt^2=(\gamma^r dr+i\gamma^t dt)^2$ if $dr^2\equiv dx^2+dy^2+dz^2$ (orthogonality) so that $\gamma^r dr\equiv \gamma^x dx+\gamma^y dy+\gamma^z dz$, $\gamma^i \gamma^j+\gamma^j \gamma^i=0$, $i\neq j, (\gamma^i)^2=1$ (B2), rewritten (with eq14)

$(\gamma^x \sqrt{\kappa_{xx}} dx+\gamma^y \sqrt{\kappa_{yy}} dy+\gamma^z \sqrt{\kappa_{zz}} dz+\gamma^t \sqrt{\kappa_{tt}} dt)^2=\kappa_{xx} dx^2+\kappa_{yy} dy^2+\kappa_{zz} dz^2-\kappa_{tt} dt^2=ds^2$. Multiply both sides by $1/ds^2$ & $(\delta z/\sqrt{dV})^2\equiv \psi^2$ and using operator eq 11 inside the brackets() get

Newpde $\gamma^\mu(\sqrt{\kappa_{\mu\mu}})\partial\psi/\partial x_\mu=(\omega/c)\psi$ for e, ν , $\kappa_{oo}=1-r_H/r=1/\kappa_{rr}$ $r_H=e^2 X 10^{40N}/m$ ($N=-, 0, 1, \dots$) (16)
 $=C_M/\xi_1$ (from* eq.13) C_M =Fiegenbaum point. So: **postulate 1** \rightarrow **Newpde.** syllogism

* C_M/ξ_1 is ξ small C boost for $z=zz$ so postulate1 from Newpde $r=r_H$ $2P_{3/2}$ stable state. See fig6. The 4 eq.12 Newpde e, ν rotations at $r=r_H$ are the 4 W^+, γ, W^-, Z_0 SM Bosons (appendixA).

So Penrose's intuition(6) was right on! There *is* physics in the Mandelbrot set, all of it.

2.1 Oscillation of $\delta z(\equiv\psi)$ on a given fractal scale

From Newpde (eg., eq.1.13 Bjorken and Drell) $i\hbar \frac{\partial\psi}{\partial t} = \frac{\hbar c}{i} \left(\alpha_1 \frac{\partial\psi}{\partial x^1} + \alpha_2 \frac{\partial\psi}{\partial x^2} + \alpha_3 \frac{\partial\psi}{\partial x^3} \right) +$

$\beta mc^2 \psi = H\psi$. For electron at rest: $i\hbar \frac{\partial\psi}{\partial t} = \beta mc^2 \psi$ so: $\delta z = \psi_r = w^r(0)e^{-i\varepsilon_r \frac{mc^2}{\hbar} t}$ $\varepsilon_r=+1,$

$r=1,2; \varepsilon_r=-1, r=3,4$.): This implies an oscillation frequency of $\omega=mc^2/\hbar$. which is fractal here. So

the eq.12 the 45° line has this ω oscillation as a (that eq.7-9 δz variation) rotation. On our own fractal cosmological scale we are in the expansion stage of one such oscillation. Thus the

fractalness of the Newpde explains cosmology. The next higher cosmological scale is independent (but still connected by superposition of speeds implying a separation of variables

result: $i\hbar \frac{\partial\psi}{\partial t} = \beta \sum_N (10^{40N} (\omega t)_{\varepsilon+\Delta\varepsilon}) \psi = \beta \sum_N (10^{40N} m_{\varepsilon+\Delta\varepsilon} c^2/\hbar) \psi$). By the way fractal

scale $N=1$ the 45° small Mandelbulb chord ε (Fig6) is now, given this ω , getting larger with

time so $t \propto \varepsilon$. So cosmologically for stationary $N=1$ $\delta z=\sqrt{\kappa_{oo}}dt= e^{-i\varepsilon_r \frac{mc^2}{\hbar} t} \rightarrow e^{i(\varepsilon+\Delta\varepsilon)}$ (17)

This $N=0$ and $N=-1$ δz is the source of the small rotation in eq.12. Later we see that high energy scattering drives the $\delta\delta z$ term ($/ds$) to the big $\Delta 45^\circ$ extreme (so preferred) jumps (appendixA).

2.2 $N=0$ small C boost so postulate observable1 (e) Postulate 1 (observable) requires

that $C\approx 0$ in equation 1. Recall eq.3 $\delta z+\delta z\delta z=C$. So for $N=1$ observer $|\delta z| \gg 1$ so $\delta z\delta z=C$. Given

eq.3 for $N=0$ $|\delta z| \gg |\delta z\delta z|$, $C\approx\delta z$ sect.1 for $N=0$. Note also our above circle e electron $\Delta\varepsilon$

intersection ground state is at 45° (2^{nd} & 3^{rd} quadrants is from minimum ds^2). Following the

energy increase for Newpde states $\mu 1S_{1/2}$ then is not a constant in time because of eq.12 angle Newpde zitterbewegung variable contribution (sect.2.1) to the δz chord perturbation of the 45° (fig6 below.) The $68.7^\circ = \text{Arctan}(\delta z/C_M)$ is from eq.4 quadratic equation at the Fiegenbaum point.(so it gives our *fundamental* (excited state Mandelbulb) mass $\tau 2S_{1/2}$ that never changes allowing us to normalize it to 1). Next is the Newpde composite $3e 2P_{3/2}$ state at $r=r_H$. with its quantized B flux giving a large γ for each of the 2 positrons. Since the S and P states don't interact here (one is free space) these states are degenerate giving the same reduced mass for each $(\tau+\mu)/2 = \gamma m_e + \gamma m_e = m_p$. Note also that the **real** component of eq.5 is the Minkowski metric implying these γ boosts. For $N=1$ eq.3 implies $C=\delta z \delta z/\xi$ so that $1/\delta z \delta z = 1/(\text{Mandelbulb radius})^2 = \text{mass}$ (from fig6). or as a fraction of τ , electron $\Delta \epsilon = .00058$ (18)

But γ (observer) $= \gamma$ (observable) so for the $N=0$ observable we got the γ from the $N=1$ observer case in $r_H = C_M/\gamma = C_M/\xi = C$ for small C and so postulate 1.

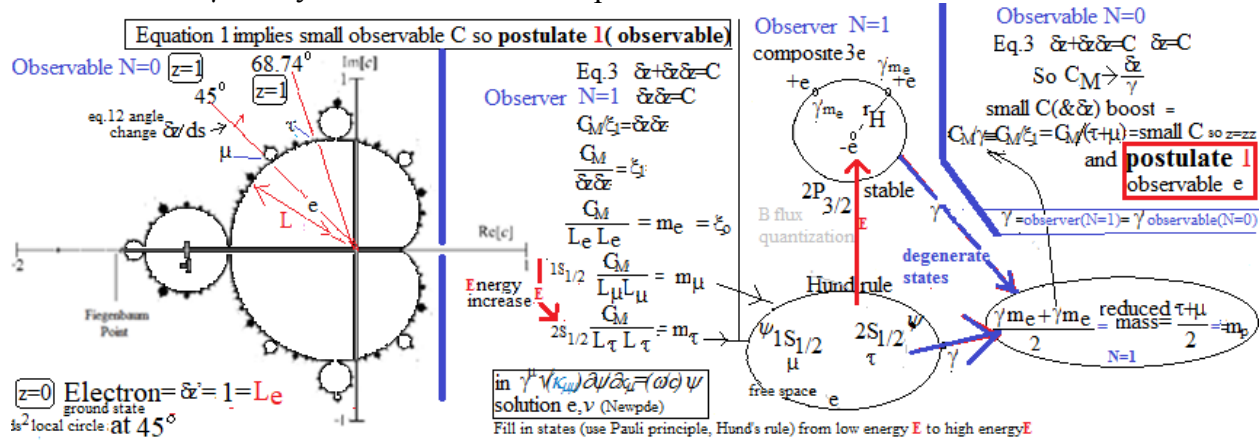


Fig.6 Conclusion

So the **small C** at the end was required. So we really did just **postulate 1**

So we just do *what is simplest* (let Occam be your guide), just **postulate 1**: the physics (Newpde) will then follow, top down:

* **Ultimate Occam's Razor** (observable)

It means here *ultimate* simplicity, the *simplest* idea imaginable. So for example $z=zz$ is *simpler* than $z=zzzz$. Therefore **1** in this context (uniquely algebraically defined by $z=zz$) is this ultimate Occam's razor **postulate** since 0 (also from $z=zz$) postulates literally *nothing*.

3e Stability: We can actually calculate m_p from the quantization of the magnetic flux

$h/2e = \Phi_0 = BA$..Using the Mandelbrot set $2m_p = \tau + \mu + e = \xi_1$ (18)

written as $e^{i(\epsilon + \Delta \epsilon)}$ (eq.D9 outside r_H in D18 hence the 'i') This relation with h just sets the h . and

the Mercuron equation D15 for μ requires the location of object B to find the actual magnitude of m_e .(eq.D9). *Stability* is also implied by $(dt')^2 = (1-r_H/r)dt^2$ since clocks stop at $r=r_H$. That 3rd mass

also reverses the pair annihilation with virtual pair creation inside the r_H volume given

$\sigma = \pi r_H^2 \approx (1/20)\text{barns}$ which is the reason why only composite $3e$ gives stability and not other

larger composites(except multiples of $3e$ itself i.e.,nucleons). Note here we also derived baryon

physics (m_p). See Chapters 8,9,10 Frobenius series solution. This composite $3e$ stability then

makes this small C limit always true and so our **postulate 1** is always true.

2.2 Left end small $d\text{rdt}$ (eq.6) extremum **Fiegenbaum point** Fractalness

The Feigenbaum point (11a) is the only part of the Mandelbrot set we use. At the Feigenbaum point (imaginary) time $X10^{-40}=\Delta$ and real -1.40115 . Since $|C_M| \gg 0$ in eq.2 postulated eq.1 $z=zz$ implies a boosted SR Lorentz transformation universal reference frame to random (since this transformation cancels noise C in eq.2, fig6), small C_M subset $C \approx \delta z'$ (from eq.3) =real distance =real $\delta z/\gamma = 1.4011/\gamma = C_M/\gamma \equiv C_M/\xi_1$ using large ξ_1 . Note at the Feigenbaum point distance $1.4011/\gamma$ shrinks a lot but time $X10^{-40}\gamma$ doesn't get much bigger since it was so small to begin with at the Feigenbaum point. Eq.1 then means we have Ockam's razor optimized **postulated 1**. Given the New pde r_H we only see the $r_H = e^2 10^{40N}/m$ sources from our $N=0$ observer baseline. We never see the $r < r_H$ <http://www.youtube.com/watch?v=0jGaio87u3A> which explores the Mandelbrot set interior near the Feigenbaum point. Reset the zoom start at such extremum $S_N C_M = 10^{40N} C_M$ in eq.13. The splits are in 3 directions from the orbs. There appear to be about 2.5 splits going by each second (given my PC baud rate) and the next Mandelbrot set comes up in about 62 seconds. So $3^{2.7 \times 62} = 10^N$ so $172 \log 3 = N = 82$. So there are 10^{82} splits. So there are about 10^{82} splits per initial split. But each of these Mandelbrot set Feigenbaum points is a $C_M/\xi \equiv r_H$ in electron (eq.13 above). So for each larger electron there are **10^{82} constituent electrons**. Also the scale difference between Mandelbrot sets as seen in the zoom is about **10^{40} , the scale change** between the classical electron radius and 10^{11}ly with the C noising giving us our fractal universe.

Recall again we got from eq.3 $\delta z + \delta z \delta z = C$ with quadratic equation result:

$\delta z = \frac{-1 \pm \sqrt{1-4C}}{2}$. is real for noise $C < 1/4$ creating our noise on the $N=0$ th fractal scale. So $1/4 = (3/2)kT/(m_p c^2)$. So T is 20MK. So here we have *derived the average temperature of the universe* (stellar average). That $z' = 1 + \delta z$ substitution also introduces Lorentz transformation rotational and translation noise that does not effect the number of splits, analogous to how a homeomorphism does not change the number of holes (which is a Topological invariant). So the excess C noise (due to that small C' boost) causes the Feigenbaum point neighborhood internal structure to become randomized (as our present universe is) but the number of electrons (10^{82}) remains invariant. See appendix D mixed state case2 for further organizational effects. $N = r^D$. So the **fractal dimension** = $D = \log N / \log r = \log(\text{splits}) / \log(\#r_H \text{ in scale jump})$ = $\log 10^{80} / \log 10^{40} = \log(10^{40})^2 / \log(10^{40}) = 2$. (See appendix E for Hausdorff dimension & measure) which is the same as the 2D of eq.4 and the Mandelbrot set. The next smaller (subatomic) fractal scale $r_1 = r_H = 2e^2/m_e c^2$, $N=0$ th, $r_2 = r_H = 2GM/c^2$ is defined as the $N=1$ th where $M = 10^{82} m_e$ with $r_2 = 10^{40} r_1$ So the Feigenbaum pt. gave us a lot of physics:
eg. #of electrons in the universe, the universe size, temp.

2.2 Results: What makes this all work is **postulate 1** ($z=zz$ so $z=1,0$). Add constant C ($\delta C=0$) and get $z' = z'z' + C$ (eq1). $z=zz$ postulated so $z=1,0 \in \{z'\}$. (Get **Postulate 1** \rightarrow Newpde). Postulate 1 is the simplest idea imaginable: a Occam's razor theory. We *also* get the *actual* physics with the Newpde (Therefore the usual postulating of hundreds of Lagrange densities, free parameters, dimensions ,etc., is senseless.).

For example (appendixC) *Newpde composite 3e* $2P_{3/2}$ at $r=r_H$ is the proton: That B flux quantization(C3) implies a big proton mass implying 2 high speed $\gamma=917$ positrons and so the Fitzgerald **contracted E field lines are the strong force**: we finally understand the strong force! So these two positrons then have big mass *two body* motion(partII) so **ortho(s,c,b) and para(t) excited (multiplet) states** understood. Eq.12 implies *Composite e,v* at $r=r_H$ is **the electroweak SM** (appendixA) **Special relativity** is that Minkowski result. **With the Eqs.16 ψ** (appendix C)

we finally understand Quantum Mechanics for the first time and eq.4 gave us a first principles derivation of r, t space-time for the first time. That Newpde $\kappa_{\mu\nu}$ metric (In eq.14), on the $N=-1$ next smaller fractal scale(1) so $r_H=10^{-40}2e^2/m_e c^2 \equiv 2Gm_e/c^2$, is the Schwarzschild metric since $\kappa_{00}=1-r_H/r=1/\kappa_{rr}$ (15): we just derived General Relativity(gravity) from quantum mechanics in one line. The Newpde zitterbewegung expansion component ($r < r_c$) on the next larger fractal scale ($N=1$) is the universe expansion: we just derived the expansion of the universe in one line. The Newpde appendix C derivation of those precision QED values (eg.,Lamb shift sect.D) allow us to abolish the renormalization and infinities. So there is no need for those many Lagrangian density postulates anymore, just postulate 1 instead.

Real# Mathematics from Postulate 1

The postulate 1 also gives the *list-define* math (B2) *list* cases $1 \cup 1 \equiv 1+1 \equiv 2$, *define* $a=b+c$ (So no other math axioms but 1.) and Cauchy sequence proof (2) of real number eigenvalues (sect.2.1,Ch.2) from a Cauchy sequence of rational numbers as a special case of the Mandelbrot set iteration formula starting $-1/4$. That means the **mathematics and the physics** come from (postulate 1 \rightarrow Newpde): *everything*. Recall from eq.7 that $dr+dt=ds$. So combining in quadrature eqs 7&11 $SNR\delta z=(dr/ds+dt/ds)\delta z=((dr+dt)/ds)\delta z=(1)\delta z$ (11a,append) and so having come *full circle* back to postulate 1 as a real eigenvalue ($1 \equiv$ Newpde electron). So we really do have a binary physics signal. So, having come *full circle* then: (postulate 1 \Leftrightarrow Newpde)

Mathematical Notion (of postulate 1 \Leftrightarrow Newpde)

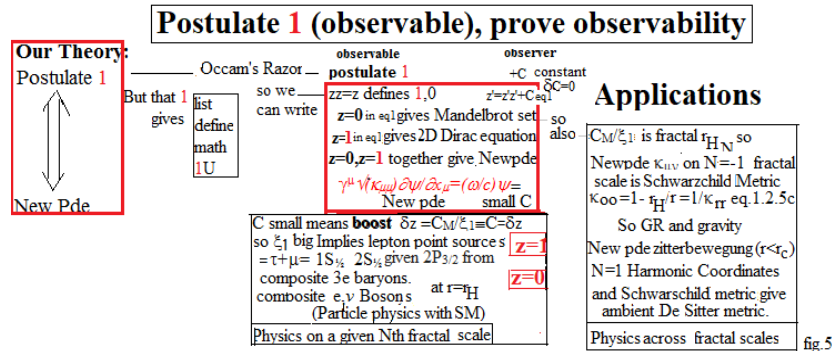


fig.5

Intuitive Notion (of postulate $1 \leftrightarrow \text{Newpde}$)

The Mandelbrot set introduces that $r_H = C_M / \xi_1$ horizon in $\kappa_{00} = 1 - r_H/r$ in the Newpde, where C_M is fractal by $10^{40} \times$ scale change (fig.2). So we have found (davidmaker.com) that: Given that fractal selfsimilarity astronomers are observing from the inside of what particle physicists are studying from the outside, that **ONE** *New pde* e electron r_H , **one** thing (fig.1). *Everything* we observe big (cosmological) and small (subatomic) is then that (*New pde*) r_H , even baryons are composite $3e$. So we understand, *everything*. This is the only Occam's razor first principles theory

Summary: So instead of doing the usual powers of 10 simulation we do a single power of 10^{40} simulation and we are immediately back to where we started!

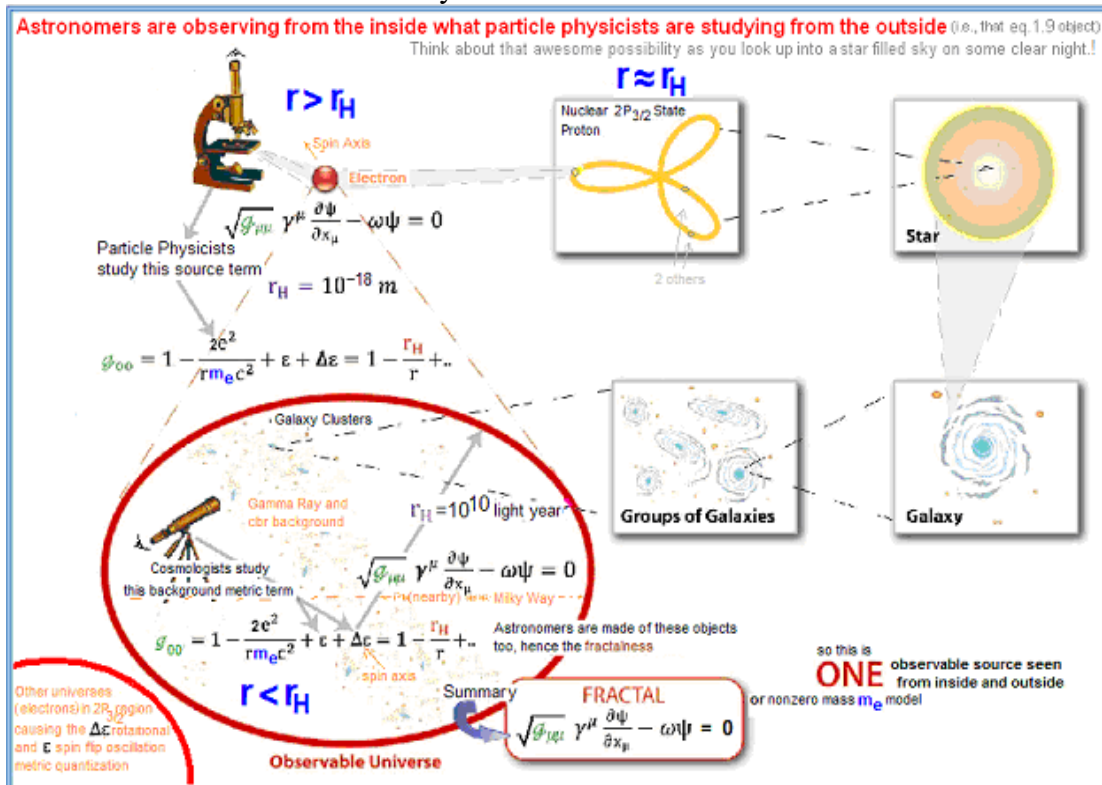


fig2

(↑lowest left corner) Object B caused perturbation structure jumps: void→galaxy→globular,,etc.

References

- (6) Penrose in a utube video implied that the Mandelbrot set might contain physics. Here we merely showed how to find it. For example the (fractal Mandelbulb neighborhood area $|drdt| > 0$ of the) Fiegenbaum point is a subset (containing that 10^{40} X selfsimiilar scale jump: Fig1)
- (7) Cantor: Ueber die Ausdehnung eines Satzes aus der Theorie der trigonometrischen Reihen, “Ueber eine elementare Frage der Mannigfaltigkeitslehre” Jahresbericht der Deutschen Mathematiker-Vereinigung.” Mandelbrot set iteration sequence z_n $C_M = -1/4$, $z_0 = 0$ same as Cauchy seq. since it begins with rational number $-1/4$, allowing the (C’ uncertainty) dr neighborhood of 0 to have a nonzero probability of being a real number and thereby giving real eigenvalues to the equation 11 operator formalism after the small C’ boost to get observability around $dr=0$. $dr=0$.
- (8) Tensor Analysis, Sokolnikoff, John Wiley
- (9) The Principle of Relativity, A Einstein, Dover
- (10) Quantum Mechanics, Merzbacher, John Wiley
- (11) lemniscate circle sequence (Wolfram, Weisstein, Eric
- (12) **Bosons: Equation 12 implications for Rotation**

Recall from sect.1 eq.3 that $\delta C = \delta(\delta z + \delta z \delta z) = \delta \delta \delta z(1) + \delta \delta z(\delta z) + (\delta z)\delta \delta z = \delta C = 0$ so C is split between $\delta \delta z$ noise and $\delta z \delta z$ classical invariance ds^2 proper time.

Recall at $N=0$ the $N=1$ $|\delta z| \gg 1$ & $C_M \gg 1$. So $\delta z \delta z \approx C_M$ there. So equation 5 holds then. But $\frac{\delta z'}{ds} = \pm 45^\circ$ ($\pi/4$) extremum at $N=0$ (SM) and $N=-1$ (GR) is also a solution for observer $N=1$.

So as the γ boosted observer $\delta z/\gamma$, C/γ , gets smaller than the huge $N=1$ scale (so higher energy, smaller wavelength, beam probes) $\delta \delta z(1)/ds$ noise angle gets relatively larger (relative to $\delta(\delta z \delta z)/ds$, sect.1) until finally the next smaller $N=0$ (and next smaller one after that, $N=-1$) is $N=0$ fractal scale in that sect.1 big angle $\pm 45^\circ$ required extremum solution (Recall ‘extremum’s are our solutions.) $45^\circ = \pi/4 \approx 1 \approx \delta z'/ds(\text{observable}) = C_{Mend}/ds \equiv \theta$ (in equation 12). So here all four $\theta \pm 45^\circ \times 2$ rotations of **Composite e,v** implied by eq.12. So we have the $N=0$, $N=-1$ solutions for $\delta z'$ angle perturbation of $N=1$ for big scattering energies. So observer $\gamma = \text{observed } \gamma$

I \rightarrow II, II \rightarrow III, III \rightarrow IV, IV \rightarrow I rotations in eq.7-9 plane Give SM Bosons

For $z=0$ $\delta z'$ is big in $z' = 1 + \delta z$ and so we have again $\pm 45^\circ$ min ds and so two possible 45° rotations so through a total of two quadrants for $\pm \delta z'$ in eq.12. one such rotation around an around a axis (SM) and the other around a diagonal (SC). Note in fig.3 dr, dt is also a rotation. and so has an eq.11 rotation operator observable θ . Thus from equation 11 for (θ) angle rotations $\theta \delta z \equiv (dr/ds)\delta z = -i\partial(\delta z)/\partial r$ for the first 45° rotation. So we got through one Newpde derivative for each 45° rotation. For the next 45° rotation in fig.4 it is then a second derivative $\theta \theta \delta z' = e^{i\theta p} e^{i\theta} \delta z = e^{i(\theta p + \theta)} \delta z = (dr/ds)((dr/ds)dr') = -i\partial(-i\partial(dr'))/\partial r \partial r = -\partial^2(dr')/\partial r^2$ large angle rotation in figure 3. In contrast for $z=1$, $\delta z'$ small so 45° - 45° small angle rotation in figure 3 (so then $N=-1$). Do the same with the time t and get for $z=0$ rotation of $45^\circ + 45^\circ$ (fig.4) then $\theta \theta \delta z' = (d^2/dr^2)z' + (d^2/dt^2)\delta z'$ (A1)

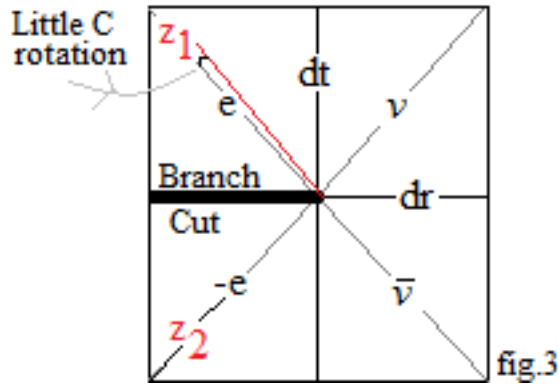


fig.3. for $45^\circ-45^\circ$ So two body (e,v) singlet $\Delta S = \frac{1}{2} - \frac{1}{2} = 0$ component so pairing interaction (sect.4.5). Also ortho $\Delta S = \frac{1}{2} + \frac{1}{2} = 1$ making 2 body (at $r=r_H$) $S=1$ Bosons and so a field theory. Note we also get these Laplacians characteristic of the Boson field equations by those $45^\circ+45^\circ$ rotations so eq.16 implies Bosons accompany our leptons (given the $\delta z'$), **so these leptons exhibit “force”**.

Newpde $r=r_H, z=0$, $45^\circ+45^\circ$ rotation of composites e,v implied by Equation 12

So $z=0$ allows a large C z rotation application from the 4 different axis' max extremum (of eq.16) branch cuts gives the 4 results: Z,+W, photon bosons of the Standard Model. So we have derived the Standard Model of particle physics in this very elegant way (from the four axis'). You are physically at $r=r_H$ if you rotate through the electron quadrants (I, IV). of eq.7-9. So we have large C_M dichotomic 90° rotation to the next Reimann surface of eq.12, eq.A1 $(dr^2+dt^2)z''$ from some initial extremum angle(s) θ . Eq.12 solutions imply complex 2D plane Stern Gerlach dichotomic rotations using eq.A1 thereby using Pauli matrices σ_i algebra, which maps one-to-one to the quaternionA algebra. Using eq.12 we start at some initial angle θ and rotate by 90° the noise rotations are: $C=\delta z'' = [e_L, v_L]^T \equiv \delta z'(\uparrow) + \delta z'(\downarrow) \equiv \psi(\uparrow) + \psi(\downarrow)$ has a eq.12 infinitesimal unitary generator $\delta z'' \equiv U = 1 - (i/2)\epsilon n \cdot \sigma$, $n \equiv \theta/\epsilon$ in $ds^2 = U^\dagger U$. But in the limit $n \rightarrow \infty$ we find, using elementary calculus, the result $\exp(-(i/2)\theta \cdot \sigma) = \delta z''$. We can use any axis as a branch cut since all 4 are eq.16 large extremum so for the 2nd rotation we move the branch cut 90° and measure the angle off the next diagonal since Pauli matrix dichotomic rotations are actually axis rotations, leaving our e and v directions the same. In any case $(dr+dt)z''$ in eq.16 can then be replaced by eq.A1 $(dr^2+dt^2+..) \delta z'' = (dr^2+dt^2+..) e^{\text{quaternionA}}$ Bosons because of eq.A1. A2 Then use eq. 12 and quaternions to rotate $\delta z''$ since the quaternion formulation is isomorphic to the Pauli matrices. $dr' = \delta z_r = \kappa_r dr$ for **Quaternion A** $\kappa_{ij} = e^{iA_i}$.

Appendix A Quaternion ansatz $\kappa_r = e^{iA_r}$ instead of $\kappa_r = (dr/dr')^2$ in eq.14. $N=0$.

A1 for the eq.12: large $\theta = 45^\circ+45^\circ$ rotation (for $N=0$ so large $\delta z' = \theta r_H$). Instead of the equation 13,15 formulation of κ_{ij} for small $\delta z'$ ($z=1$) and large $\theta = 45^\circ+45^\circ$ we use A_r in dr direction with $dr^2 = x^2 + dy^2 + dz^2$. So we can again use 2D (dr,dt) $E = 1/\sqrt{\kappa_{00}} = 1/\sqrt{e^{iA_i}} = e^{-iA/2}$. The 1 is mass energy

and the first real component after that in the Taylor expansion is field energy A^2

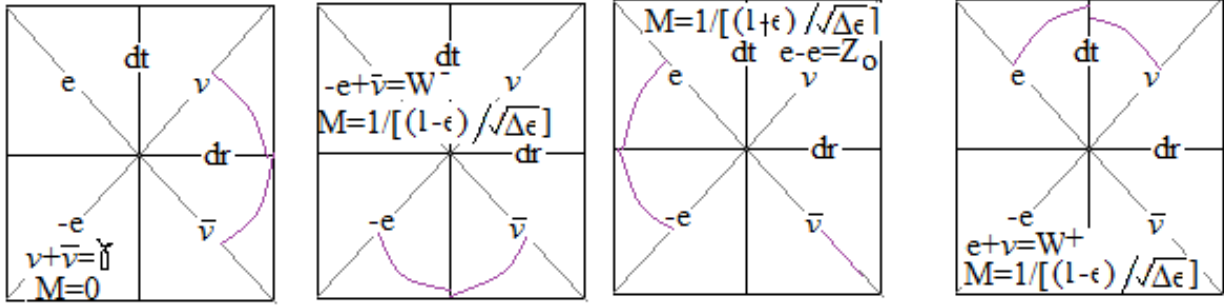


fig4

Fig.4 applies to eq.9 $45^\circ+45^\circ=90^\circ$ case: **Bosons**.

A2 These quadrants were defined in eq.7-9 and used in eq.12. The Appendix A4 derivation applies to the far right side figure. Recall from eq.12 $z=0$ result $C_M=45^\circ+45^\circ=90^\circ$, gets Bosons. $45^\circ-45^\circ$ leptons. The v in quadrants II (eq.5) and III (eq.9). e in quadrants I (eq.7) and IV (eq.7). Locally normalize out $1+\epsilon$ (appendix D). For the **composite** e, v on those required large $z=0$ eq.9 rotations for $C \rightarrow 0$, and for stability $r=r_H$ (eg., for $2P_{1/2}$, $I \rightarrow II$, $III \rightarrow IV$, $IV \rightarrow I$) unless $r_H=0$ ($II \rightarrow III$) Example:

A4 Quadrants II \rightarrow III rotation eq.A2 $(dr^2+dt^2+..)e^{\text{quaternion } A} = \text{rotated through } C_M$ in eq.16.

example C_M in eq.A1 is a 90° CCW rotation from 45° through v and antiv

A is the 4 potential. From eq.9b we find after taking logs of both sides that $A_0=1/A_r$ (A2)

Pretending we have a only two i, j quaternions but still use the quaternion rules we first do the r derivative: From eq. A1 $dr^2\delta z = (\partial^2/\partial r^2)(\exp(iA_r+jA_0)) = (\partial/\partial r)[(i\partial A_r/\partial r + \partial A_0/\partial r)(\exp(iA_r+jA_0))] = \partial/\partial r[(\partial/\partial r)iA_r + (\partial/\partial r)jA_0](\exp(iA_r+jA_0)) + [i\partial A_r/\partial r + j\partial A_0/\partial r]\partial/\partial r(\exp(iA_r+jA_0)) + (i\partial^2 A_r/\partial r^2 + j\partial^2 A_0/\partial r^2)(\exp(iA_r+jA_0)) + [i\partial A_r/\partial r + j\partial A_0/\partial r][i\partial A_r/\partial r + j\partial A_0/\partial r]\exp(iA_r+jA_0)$ (A3)

Then do the time derivative second derivative $\partial^2/\partial t^2(\exp(iA_r+jA_0)) = (\partial/\partial t)[(i\partial A_r/\partial t + \partial A_0/\partial t)(\exp(iA_r+jA_0))] = \partial/\partial t[(\partial/\partial t)iA_r + (\partial/\partial t)jA_0](\exp(iA_r+jA_0)) + [i\partial A_r/\partial t + j\partial A_0/\partial t]\partial/\partial t(\exp(iA_r+jA_0)) + (i\partial^2 A_r/\partial t^2 + j\partial^2 A_0/\partial t^2)(\exp(iA_r+jA_0)) + [i\partial A_r/\partial t + j\partial A_0/\partial t][i\partial A_r/\partial t + j\partial A_0/\partial t]\exp(iA_r+jA_0)$ (A4)

Adding eq. A2 to eq. A4 to obtain the total D'Alambertian $A_3+A_4=$

$$[i\partial^2 A_r/\partial r^2 + i\partial^2 A_r/\partial t^2] + [j\partial^2 A_0/\partial r^2 + j\partial^2 A_0/\partial t^2] + ii(\partial A_r/\partial r)^2 + ij(\partial A_r/\partial r)(\partial A_0/\partial r) + ji(\partial A_0/\partial r)(\partial A_r/\partial r) + jj(\partial A_0/\partial r)^2 + ii(\partial A_r/\partial t)^2 + ij(\partial A_r/\partial t)(\partial A_0/\partial t) + ji(\partial A_0/\partial t)(\partial A_r/\partial t) + jj(\partial A_0/\partial t)^2 .$$

Since $ii=-1$, $jj=-1$, $ij=-ji$ the middle terms cancel leaving $[i\partial^2 A_r/\partial r^2 + i\partial^2 A_r/\partial t^2] + [j\partial^2 A_0/\partial r^2 + j\partial^2 A_0/\partial t^2] + ii(\partial A_r/\partial r)^2 + jj(\partial A_0/\partial r)^2 + ii(\partial A_r/\partial t)^2 + jj(\partial A_0/\partial t)^2$

Plugging in A2 and A4 gives us cross terms $jj(\partial A_0/\partial r)^2 + ii(\partial A_r/\partial t)^2 = jj(\partial(-A_r/\partial r))^2 + ii(\partial A_r/\partial t)^2 = 0$. So $jj(\partial A_r/\partial r)^2 = -jj(\partial A_0/\partial t)^2$ or taking the square root: $\partial A_r/\partial r + \partial A_0/\partial t = 0$ (A5)

$i[\partial^2 A_r/\partial r^2 + i\partial^2 A_r/\partial t^2] = 0$, $j[\partial^2 A_0/\partial r^2 + i\partial^2 A_0/\partial t^2] = 0$ or $\partial^2 A_\mu/\partial r^2 + \partial^2 A_\mu/\partial t^2 + .. = 1$ (A6)

A4 and A5 are **Maxwell's equations** (Lorentz gauge formulation) in free space, if $\mu=1,2,3,4$.

$$\square^2 A_\mu = 1, \quad \square \bullet A_\mu = 0 \quad (A7)$$

The Lorentz gauge is the only gauge hence it is no gauge at all and we have avoided the Maxwell overdeterminism problem (8eq., 6 unknowns E_i, B_i). Must use Newpde 4D orthogonalization here Amplitudes of physical processes in QED in the noncovariant Coulomb gauge coincide with those in the covariant Lorenz gauge. The Aharonov–Bohm effect depends on a line integral of A around a closed loop, and this integral is not changed by $A \rightarrow A + \nabla\psi$ which doesn't change

$B=\nabla XA$ either. So formulation in the Lorentz gauge mathematics works so it is no longer a gauge, we are gaugeless.

A5 Other 45°+45° Rotations (Besides above quadrants II→III)

For the **composite e,v** on those required large $z=0$ eq.12 rotations for $C\approx 0$, and for stability $r=r_H$ for $2P_{1/2}$ (I→II, III→IV, IV→I) unless $r_H=0$ (II→III) are:

Ist→IIrd quadrant rotation is the W^+ at $r=r_H$. Do similar math to A2-A7 math and get instead a Proca equation The limit $\varepsilon \rightarrow 1=\tau$ (D13) in ξ_1 at $r=r_H$. since Hund's rule implies $\mu=\varepsilon=1S_{1/2} \leq 2S_{1/2}=\tau=1$. So the ε is negative in $\Delta\varepsilon/(1-\varepsilon)$ as in case 1 charged as in appendix C1 case 2.

$E=1/\sqrt{(\kappa_{00})} -1 = [1/\sqrt{(1-\Delta\varepsilon/(1-\varepsilon))-r_H/r}] -1 = [1/\sqrt{(\Delta\varepsilon/(1-\varepsilon))}] -1$. $E_t=E+E=2/\sqrt{(\Delta\varepsilon/(1-\varepsilon))}=W^+$ mass.

$E_t=E-E$ gives E&M that also interacts weakly with weak force.

IIIrd →IV quadrant rotation is the W^- . Do the math and get a Proca equation again.

$E=1/\sqrt{(\kappa_{00})} -1 = [1/\sqrt{(1-\Delta\varepsilon/(1-\varepsilon))-r_H/r}] -1 = [1/\sqrt{(\Delta\varepsilon/(1-\varepsilon))}] -1$. $E_t=E+E=2/\sqrt{(\Delta\varepsilon/(1-\varepsilon))}=W^-$ mass.

$E_t=E-E$ gives E&M that also interacts weakly with weak force.

IVth → Ist quadrant rotation is the Z_0 . Do the math and get a Proca equation. C_M charge cancelation. D14 gives $1/(1+\varepsilon)$ gives 0 charge since $\varepsilon \rightarrow 1$ to case 1 in appendix C2.

$E=1/\sqrt{(\kappa_{00})} -1 = [1/\sqrt{(1-\Delta\varepsilon/(1+\varepsilon))-r_H/r}] -1 = [1/\sqrt{(\Delta\varepsilon/(1+\varepsilon))}] -1$. $E_t=E+E=2/\sqrt{(\Delta\varepsilon/(1+\varepsilon))}-1=Z_0$ mass.

$E_t=E-E$ gives E&M that also interacts weakly with weak force. Seen in small left handed polarization rotation of light.

IInd→IIIrd quadrant rotation through those 2 neutrinos gives 2 objects. $r_H=0$

$E=1/\sqrt{\kappa_{00}} -1 = [1/\sqrt{(1-\Delta\varepsilon/(1+\varepsilon))}] -1 = \Delta\varepsilon/(1+\varepsilon)$. Because of the \pm square root $E=E+-E$ so E rest mass is 0 or $\Delta\varepsilon=(2\Delta\varepsilon)/2$ reduced mass.

$E_t=E+E=2E=2\Delta\varepsilon$ is the pairing interaction of SC. The $E_t=E-E=0$ is the 0 rest mass photon

Boson. Do the math (eq.A2-A7) and get Maxwell's equations. Note there was no charge C_M on the two ν s. Note we get SM particles out of composite e,v using required eq.9 rotations for

A6 Object B Effect On Inertial Frame Dragging (from appendix D)

The fractal implications are that we are inside a cosmological positron inside a proton $2P_{3/2}$ at $r=r_H$ state. The cosmological object (electron) we are inside of is a positron and call it object A which orbits electron object B with a given distant 3rd object C. Object B is responsible for the mass of the electron since it's frame dragging creates that Kerr metric $(a/r)^2=m_e c^2$ (D9) result used in eq.D9. So Newpde ground state $m_e c^2 \equiv \langle H_e \rangle$ is the fundamental Hamiltonian eigenvalue defining idea for composite e,v, $r=r_H$ implying Fermi 4 point $E = \int \psi^\dagger H \psi dV = \int \psi^\dagger \psi H dV = \int \psi^\dagger \psi G$

Recall for composite e,v all interactions occur inside r_H $(4\pi/3)\lambda^3 = V_{rH}$. $\frac{1}{V^{1/2}} = \psi_e = \psi_3 \frac{1}{V^{1/2}} =$

$\psi_v = \psi_4$ so 4pt $\iiint_0^{r_H} \psi_1 \psi_2 \psi_3 \psi_4 dV = 2G \iiint_0^{r_H} \psi_1 \psi_2 \frac{1}{V^{1/2}} \frac{1}{V^{1/2}} V$

$\equiv \iiint_0^{r_H} \psi_1 \psi_2 G \equiv \iiint_0^{r_H} \psi_1 \psi_2 (2m_e c^2) dV_{rH} = \iiint_0^{V_{rH}} \psi_1 (2m_e c^2) \psi_2 dV_{rH}$ (A8)

Application of Eq.A8 To Ortho states (created by that $2P_{3/2}$ 2body motion at $r=r_H$)

The composite 3e ortho state (partII) operator adds spin (eg., as in 2nd derivative eq.A1) so 2nd derivative $\Sigma((\gamma^\mu \sqrt{\kappa_{\mu\mu}} dx_\mu) - i\kappa)(\gamma^\nu \sqrt{\kappa_{\nu\nu}} dx_\nu + i\kappa)\chi = \Sigma((\gamma^\mu \sqrt{\kappa_{\mu\mu}} dx_\mu) - i\kappa)\psi$ so $\frac{1}{2}(1 \pm \gamma^5)\psi = \chi$. In that regard the expectation value of γ^5 is speed and varies with $e^{i3\phi/2}$ in the trifoldium. The spin $^{1/2}$ decay proton $S_{1/2} \propto e^{i\phi/2} \equiv \psi_1$, the original ortho $2P_{1/2}$ particle is chiral $\chi = \psi_2 \equiv \frac{1}{2}(1 - \gamma^5)\psi = \frac{1}{2}(1 - \gamma^5 e^{i3\phi/2})\psi$. Initial

$$\begin{aligned} \text{A8 to read} &= \iiint_0^{V_{rH}} \psi_1 \psi_2 (2m_e c^2) dV_{rH} = \iint \psi_{S1/2}^* (2m_e c^2 V_{rH}) \chi dV_\phi = K \int \langle e^{i\frac{\phi}{2}} [\Delta \varepsilon V_{rH}] (1 - \\ &\gamma^5 e^{i\frac{\phi}{2}}) \psi \rangle d\phi = K G_F \int \langle e^{i\phi/2} - \gamma^5 i e^{i(4/2)\phi} \rangle d\phi = K G_F \left(\frac{2e^{i\phi}}{i} \Big|_0^{2\pi+C} - \frac{2\gamma^5 e^{i4\phi}}{i4} \Big|_0^{2\pi+C} \right) = \end{aligned}$$

Review of 2P_{3/2} Next higher fractal scale (X10⁴⁰), cosmological scale. Recall from D9 $m_e c^2 = \Delta \epsilon$ is the energy gap for object B vibrational stable eigenstates of composite 3e (vibrational perturbation r is the only variable in Frobenius solution, part II Ch.8,9,10) proton. Observer in object A. From fig.7 $v \sin 30^\circ + v \sin 30^\circ = v$. From fig 7 $r^2 = 1^2 + 1^2 + 2(1)(1)\cos 120^\circ = 3$, so $r = \sqrt{3}$.

Recall for the positron motion $\gamma = \frac{1}{\sqrt{1-\frac{v^2}{c^2}}} = 917$. So Fitzgerald contract $r_{CA} = \sqrt{1 - \frac{\cos^2 30^\circ c^2}{c^2}} \sqrt{3} =$

.866=cos30°. The E field in the forward or backward direction of the CA line (the weakest) due to a charge moving away is $E'=(1-v^2/c^2)E=(1/\gamma^2)E=(1/917^2)E$ (from Feynman's lectures) where $E=q/r^2$. For circular motion in the proton around the central electron $\frac{mv^2}{r} = qE$ **so that** $\Delta mc^2 = KE = \frac{mv^2}{2} = qr_{CA}E_{CA} \frac{1}{2} = ((1/\gamma^2)q^2/r_{CA}^2)(r_{CA}))^{1/2} = (1/\gamma^2)q(q/1^2))(1)^{1/2}/r_{CA} = ((1/\gamma^2)/r_{CA})[(qE_{AB}r_{AB})^{1/2}] = ((1/\gamma^2)/r_{CA})(m_e c^2) = \Delta mc^2$ in summary = object C scissors eigenstates.

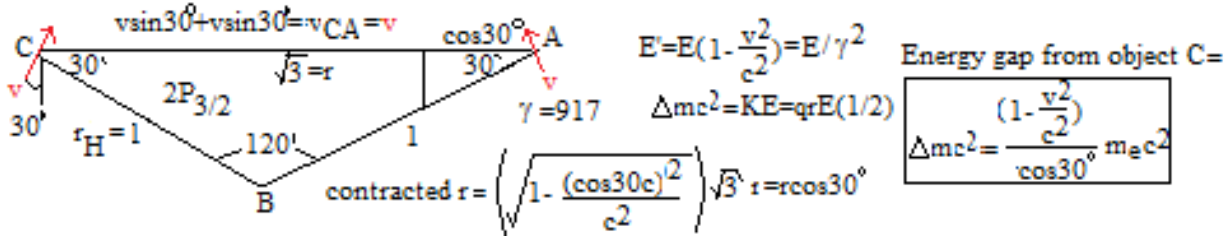


Fig7 Allowing us to finally compare the energy gap caused by object C to the energy gap caused by object B (A8). So to summarize $E_{qr} = \Delta E = (m_e c^2 / ((\cos 30^\circ) 917^2)) = m_e c^2 / 728000$. So energy gap caused by object C is $\Delta E = (m_e c^2 / ((\cos 30^\circ) 917^2)) = m_e c^2 / 728000$. The weak interaction occurs inside of r_H with those electrons m_e . The G can be written for E&M decay as $(2m_e c^2) X V_{r_H} = 2m_e c^2 [(4/3) \pi r_H^3]$. So for weak decay from equation A8 it is $G_F = (2m_e c^2 / 728,000) V_{r_H} = G_F$ **the strength of the Fermi weak interaction constant** which is the coupling constant for the Fermi 4 point weak interaction integral. Note $2m_e c^2 / 729,000 = 1.19 \times 10^{-19} \text{J}$. So $\Delta E = 1.19 \times 10^{-19} / 1.6 \times 10^{-19} = .7 \text{eV}$ which is our ΔE gap for the weak interaction inside the integral for G_F .

A8 Derivation of the Standard Model from Newpde but with No Free parameters

Since we have now derived M_W , M_Z and their associated Proca equations, and Dirac equations for m_τ, m_μ, m_e etc., and G, G_F, ke^2 Maxwell's equations, etc. we can now write down the usual Lagrangian densities that implies these results. In the formulation $M_Z=M_W/\cos\theta_W$ you can find the Weinberg angle θ_W , $g\sin\theta_W=e$, $g'\cos\theta_W=e$; solve for g and g' , etc., We will have thereby

derived the standard model from first principles (i.e., postulate 1). It no longer contains free parameters.

Appendix B ultimate Occam's razor (observable) also implies the underlying relational algebra (without all those many Relational axioms). Eg., $1 \cup 1 \equiv 1 + 1$ (Ch.2).

Postulate 1 (observable) so observer C so $1 \cup C \equiv 1 + C$. with algebraic definition of $1 \cdot z = zz$ having both 1,0 as solutions so defining negation \sim with $0 = 1 - 1$. Thus we can define $\sim((A \cup B) \sim B \sim A) \equiv A \cap B$. So we have defined intersection \cap so we have derived set theory. So in postulate 1 $z = zz$ why did 0 come along for the ride? There is a deeper reason in set theory. Note \emptyset and 0 aren't really new postulates since they postulate literally "nothing".

Recall we just derived set theory from the postulate of 1 (observable).

The null set \emptyset is the subset of every set. In the more fundamental set theory formulation

$\{\emptyset\} \subset \{\text{all sets}\} \Leftrightarrow \{0\} \subset \{1\}$ since $\emptyset = \emptyset \cup \emptyset \Leftrightarrow 0 = 0 + 0$, $\{\{1\} \cup \emptyset\} = \{1\} \Leftrightarrow 1 + 0 = 1$.

So list $1 \cup 1 \equiv 1 + 1 \equiv 2$, $2 \cup 1 \equiv 1 + 2 \equiv 3$, ... all the way up to 10^{82} (see Fiegenbaum point) and **define** all this list as $a + b = c$, etc., to create our algebra and numbers which we use to write equation 1 $z = zz + C$, $\delta C = 0$ for example. Recall every set has the null set as a subset. So from above set $\{1\}$ (ξ_1 for $z = 1$) has the 0 (ξ_0 for $z = 0$ ground state) as a subset. So $\xi_1 = \xi_{2S\frac{1}{2}} + \xi_{1S\frac{1}{2}} + \xi_0 = \tau + \mu + m_e$. (B1)

B2 2D+2D→4D Orthogonality

Note adding the $N=0$ fractal scale 2D δz perturbation to $N=1$ eq.7 2D gives curved space 4D. So $(dx_1 + i dx_2) + (dx_3 + i dx_4) \equiv dr + i dt$ given (eqs 5, 7a) $dr^2 - dt^2 = (\gamma^r dr + i \gamma^t dt)^2$ if $dr^2 \equiv dx^2 + dy^2 + dz^2$ so that $\gamma^r dr \equiv \gamma^x dx + \gamma^y dy + \gamma^z dz$, $\gamma^i \gamma^j + \gamma^j \gamma^i = 0$, $i \neq j$, $(\gamma^i)^2 = 1$, rewritten (with curved space $\kappa_{\mu\nu}$ eq.13-15) $(\gamma^x \sqrt{\kappa_{xx}} dx + \gamma^y \sqrt{\kappa_{yy}} dy + \gamma^z \sqrt{\kappa_{zz}} dz + \gamma^t \sqrt{\kappa_{tt}} dt)^2 = \kappa_{xx} dx^2 + \kappa_{yy} dy^2 + \kappa_{zz} dz^2 - \kappa_{tt} dt^2 = ds^2$.

More fundamentally satisfying this 4D Clifford algebra and complex orthogonalization requirement is a special case of any 2 $x_i x_j$ in eq.3 (directly from postulate 1): Imposing orthogonality thereby creates 6 pairs of eqs.3&5. So each particle carries around it's own $dr + i dt$ complex coordinates with them on their world lines. Alternatively this 2D $dr + i dt$ is a 'hologram' 'illuminated' by a modulated $dr^2 + dt^2 = ds^2$ 'circle' wave (as 2nd derivative wave equation operators from eq.11 circle) since 4Degrees of freedom are imbedded on a 2D (dr, dt) surface here, with observed coherent superposition output as eq.16 solutions. A more direct way is to simply write the 4Degrees of freedom on the 2D surface as $dr + i dt = (dr_1 + i dt_1) + (dr_2 + i dt_2) = (dr_1, \omega dt_2), (dr_2, i dt_2) = (x, z, y, i dt) = (x, y, z, i dt)$, where $\omega dt \equiv dz$ is the z direction spin $\frac{1}{2}$ component ω (angular velocity) axial vector of the Newpde lepton (eqs.7-9); which we get anyway from lepton equation eq.16.

Appendix C

Quantum Mechanics Is The Newpde ψ (for each N fractal scale)

Recall the solution to (postulate 1) $z = zz$ is 1,0. In $z = 1 - \delta z$, $\delta z^* \delta z$ is (defined as) the probability of z being 0. Recall $z = 0$ is the $\xi_0 = m_e$ solution to the new pde so $\delta z^* \delta z$ is the probability we have just an electron (sect.3). Note $z = zz$ also thereby conveniently provides us with an automatic normalization of δz . Note also that $(\delta z^* \delta z)/dr$ is also then a one dimensional probability 'density'. So Bohr's probability density "postulate" for $\psi^* \psi \equiv (\delta z^* \delta z)$ is derived here. It is not a postulate anymore. (So Bohr was very close to the postulate of 1, and so using $z = zz$ here.)

Note the electron eq.7 has two parts (i.e., $dr + dt$ & $dr - dt$), that solve eq.3 together, same kind of $\delta(p_A - p_B)$ conservation relation as between Alice and Bob; signal, idler, Bell's stuff. We could then label these two parts *observer* and *object* with associated eq.7 wavefunctions $\delta z \equiv \psi_1$,

$\delta z \equiv \psi_2$. So if there is no observer eq.7 (So no ψ_1) then eq.3 doesn't hold at all and so there is no object "observed" wavefunction. ψ_2 . Thus the object wave function ψ_2 "collapses" to the wavefunction 'observed' ψ_2 (or eq.5 and so **postulate 1** does not even hold), if "observed" ψ_1 exists. Then apply the same mathematical reasoning to every other $\delta(p_A - p_B)$ situation and we will also have thereby derived Bell's theorem and its general cases. Thus we derived the Copenhagen interpretation of Quantum Mechanics QM mathematically, from eq.16 and so first principles **postulate 1**, not from the usual hand waving arguments.

Recall from appendix A $dr^2 + dt^2$ is a second derivative *operator* wave equation(A1), that holds all the way around the circle(even for the eq.10 vacuum solutions), gives waves. In eq.12, error magnitude C (sect.2.3) is also a δz angle measure on the dr, idt plane. One extremum ds ($z=0$) is at 45° so the largest C is on the diagonals (45°) where we have eq.4 extremum holding: particles. So a wide slit has high uncertainty, so large C (rotation angle) so we are at 45° (eg., particles, eq.16 photoelectric effect). For a *small slit* we have less uncertainty so smaller C, not large enough for 45° , so only the *wave equation* A1 holds (small slit diffraction). Thus we derived wave particle duality here.

Recall wave equation eq.A1 iteration of the New pde with eq.11 operator formalism. So $dr/ds = k$ in the sect.1 $\delta z = ds e^{i\theta}$ θ exponent then becomes $k = 2\pi/\lambda$. Multiplying both sides by \hbar with $\hbar k \equiv mv$ as before we then have the DeBroglie equation that relates particle momentum to wavelength in quantum mechanics. Equation 8a (sect.1) then counts units N of $(dt/ds) = \hbar\omega = \hbar ck$ on the diagonal so that $E = p v = \hbar\omega$ for all energy components, universally. Thus this eq.11a counting N does not require the (well known) quantization of the E&M field with SHM. First, set the unit of distance r_H on our baseline fractal scale: (eq.1 $N=0$. See figure 1 attachment.). The 4X Mandelbrot set formulation allows only these finite extremum.

$\delta z \equiv \psi$

Appendix D. $N=0$ (eq.13,14,15 give our Newpde metric $\kappa_{\mu\nu}$ at $r < r_H$, $r > r_H$)

Found GR from eq.13 and eq.14 so we can now write the Ricci tensor $R_{\mu\nu}$ (and self similar perturbation Kerr metric since frame dragging decreased by external object B, sect.A6). Also for fractal scale $N=0$, $r_H = 2e^2/m_e c^2$, and for $N=-1$ $r'_H = 2Gm_e/c^2 = 10^{-40} r_H$.

Nonzero Generic maximally symmetric (MS) ambient metric generated by object B

$N=2$ big guy sees us from the outside and so sees sine. To see what we see he multiplies \sin by i and u by I since we are inside (so $-i \sin u \rightarrow \sin hu$). So start with complete frame dragging suppression eq.13, 15 but with ambient metric (provided by later **perturbation** $a < r$ **provided by some rotation**) metric ansatz: $ds^2 = -e^\lambda (dr)^2 - r^2 d\theta^2 - r^2 \sin\theta d\phi^2 + e^\mu dt^2$ so that $g_{00} = e^\mu$, $g_{rr} = e^\lambda$. From eq. $R_{ij} = 0$ for spherical symmetry in free space

$$R_{11} = \frac{1}{2}\mu'' - \frac{1}{4}\lambda' \mu' + \frac{1}{4}(\mu')^2 - \lambda'/r = 0 \quad (D1)$$

$$R_{22} = e^{-\lambda} [1 + \frac{1}{2} r(\mu' - \lambda')] - 1 = 0 \quad (D2)$$

$$R_{33} = \sin^2\theta \{ e^{-\lambda} [1 + \frac{1}{2} r(\mu' - \lambda')] - 1 \} = 0 \quad (D3)$$

$$R_{00} = e^{\mu-\lambda} [-\frac{1}{2}\mu'' + \frac{1}{4}\lambda' \mu' - \frac{1}{4}(\mu')^2 - \mu'/r] = 0 \quad (D4)$$

$$R_{ij} = 0 \text{ if } i \neq j$$

(eq. D1 -D4 from pp.303 Sokolnikof(8)): Equation D2 is a mere repetition of equation D3. We thus have only three equations on λ and μ to consider. From equations D1, D4 we deduce that $\lambda' = -\mu'$ so that radial $\lambda = -\mu + \text{constant} = -\mu + C$ where C represents a possible \sim constant ambient metric contribution which could be imaginary in the case of the slowly oscillating ambient metric of nearby object B. So $e^{-\mu+C} = e^\lambda$. Then D2 can be written as:

$$e^{-C}e^{\mu}(1+r\mu')=1 \quad (D5)$$

Set $e^{\mu}=\gamma$. So $e^{-\lambda}=\gamma e^{-C}$ and so integrating this first order equation (equation.D11) we get:

$$\gamma=-2m/r+e^C \equiv e^{\mu}=g_{00} \text{ and } e^{-\lambda}=(-2m/r+e^C)e^{-C}=1/g_{rr} \quad (D6)$$

From equation D6 we can identify radial C with also rotational Kerr metric oblateness perturbation Mandlebulb component here (D8 below) of Mandelbrot set Fig.6 eq.18

$2m/r=r_H/r=C_M/\xi r=e^{-C}=e^{-(\varepsilon+\Delta\varepsilon)}=\tau+\mu+\Delta\varepsilon$. (eq.17). We end up being at the horizon r_H in equation D8. So $2m/r$ is set equal to e^C in eq. D6. So at the end, at the horizon r_H , in eq.D8, $2m/r$ is set equal to $e^C=e^{-(\varepsilon+\Delta\varepsilon)}$ in D6. So $\kappa_{00}=1-e^{-(\varepsilon+\Delta\varepsilon)}-2m/r$. from eq.17. Given external object B oscillating zitterbewegung for $r<r_C$ then $e^{-(\varepsilon+\Delta\varepsilon)} \rightarrow e^{-i(\varepsilon+\Delta\varepsilon)}$ so that $\kappa_{00}=1-e^{-i(\varepsilon+\Delta\varepsilon)}-2m/r$ (D7) So: $e^{-\lambda}=1/\kappa_{rr}=1/(1-2m'/r)$

Perturbative self similar rotation providing the ambient metric Generated by object B

Our new pde has spin S and so the self similar ambient metric on the $N=0$ th fractal scale is the Kerr metric which contains those **perturbation rotations** ($d\theta/dt$ T violation so (given CPT) then

$$\textbf{CP violation) } ds^2 = \rho^2 \left(\frac{dr^2}{\Delta} + d\theta^2 \right) + (r^2 + a^2) \sin^2 \theta d\phi^2 - c^2 dt^2 + \frac{2mr}{\rho^2} (a \sin^2 \theta d\theta - c dt)^2, \quad (D8)$$

where $\rho^2(r, \theta) \equiv r^2 + a^2 \cos^2 \theta$; $\Delta(r) \equiv r^2 - 2mr + a^2$, In our 2D $d\phi=0, d\theta=0$ Define:

$$\left(\frac{r^2 + a^2 \cos^2 \theta}{r^2 - 2mr + a^2} \right) dr^2 + \left(1 - \frac{2m}{r^2 + a^2 \cos^2 \theta} \right) dt^2 \quad \theta \neq 0$$

$\rho^2(r, \theta) \equiv r^2 + a^2 \cos^2 \theta$; $\Delta(r) \equiv r^2 - 2mr + a^2$, $r^{\wedge 2} \equiv r^2 + a^2 \cos^2 \theta$, $r'^2 \equiv r^2 + a^2$. Inside r_H $a \ll r, r' \gg 2m$

$$\left(\frac{(r^{\wedge})^2}{(r')^2 - 2mr} \right) dr^2 + \left(1 - \frac{2mr}{(r^{\wedge})^2} \right) dt^2 + \dots = \left(\frac{1}{\frac{(r')^2}{(r^{\wedge})^2} - \frac{2mr}{(r^{\wedge})^2}} \right) dr^2 + \left(1 - \frac{2mr}{(r^{\wedge})^2} \right) dt^2.$$

The $(r^{\wedge}/r')^2$ term is

$$\frac{(r')^2}{(r^{\wedge})^2} = \frac{r^2 + a^2}{r^2 + a^2 \cos^2 \theta} = \frac{1 + \frac{a^2}{r^2}}{1 + \frac{a^2}{r^2} \cos^2 \theta} \approx 1/g_{rr} (\approx g_{00}) \quad \text{From D7: } \xi_1 = e^{i(\varepsilon + \Delta\varepsilon)} \text{ for } e^C = e^{i(\varepsilon + \Delta\varepsilon)}$$

$= \tau + \mu + \Delta\varepsilon = \text{zitterbewegung from D6. } 2m/r + e^C$

$$\begin{aligned} \left(1 + \frac{a^2}{r^2} \right) \left(1 - \frac{a^2}{r^2} \cos^2 \theta \right) + \dots &= 1 - \frac{a^4}{r^4} \cos^2 \theta - \frac{a^2}{r^2} \cos^2 \theta + \frac{a^2}{r^2} + \dots = 1 + \frac{a^2}{r^2} (1 - \cos^2 \theta) + \dots \\ &= 1 + \frac{a^2}{r^2} \sin^2 \theta + \dots \equiv 1 + \frac{a^2}{r^2} u^2 = (D7,17) = 1 + e^C = 1 + e^{i(\varepsilon + \Delta\varepsilon)} = \end{aligned}$$

(Replace a^2/r^2 Kerr object B term with inertial frame D7 dragging mass ξ_1 . In eq.D8 subtract $2mr/(r')^2 = r_H/r_H$). In From eq.17 general the closer object B is the larger e^C is.

$$= 1 + \xi_1 - \frac{r_H}{r_H} = e^C = 1 + \varepsilon + \Delta\varepsilon + \dots = e^{i(\varepsilon + \Delta\varepsilon)} \quad (D9)$$

So this is a Kerr metric inertial frame dragging suppression due to outside object B of magnitude $((a/r) \sin \theta)^2 = 1/g_{rr} = e^{i\varepsilon}$ from D7 in the proper frame. Inside object A. ε also changes with time (Mercuron equation D15).

Object B oscillation sound wave observed compression in Shapely, Bootes, rarefaction in Eridanis.

D2 Examples of this ambient metric. N=0 Composite 3e

Introduction: $N=0$ Frobenius solution is for constant ψ (and so constant ε) just inside r_H .

Equations D6,D7,D9 provide the $e^{i(\varepsilon + \Delta\varepsilon)}$ contributions from each maximal symmetry ε source, with the B flux quantization causing the $n\varepsilon$ quantization of the ambient metric. There appear to

be 2 B field sources, the two fast moving positrons (are right on r_H and so are close to these boundaries) creating that huge internal magnetic field. So for the inside $1+2(\varepsilon + \Delta\varepsilon)$ get added and we normalize the maximal symmetry B field away for the observer 2nd positron by dividing by $1+\varepsilon$.

In contrast for just *outside* r_H the flux is canceled out because of the frequent creation and annihilation events inside resulting in a Faraday's law B flux change cancellation application that gives the Meisner effect zero point energy (eq.9.22) pion ε' cloud who's energy is thereby added to $2m/r=r_H/r$ as implied by eq. D6. Thus:

For $z=0$ just inside r_H , the two positrons each have constant ψ ($N=0$ ch.8,9) inside r_H . So from eq.D9 divide κ_{rr} by $1+\varepsilon+\varepsilon=1+2\varepsilon=e^C$ So $\frac{1}{\kappa_{rr}} = (1)(1 + 2\varepsilon) \equiv 1 + 2(\varepsilon + \Delta\varepsilon)$ (D9a)

Note negative potential energy here. Normalize out the $\kappa_{\theta\theta}$ magnetic field maximak symmetry of the observer by multiplying $\kappa_{\theta\theta}$ by $1+\varepsilon=e^C$ for the magnetic (see partII flux of B)

$$\begin{aligned} \frac{1}{\left(\frac{1+2\varepsilon+\Delta\varepsilon}{1+\varepsilon}-2m/\xi_0 r\right)} dr^2 + (1 - 2m/r\xi_0) dt^2 &= \frac{1}{\left(1+\frac{\varepsilon}{1+\varepsilon}-2m/\xi_0 r\right)} dr^2 + \left(1 - \frac{2m}{r\xi_0}\right) dt^2 \\ &= \frac{1}{(1+\varepsilon'-2m/\xi_0 r)} dr^2 + \left(1 - \frac{2m}{r\xi_0}\right) dt^2, \quad \varepsilon' \equiv \varepsilon/(1+\varepsilon). \end{aligned} \quad (D10)$$

For $z=0$ just outside r_H , Since randomly the B field disappears ($dB/dt \neq 0$) due to that creation-annihilation we have a Faraday's law Meisner effect. With outside r_H B results, just divide by $1+\varepsilon'$ (D9) for zero point energy $\varepsilon''=0.8 \pi^\pm$ of eq.9.22 (partII) which has to itself increase and decrease with (see D9) each of these annihilation events and π^\pm exists just outside r_H (from our Frobenius solution): $\frac{1}{(1+\varepsilon''-2m/\xi_0 r)} dr^2 + ((1 - 2m/\xi_0 r)) dt^2 = ds^2$ (D11)

For $z=0 \rightarrow z=1$ $r > r_H$ then free space boost sect.2 $\xi_0 \rightarrow \tau$. Define $\varepsilon' \equiv \frac{\varepsilon}{1+\varepsilon}$. Must normalize again (for local ambient metric $\Delta\varepsilon$ change contributions) so multiply by $\frac{1}{1+\varepsilon'}$ (see D9 for $z=1$ outside)

$$\frac{1}{\left(1+\frac{\Delta\varepsilon}{1+\varepsilon}-2m/\xi_1 r\right)} dr^2 + (1 - 2m/r\xi_1) dt^2 = \frac{1}{\left(1+\frac{\Delta\varepsilon}{1+\varepsilon}-2m/\xi_1 r\right)} dr^2 + \left(1 - \frac{2m}{r\xi_1}\right) dt^2 \quad (D12)$$

D3 A $N=0$ Application example: (mentioned on first page)

Separation Of Variables On New Pde

After separation of variables the "r" component of equation 16 (Newpde) can be written as:

$$\left[\left(\frac{dt}{ds} \sqrt{\kappa_{00}} m_p \right) + m_p \right] F - \hbar c \left(\sqrt{\kappa_{rr}} \frac{d}{dr} + \frac{j+3/2}{r} \right) f = 0 \quad D13$$

$$\left[\left(\frac{dt}{ds} \sqrt{\kappa_{00}} m_p \right) - m_p \right] f + \hbar c \left(\sqrt{\kappa_{rr}} \frac{d}{dr} - \frac{j-1/2}{r} \right) F = 0. \quad D14$$

Using the above Dirac equation component we find the anomalous gyromagnetic ratio Δgy for the spin polarized $F=0$ case. Recall the usual calculation of rate of the change of spin S gives $dS/dt \propto m \propto gyJ$ from the Heisenberg equations of motion. We note that $1/\sqrt{\kappa_{rr}}$ rescales dr in

$\left(\sqrt{\kappa_{rr}} \frac{d}{dr} + \frac{j+3/2}{r} \right) f$ in equation C5. Thus to have the same rescaling of r in the second term we

must multiply the second term denominator (i.e., r) and numerator (i.e., $J+3/2$) each by $1/\sqrt{\kappa_{rr}}$ and set the numerator ansatz equal to $(j+3/2)/\sqrt{\kappa_{rr}} \equiv 3/2 + J(gy)$, where gy is now the gyromagnetic ratio. This makes our equation D13, D14 compatible with the standard Dirac equation allowing us to substitute the gy into the Heisenberg equations of motion for spin S: $dS/dt \propto m \propto gyJ$ to find the correction to dS/dt . Thus again:

$$\begin{aligned} [1/\sqrt{\kappa_{rr}}] (3/2 + J) &= 3/2 + Jgy, \text{ Therefore for } J = 1/2 \text{ we have:} \\ [1/\sqrt{\kappa_{rr}}] (3/2 + 1/2) &= 3/2 + 1/2 gy = 3/2 + 1/2 (1 + \Delta gy) \end{aligned} \quad D15$$

Then we solve for Δg_y and substitute it into the above dS/dt equation.

Thus solve eq. D12, D15 with eq.18 values in $\sqrt{\kappa_{rr}} = 1/\sqrt{(1+\Delta\epsilon/(1+\epsilon))} = 1/\sqrt{(1+\Delta\epsilon/(1+0))} = 1/\sqrt{(1+.0005799/1)}$. Thus from equations C1,D13,D15:

$[\sqrt{(1+.0005799)}](3/2 + 1/2) = 3/2 + 1/2(1+\Delta g_y)$. Solving for Δg_y gives anomalous **gyromagnetic ratio correction of the electron** $\Delta g_y = .00116$.

If we set $\epsilon \neq 0$ (so $\Delta\epsilon/(1+\epsilon)$) instead of $\Delta\epsilon$ in the same κ_{oo} in eq.16 we get the anomalous **gyromagnetic ratio correction of the muon** in the same way.

Composite 3e: Meisner effect For B just outside r_H . (where the zero point energy particle eq. 9.22 is $.08=\pi^\pm$) See D11

Composite 3e CASE 1: Plus $+r_H$, therefore is the proton + charge component. Eq.C1 & D11

$1/\kappa_{rr} = 1 + r_H/r_H + \epsilon'' = 2 + \epsilon''$. $\epsilon'' = .08$ (eq.9.22). Thus from eq.C7: $\sqrt{2 + \epsilon''}(1.5+.5) = 1.5+.5(gy)$, $gy = 2.8$

The gyromagnetic ratio of the proton

Composite 3e CASE 2: negative r_H , thus charge cancels, zero charge:

$$1/\kappa_{rr} = 1 - r_H/r_H + \epsilon'' = \epsilon'' \quad \text{Therefore from equation D15 and case 1 eq.12 } 1/\kappa_{rr} = 1 - r_H/r_H + \epsilon''$$

$$\sqrt{\epsilon''}(1.5+.5) = 1.5+.5(gy), \quad gy = -1.9.$$

the **gyromagnetic ratio of the neutron** with the other charged and those ortho neutral hyperon magnetic moments scaled using their masses by these values respectively.

D4 Separation of Variables

After separation of variables the “r” component of equation 16 (Newpde) can be written as

$$\left[\left(\frac{dt}{ds} \sqrt{\kappa_{00}} m_p \right) + m_p \right] F - \hbar c \left(\sqrt{\kappa_{rr}} \frac{d}{dr} + \frac{j+3/2}{r} \right) f = 0 \quad D16$$

$$\left[\left(\frac{dt}{ds} \sqrt{\kappa_{00}} m_p \right) - m_p \right] f + \hbar c \left(\sqrt{\kappa_{rr}} \frac{d}{dr} - \frac{j-1/2}{r} \right) F = 0. \quad D17$$

Comparing the flat space-time Dirac equation to the left side terms of equations C5 and C6:

$$(dt/ds)\sqrt{\kappa_{00}} = (1/\kappa_{00})\sqrt{\kappa_{00}} = (1/\sqrt{\kappa_{00}}) = \text{Energy} = E \quad D18$$

Note for electron motion around hydrogen proton $mv^2/r = ke^2/r^2$ so $KE = 1/2 mv^2 = (1/2)ke^2/r = PE$ potential energy in $PE + KE = E$. So for the electron (but not the tauon or muon that are not in this orbit) $PE_e = 1/2 e^2/r$. Here write the hydrogen energy and pull out the electron contribution. So in eq.B1 and D18 $r_H = (1+1+.5)e^2/(m_\tau + m_\mu + m_e)/2 = 2.5e^2/(2m_p c^2)$. D19

Variation $\delta(\psi^*\psi) = 0$ At $r = n^2 a_0$

Next note for the variation in $\psi^*\psi$ is equal to zero at maximum $\psi^*\psi$ probability density where for the hydrogen atom is at $r = n^2 a_0 = 4a_0$ for $n=2$ and the $\psi_{2,0,0}$ eigenfunction. Also recall eq.B1 $\xi_1 = m_L c^2 = (m_\tau + m_\mu + m_e)c^2 = 2m_p c^2$ normalizes $1/2 ke^2$ (Thus divide $\tau + \mu$ by 2 and then multiply the whole line by 2 to normalize the $m_e/2$. result. $\epsilon = 0$ since no muon ϵ here.): Recall in eq.19 ξ_0 has to be pulled in a Taylor expansion as an operator since it a separate observable. So substituting eqs.D16,C1 and eq.D12 for κ_{00} , and B1,eq19 values in eq.D18:

$$E_e = \frac{(tauon + muon)(\frac{1}{2})}{\sqrt{1 - \frac{r_{H'}}{r}}} - (tauon + muon + PE_\tau + PE_\mu - m_e c^2) \frac{1}{2} =$$

$$2(m_\tau c^2 + m_\mu c^2) \frac{1}{2} + 2 \frac{m_e c^2}{2} + 2 \frac{2.5e^2}{2r(m_L c^2)} m_L c^2 - 2 \frac{2e^2}{2r(m_L c^2)} m_L c^2 - 2 \frac{3}{8} \left(\frac{2.5e^2}{r m_L c^2} \right)^2 m_L c^2$$

$$- 2(m_\tau c^2 + m_\mu c^2) \frac{1}{2}$$

$$= \frac{2m_e c^2}{2} + 2 \frac{e^2}{4r} - 2 \frac{3}{8} \left(\frac{2.5}{r m_L c^2} \right)^2 m_L c^2 = m_e c^2 + \frac{e^2}{2r} - 2 \frac{3}{8} \left(\frac{2.5e^2}{r m_L c^2} \right)^2 m_L c^2$$

So: $\Delta E_e = 2 \frac{3}{8} \left(\frac{2.5}{r m_L c^2} \right)^2 m_L c^2 = (\text{Third order } \sqrt{\kappa_{\mu\mu}} \text{ Taylor expansion term}) =$

$$\Delta E = 2 \frac{3}{8} \left[\frac{2.5(8.89 \times 10^9)(1.602 \times 10^{-19})^2}{(4(.53 \times 10^{-10}))^2 ((1.67 \times 10^{-27})(3 \times 10^8)^2)} \right]^2 (2(1.67 \times 10^{-27})(3 \times 10^8)^2)$$

$= hf = 6.626 \times 10^{-34} \times 27,360,000$ so that $f = 27 \text{ MHz}$ Lamb shift.

The other 1050 MHz comes from the zitterbewegung cloud.

Note: Need infinities if **flat space Dirac 1928 equation**. For flat space $\partial g_{ik}/\partial x^j = 0$ as a limit. Then must take field $g^{km} = 1/0 = \infty$ to get finite Christoffel symbol $\Gamma^{m}_{ij} \equiv (g^{km}/2)(\partial g_{ik}/\partial x^j + \partial g_{jk}/\partial x^i - \partial g_{ij}/\partial x^k) = (1/0)(0) = \text{undefined}$ but still implying *nonzero* acceleration on the left side of the

geodesic equation: $\frac{d^2 x^\mu}{ds^2} = -\Gamma^{\mu}_{\nu\lambda} \frac{dx^\nu}{ds} \frac{dx^\lambda}{ds}$ So we need infinite fields for flat space. Thus QED

requires (many such) infinities. But we have in general curved space $g_{ij} = \kappa_{ij}$ in the New pde so do not require that anything be infinite and yet we still obtain for the third order Taylor expansion term of $\sqrt{\kappa_{\mu\nu}}$ the Lamb shift and anomalous gyromagnetic ratio correction (see above sections C2, C3).

So renormalization is a perturbative way (given it's flat space Dirac equation and minimal interaction gauge origins) of calculating these (above) same, *NON* perturbative results, it's a perturbative GR theory. But renormalization gives lots of wrong answers too, eg., $10^{96} \text{ grams/cm}^3$ vacuum density for starters. (So we drop it here since we don't need it any longer for the high precision QED results.) In contrast note near the end of reference 5 our $G_{00} = 0$ for a 2D MS. Thus a vacuum really is a vacuum. Also that large $\xi_1 = \tau(1 + \varepsilon')$ in r_H in eq.14 is the reason leptons appear point particles (in contrast to the small ξ_0 in the composite 3e baryons).

D5 N=1 internal Observer cosmological physics from Observer at N=2

From Newpde (eg., eq.1.13 Bjorken and Drell) $i\hbar \frac{\partial \psi}{\partial t} = \frac{\hbar c}{i} \left(\alpha_1 \frac{\partial \psi}{\partial x^1} + \alpha_2 \frac{\partial \psi}{\partial x^2} + \alpha_3 \frac{\partial \psi}{\partial x^3} \right) +$

$\beta m c^2 \psi = H \psi$. For electron at rest: $i\hbar \frac{\partial \psi}{\partial t} = \beta m c^2 \psi$ so: $\delta z = \psi_r = w^r(0) e^{-i\varepsilon_r \frac{mc^2}{\hbar} t}$ $\varepsilon_r = +1,$

$r=1,2; \varepsilon_r=-1, r=3,4$.): This implies an oscillation frequency of $\omega = mc^2/\hbar$. So the eq.12 the 45° line has this ω oscillation on that δz rotation. The next higher cosmological independent (but still connected by superposition of speeds) fractal scale N=1 the 45° small Mandelbulb chord ε (Fig6)

is now getting smaller with time $t \propto \varepsilon$ as in a separation of variables result: $i\hbar \frac{\partial \psi}{\partial t} =$

$$\beta \sum_N (10^{40N} (\omega t)_{\varepsilon + \Delta \varepsilon}) \psi = \beta \sum_N (10^{40N} m_{\varepsilon + \Delta \varepsilon} c^2 / \hbar) \psi \text{ and so for stationary } N=1 \delta z = \sqrt{\kappa_{00}} dt =$$

$$e^{-i\varepsilon_r \frac{mc^2}{\hbar} t} \rightarrow e^{i(\varepsilon + \Delta \varepsilon)} \quad (18)$$

On our own fractal cosmological scale we are in the expansion stage of one such oscillation.

Recall $N > 0 \equiv \text{observer}$. Here we find what that N=2 fractal scale observer sees what we see if $\sin \mu \rightarrow \sinh \mu$ for $r > r_H$ going to $r < r_H$ in $E = 1/\sqrt{\kappa_{00}} = 1/\sqrt{(1 - r_H/r)}$ since the E in $\delta z = e^{iEt} \equiv e^{i\mu}$ and so μ

then becomes imaginary. Recall limit R_{ij} as $r \rightarrow 0$ is the source, where gravity creates gravity in the Einstein equations which becomes the modulation of the DeSitter ball. (6.14.2).

$R_{22} = e^{-\lambda} [1 + \frac{1}{2} r(\mu' - \nu')] - 1$ with $\mu = \nu$ (spherical symmetry) and $\mu' = -\nu'$. So as $r \rightarrow 0$, $\text{Im} R_{22} =$.

$\text{Im}(e^\mu - 1) = \mu + \dots = \sin \mu = \mu + \dots$ for outside r_H imaginary μ for small r (at the source) so $\sin \mu$ becomes a gravitational source (gravity itself can create gravity as a feedback mechanism). The $N=2$ observer then multiplies by $i R_{22}$, $-\sin \mu$ and μ to get $R_{22} = -\sinh \mu$ to see what the $N=2$ observer sees that we see inside r_H so:

$R_{22} = e^{-\nu} [1 + \frac{1}{2} r(\mu' - \nu')] - 1 = -\sinh \nu = -(e^\nu - e^{-\nu})/2$, $\nu' = -\mu'$ so

$e^{-\mu} [-r(\mu')] = -\sinh \mu - e^{-\mu} + 1 = -(-e^{-\mu} + e^\mu)/2 - e^{-\mu} + 1 = -(e^{-\mu} + e^\mu)/2 + 1 = -\cosh \mu + 1$. So given $\nu' = -\mu'$

$e^{-\nu} [-r(\mu')] = 1 - \cosh \mu$. Thus

$e^{-\mu} r(d\mu/dr) = 1 - \cosh \mu$

This can be rewritten as: $e^\mu d\mu / (1 - \cosh \mu) = dr/r$ (D20)

The integration is from $\xi_1 = \mu = \varepsilon = 1$ to the present day mass of the muon $= .06$ (X tauon mass).

Integrating equation B from $\varepsilon = 1$ to the present ε value we then get:

$\ln(r_{M+1}/r_{bb}) + 2 = [1/(e^\mu - 1) - \ln[e^\mu - 1]]/2$ (D215)

then $r_{bb} \approx 50 \text{ Mkm} = \text{mercuron}$ (initial $r = r_H$ each baryon. Big bang 10^{82} baryons sect.2.3). Solve for r_{M+1} , as function of μ . Find present derivative, find du from Hubble constant normalize the number to 13.7 to find total time u . Find we are now at 370by. This long of time explains the cbr thermalization and mature galaxies at dawn (instead of $\sim 200 \text{ My}$ after bb, it is 370by). The zitterbewegung (sound wave) of object B creates the condensation (at the Shapely concentration) and rarefaction void in Eridanus: we are astronomically observing here selfsimilar object B sound waves *inside* of a proton.

After a large expansion from r_{bb} our eq.14 eq.15 Schwarzschild finally becomes **Minkowski** $ds^2 = -dx_0^2 + \sum_{i=1}^n dx_i^2$. The submanifold is $-x_0^2 + \sum_{i=1}^n x_i^2 = \alpha^2$

In static coordinates r, t : (the **New pde** zitterbewegung **harmonic coordinates** x_i for $r < r_H$)

$x_0 = \sqrt{(\alpha^2 - r^2)} \sinh(t/\alpha)$: ($\sinh t$ is small t limit of equation D15. 5Tyears is the period $\gg 370 \text{ by}$)

$x_1 = \sqrt{(\alpha^2 - r^2)} \cosh(t/\alpha)$:

$x_i = r z_i$ $2 \leq i \leq n$ z_i is the standard imbedding $n-2$ sphere. R^{n-1} which also implies the **De Sitter** metric: $ds^2 = -(1 - r^2/\alpha^2) dt^2 + (1 - r^2/\alpha^2)^{-1} dr^2 + d\Omega_{n-2}^2$ (D16) **our observed ambient metric**

D6 Mixed states of $\Delta \varepsilon$ and ε $N = -1$ outside so $1S_{1/2}$ state with r

$H_{N=-1} \Delta x \Delta(m_{N=-1} c) = \hbar/2$. $m_{N=-1} = 10^{-40} m_e$. So $\Delta x = 10^5 \text{ LY}$ galaxy. $1S_{1/2}$ state may be flattened since such states are stable since $g_{00} = \kappa_{00}$.

From D13 metric source note $\Delta \varepsilon$ and ε operators so $\Delta \varepsilon \varepsilon$ (operating on Newpde ψ_N) is a new state, a "mixed state" that in the next higher scale classical limit then is a grand canonical ensemble with nonzero chemical potential (i.e., a "mixture" of systems). 2nd derivative of $\cos x = -\cos x$ so $\Delta g_{00} = -g_{00} = \cos \Delta \varepsilon$. That $g_{00} = \kappa_{00}$ in the *halo of the Milky Way galaxy* is the fundamental equation of metric quantization for all the multiples of 100 metric quantization, but here for $r < r_H$. So in general $\kappa_{00} = e^{i(m_e + m_\mu)}$, $m_e = .000058$ is the electron mass (as a fraction of the Tauon mass.) which is the component in the resulting m_e, m_μ operator sequence.

Imaginary part R_{22} locally for 2D MS $R_{00} = \Delta g_{00} = \kappa_{00}(R/2) = \cos \Delta \varepsilon$ gives also the local mixed $\Delta \varepsilon, \varepsilon$ states of part III metric quantization. Set $\cos(\Delta \varepsilon / (1 - 2\varepsilon)) = \kappa_{00} = g_{00}$, $mv^2/r = GMm/r^2$ so $GM/r = v^2$ COM in the galaxy halo (circular orbits) $(1/(1 - 2\varepsilon))$ term from D9a just inside r_H so

Pure state $\Delta \varepsilon$ (ε excited $1S_{1/2}$ state of ground state $\Delta \varepsilon$, so not same state as $\Delta \varepsilon$)

$\text{Rel}\kappa_{00} = \cos\mu$ from D9

$$\text{Case 1 } 1-2GM/(c^2r) = 1-2(v/c)^2 = 1-(\Delta\epsilon/(1-2\epsilon))^2/2 \quad (\text{D17})$$

So $1-2(v/c)^2 = 1-(\Delta\epsilon/(1-2\epsilon))^2/2$ so $=(\Delta\epsilon/(1-2\epsilon))c/2 = .00058/(1-(.06)^2)(3 \times 10^8)/2 = 99 \text{ km/sec}$
 $\approx 100 \text{ km/sec}$ (Mixed $\Delta\epsilon, \epsilon$, states classically here are grand canonical ensembles with nonzero chemical potential.). For ringed (not hub) galaxies the radial value becomes $100/2 = 50 \text{ km/sec}$.

Mixed state $\epsilon\Delta\epsilon$ (Again $GM/r = v^2$ so $2GM/(c^2r) = 2(v/c)^2$.)

$$\text{Case 2 } g_{00} = 1-2GM/(c^2r) = \text{Rel}\kappa_{00} = \cos[\Delta\epsilon + \epsilon] = 1 - [\Delta\epsilon + \epsilon]^2/2 = 1 - [(\Delta\epsilon + \epsilon)^2/(\Delta\epsilon + \epsilon)]^2/2 = 1 - [(\Delta\epsilon^2 + \epsilon^2 + 2\epsilon\Delta\epsilon)/(\Delta\epsilon + \epsilon)]^2$$

The $\Delta\epsilon^2$ is just the above first case (Case 1) so just take the mixed state cross term $[\epsilon\Delta\epsilon/(\epsilon + \Delta\epsilon)] = c[\Delta\epsilon/(1 + \Delta\epsilon/\epsilon)]/2 = c[\Delta\epsilon + \Delta\epsilon^2/\epsilon + \dots \Delta\epsilon^{N+1}/\epsilon^N + \dots]/2 = \Sigma v_N$. Note each term in this expansion is itself a (mixed state) operator. So there can't be a single v in the large gradient 2nd case so in the equation just above we can take $v_N = [\Delta\epsilon^{N+1}/(2\epsilon^N)]c$. (D18)

From eq. D18 for example $v = m100^N \text{ km/sec}$. $m=2, N=1$ here (Local arm). In part III we list hundreds of examples of D18: (sun1,2km/sec, galaxy halos m100km/sec). The linear mixed state subdivision by this ubiquitous ~ 100 scale change factor in r_{bb} (due to above object B zitterbewegung spherical Bessel function resonance boundary conditions resulting in nodes) created the voids. Same process for $N-1$ (so 100X smaller) antinodes get galaxies, 100X smaller: globular clusters, 100X smaller solar systems, etc., So these smaller objects were also created by mixed state metric quantization (eq.D18) resonance oscillation inside initial radius r_{bb} .

We include the effects of that object B drop in inertial frame dragging on the inertial term m in the Gamow factor and so lower Z nuclear synthesis at earlier epochs ($t > 18 \text{ by}$) BCE. (see part III)

Appendix E Δ Modification of Usual Elementary Calculus ϵ, δ 'tiny' definition of the limit.

Recall that: given a number $\epsilon > 0$ there exists a number $\delta > 0$ such that for all x in S satisfying

$$|x - x_0| < \delta$$

we have

$$|f(x) - L| < \epsilon$$

Then write $\lim_{x \rightarrow x_0} f(x) = L$

Thus you can take a smaller and smaller ϵ here, so then $f(x)$ gets closer and closer to L even if x never really reaches x_0 . "Tiny" for $h \rightarrow L_1$ and $f(x+h) - f(x) \rightarrow L_2$ then means that $L=0 = L_1$ and L_2 . 'Tiny' is this difference limit.

Hausdorf (Fractal) s dimensional measure using ϵ, δ

Diameter of U is defined as $|U| = \sup\{|x - y| : x, y \in U\}$. $E \subset \cup_i U_i$ and $0 < |U_i| \leq \delta$

$$H_\delta^s(E) = \inf \sum_{i=1}^{\infty} |U_i|^s$$

analogous to the elementary $V = U^s$ where of $s=3$, $U=L$ then V is the volume of a cube $\text{Volume} = L^3$. Here however 's' may be noninteger (eg., fractional). The volume here would be the respective Hausdorf outer measure.

The infimum is over all countable δ covers $\{U_i\}$ of E .

To get the Hausdorf outer measure of E we let $\delta \rightarrow 0$ $H^s(E) = \lim_{\delta \rightarrow 0} H_\delta^s(E)$

The restriction of H^s to the σ field of H^s measurable sets is called a Hausdorf s-dimensional measure. $\text{Dim } E$ is called the Hausdorf dimension such that

$$H^s(E) = \infty \text{ if } 0 \leq s < \text{dim } E, \quad H^s(E) = 0 \text{ if } \text{dim } E < s < \infty$$

So if s implies a zero H or infinite H it is not the correct dimension. This rule is analogous to the definition of the (fractal) Mandelbrot set itself in which a C that gave infinity is rejected by the definition $\delta C=0$ we can model as a binary pulse ($z=zz$ solution is binary $z=1,0$) with

$zz=z$ (1) is the algebraic definition of 1 and can add real constant C (so $z'=z'z'-C$, $\delta C=0$ (2)), $z \in \{z'\}$

Plug $z'=1+\delta z$ into eq.2 and get

$$\delta z + \delta z \delta z = C \quad (3)$$

so

$$\delta z = (-1 \pm \sqrt{1 + 4C})/2 = dr + idt \quad (4)$$

for $C < -1/4$ so real line $r=C$ is immersed in the complex plane.

$z=z_0=0$ To find C itself substitute z' on left (eq.2) into right $z'z'$ repeatedly & get $z_{N+1}=z_N z_N - C$. $\delta C=0$ requires us to reject the C s for which

$-\delta C = \delta(z_{N+1} - z_N z_N) = \delta(\infty - \infty) \neq 0$. **$z=zz$ solution is 1,0** so initial

gets the **Mandelbrot set** C_M (fig2) out to some $\|\Delta\|$ distance from $C=0$. Δ found from $\partial C/\partial t=0$, $\delta C \equiv \delta C_r = (\partial C_M / \partial (dr dt)) dr = 0$ extreme giving the Feigenbaum point $\|C_M\| = \|-1.400115..\|$ global max given this $\|C_M\|$ is biggest of all.

If s is not an integer then the dimensionality it has a fractal dimension.

But because the Feigenbaum point Δ uncertainty limit is the r_H horizon, which is impenetrable (sect.2.5, partI), ε, δ are not dr/ds eq.11a observables for $0 < \varepsilon, \delta < r_H$. Instead $\varepsilon, \delta > \Delta = r_H$ = the next $10^{40} \times$ smaller fractal scale Mandelbrot set at the Feigenbaum point.

Appendix F

Review This is an Occam's razor *optimized* (i.e., $\delta C=0$, $\|C\|=\text{noise}$)

POSTULATE OF 1

So

$z=zz$ (1) is the algebraic definition of 1, o, add real constant C (i.e., $z'=z'z'$, $\delta C=0$) (2), $z \in \{z'\}$

Digital communication analogy: Binary ($z=zz$) 1,0 signal with white noise $\delta C=0$ in $z'+C=z'z'$. Recall the algebraic definition of 1 is $z=zz$ which has solutions 1,0. Also you could say white noise C has a variation of zero ($\delta C=0$) making it easy to filter out (eg., with a Fourier cutoff filter).

So you could easily make the simple digital communication analogy of this being a binary ($z=zz$) 1,0 signal with white noise $\delta C=0$ in $z'+C=z'z'$.

(However the noise is added a little differently here ($z+C=zz$) than in statistical mechanics signal theory (eg., There you might use deconvolved signal=convolution integral

$[(\text{transfer function})\text{signal}]dA$)). where the 'signal' actually would equal $z+C$ So this is not quite the same math as in statistical mechanics.)