

Postulate 1 (observable) and define observability

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Abstract You would think just postulating 1 is goofy. That can't lead anywhere, right? But wait a minute!

The algebraic definition of 1 is $z=zz$ (with $z=1,0$) and merely adding constant C (so $\delta C=0$) as in $z'=z'z'+C$ (eq.1)) derives the universe!!!!!!!!!!!!!!!!!!!!!!!!!!!!

It is the most amazing thing I have ever seen!!!!!!!!!!!!!!!!!!!!!!!!!!!!

Thus use the $z=1,0$ (from postulated $z=zz$ as part of the larger set $\{z'\}$) to solve for all z' and C:

The $z=0 =z_0 =z'$ iteration of eq.1 using $\delta C=0$ gets the (2D)Mandelbrot set $C=C_M$ and a δz
(Need iteration to get all the Cs because of the $\delta C=0$, See sect.1. Also 10^{40N} fractal scales)

The $z=1, z'=1+\delta z$ substitution into eq.1 using $\delta C=0$ ($N>0$ observer) gets the 2D Dirac eq.

(Minkowski (flat space) metric, Clifford algebra and eq.11 *fall right out of* resulting eq.5)

These $z=1$ and $z=0$ steps together (sect.1) get the curved space $2D+2D=4D$ **Newpde** (3)

and thus the 4D universe, no more and no less (ref.5) So just **postulate 1** !!!

All I did here is to **define observability***: 1 is the observable and C the required observer!

Furthermore $N=0$ **postulate 1** can also be used in a list-define math to get the real number algebra (without all those many Rel#math axioms).Eg., $1 \cup 1 \equiv 1+1$ (B2,Ch.2). So we get both the physics (ref5) AND (rel#)mathematics from ONE postulate1, everything!

We figured it out. We understand the universe finally! Completely! Just **postulate1**

1 (in $z=zz$) is thereby the ultimate application of Occam's razor since
0 (in $z=zz$) postulates literally *nothing* and $z=zz$ is simpler than $z=zzzz$
 $\delta C=0$ means Occam's razor optimized (lowest observer noise C)

Hey, wait a second, so do you mean to say that all I have to do is postulate 1 and get no more and no less than the entire physical universe including even the math? Yes!! davidmaker.com

*Eq.11: $p_x \psi = -i \hbar \partial \psi / \partial x$. is the well known observable p criteria (appendix). ψ is from Newpde(3)

It's Broken, fix it (We badly need that **Newpde**)

In that regard Dirac in 1928 made his equation(1) flat space(2). But space is not in general flat, there are forces.

So over the past 100 years people have had to try to make up for that mistake by adding ad hoc convoluted gauge force after gauge force until fundamental theoretical physics became a mass of confusion, a train wreck, a junk pile. So all they can do for ever and ever is to rearrange that junk pile with *zero actual progress* in this most fundamental of all sciences,.. forever. We died.

By the way note that **Newpde**(3) $\gamma^\mu \sqrt{\kappa_{\mu\mu}} \partial \psi / \partial x_\mu = (\omega/c) \psi$ is NOT flat space (4) so it cures this problem (5).

References

(1) $\gamma^\mu \partial \psi / \partial x_\mu = (\omega/c) \psi$

(2) Spherical symmetry: $(\gamma^x \sqrt{\kappa_{xx}} dx + \gamma^y \sqrt{\kappa_{yy}} dy + \gamma^z \sqrt{\kappa_{zz}} dz + \gamma^t \sqrt{\kappa_{tt}} dt)^2 = \kappa_{xx} dx^2 + \kappa_{yy} dy^2 + \kappa_{zz} dz^2 - \kappa_{tt} dt^2 = ds^2$
 $\kappa_{xx} = \kappa_{yy} = \kappa_{zz} = \kappa_{tt} = 1$ is flat space, Minkowski, as in his Dirac equation(1).

(3) Newpde: $\gamma^\mu \sqrt{(\kappa_{\mu\mu})} \partial\psi/\partial x_\mu = (\omega/c) \psi$ for e, ν . So we didn't just drop the $\kappa_{\mu\nu}$ (as is done in ref.1)

(4) Here $\kappa_{\rho\rho} = 1 - r_H/r = 1/\kappa_r$, $r_H = (2e^2)(10^{40N})/(mc^2)$. The $N = \dots -1, 0, 1, \dots$ fractal scales (appendix)

(5) $N=0$ the greater than 2nd order Taylor expansion terms of $\sqrt{\kappa_{ij}}$ give the Lamb shift and anomalous gyromagnetic ratio (without the renormalization and infinities (see appendix D)).

For $N=-1$ (i.e., $e^2 \times 10^{-40} \equiv Gm_e^2$) κ_{ij} is by inspection the Schwarzschild metric g_{ij} ; so we just derived General Relativity and the gravity constant G from Quantum Mechanics in one line.

So $\kappa_{\mu\nu}$ provides the general covariance of the Newpde. Eq. 4 even provides us space-time r, t .

For $N=1$ (so $r < r_C$) Newpde zitterbewegung expansion stage explains the universe expansion (For $r > r_C$ it's not observed, per Schrodinger's 1932 paper.).

For $N=1$ zitterbewegung harmonic coordinates and Minkowski metric submanifold (after long time expansion) gets the De Sitter ambient metric we observe (D16).

For $N=0$ Newpde $r=r_H$ $2P_{3/3}$ state composite $3e$ is the baryons (appendixC, partII) and Newpde $r=r_H$ composite e, ν is the 4 Standard electroweak Model Bosons (eq.12 \rightarrow appendixA).

We got all of physics here by mere inspection of this (curved space) Newpde with no gauges. The Newpde fixes it.

Compare and contrast

The core of mainstream physics is the Standard electroweak Model (SM) that gives us important results like Maxwell's equations and weak interaction theory that explain electricity and magnetism and some radioactive decays respectively. Add to that QCD that explains the nuclear force (NF) and baryons. General Relativity (GR) gives us gravity and mechanics. But they are not fundamental since they contain *many assumptions* (Lagrangian densities, free parameters, many dimensions, gauge symmetries,..etc..) of unknown origin.

In contrast

what if you found instead a mathematical theory with only one *simple assumption* (eg., '1', defined from $z=zz$ since $1=1 \times 1$) using a *single simple math step* top down that got a *generally covariant generalization of the Dirac equation that does not require gauges* (Newpde that in turn gave these same results (i.e., SM particles, NF, GR, QM in ref.5 & real#)? You will have a truly *fundamental* theory (See above postulate 1)

i.e., Postulate 1 \rightarrow Newpde. davidmaker.com So just postulate 1

Appendix: Postulate $z=zz$ ($z=1,0$) and merely add constant C (so $\delta C=0$) as in $z'=z'z'+C$ eq1

Details of those two $z=0, z=1$ steps (in above **definition of observability**)

But first solve equation 1 by itself (to at least see that z' can be complex)

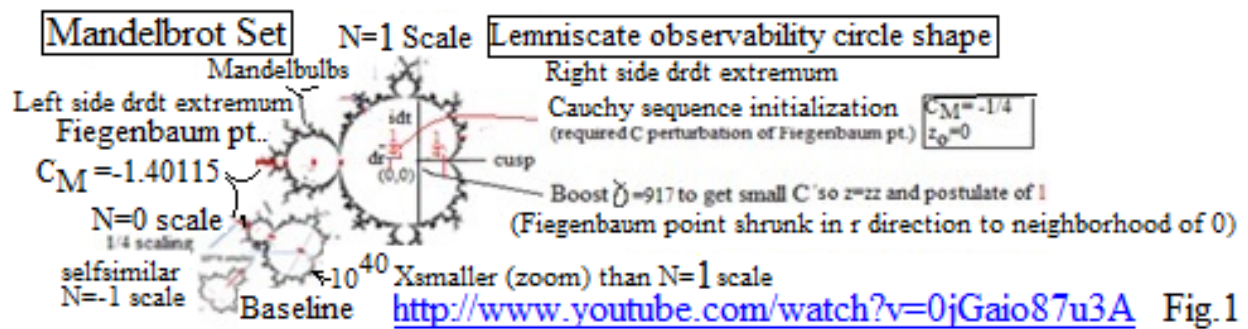
Thus plug $z'=1+\delta z$ into eq.1 and get $\delta z + \delta z \delta z = C$ (3)

For real $C < -1/4$ $\delta z = (-1 \pm \sqrt{1 + 4C})/2 = dr + idt$ (4)

is complex. Knowing z' is complex we then can then move on to the *complete* solution:

1st step: (Recall postulated $1,0 \in \{z'\}$ with rest of z' unknown)

$\mathbf{z=0=z_0=z'}$ To find all C substitute z' on left (eq.1) into right $z'z'$ repeatedly and get iteration $z_{N+1}=z_N z_N - C$. Constraint $\delta C=0$ requires us to reject the C s for which $-\delta C = \delta(z_{N+1} - z_N z_N) = \delta(\infty - \infty) \neq 0$ which gets the **Mandelbrot set** C_M (fig1) out to some $\|\Delta\|$ distance from $C=0$. That bare fig1 Mandelbrot set structure is picked out of the zoom because it is built of lemniscate circles (sect.2.2,fig.7) and so implies eq.11 observables for fig.1 shapes. Also Δ is found from $\delta C \equiv \left(\frac{\partial C}{\partial A}\right)_{dt} dr = 0$ (sect.2.1). The smallest bulbous lobe of area 'A=drdt' as a function of largest $d\|z\|$ (which is a -dr) is at the Fieigenaum point -1.400115.. So C_M ends up being real here so we really can use the above *real* $C < -1/4$ approximation and so require complex z . That $-1/4$ iteration also makes dr, dt real so $dr^2 + dt^2 = C^2$.



Note in distance units of 1 unit on the $N=0$ fractal scale, for (big C) observer $N=1$ scale lengths $\delta z \gg 1$ which applies to eq.11 *observers*.

2nd step:

$\mathbf{z=1}$ in $z'=1+\delta z$ in eq.1 get eq.3 ($\delta z \gg 1$): $\delta(\delta z + \delta z \delta z) = \delta \delta z (1) + \delta \delta z (\delta z) + (\delta z) \delta \delta z \approx \delta(\delta z \delta z) = 0 = (\text{plug in eq.4}) = \delta[(dr+idt)(dr+idt)] = \delta[(dr^2-dt^2) + i(drdt+dt dr)] = 0$ (5)
 =(Minkowski metric)+i(Clifford algebra)

Factor eq.5 **real** $\delta(dr^2-dt^2) = \delta[(dr+dt)(dr-dt)] = 0 = [[\delta(dr+dt)](dr-dt)] + [(dr+dt)[\delta(dr-dt)]] = 0$ (6)

so $(\rightarrow \pm e)$ $dr+dt=ds$, $dr-dt=ds \equiv ds_1$, for $(-dr-dt)^2=ds^2 \rightarrow$ Ist and IVth quadrant in fig3 (7)

Also note the positive scalar $drdt$ of eq.7 implies the eq.5 *non* infinite extremum **imaginary** $drdt+dt dr = 0 = \gamma^i dr^j dt + \gamma^j dt \gamma^i dr = (\gamma^i \gamma^j + \gamma^j \gamma^i) drdt$ so Clifford algebra $(\gamma^i \gamma^j + \gamma^j \gamma^i) = 0, i \neq j$. (7a)

$(\rightarrow \text{light cone } v)$ $dr+dt=ds$, $dr=-dt$, for $(-dr-dt)^2=ds^2 \rightarrow$ III quadrant (8)

“ “ $dr-dt=ds$, $dr=dt$, for $(-dr-dt)^2=ds^2 \rightarrow$ II quadrant (9)

$(\rightarrow \text{vacuum, } z=1)$ $dr=dt$, $dr=-dt$ so $dt=0=dr$ (So eigenvalues of $dt, dr=0$ in eq.11) (10)

We square eqs.7,8,9 $ds_1^2 = (dr+dt)(dr+dt) = (-dr-dt)(-dr-dt) = [dr^2+dt^2] + (drdt+dt dr) \equiv ds^2 + ds_3 = ds_1^2$. Since ds_3 (is max or min) and ds_1^2 (from eq.7,8,9) are invariant then so is **Circle** $ds^2 = dr^2 + dt^2 = ds_1^2 - ds_3$. also implying the rest of the Clifford algebra $\gamma^i \gamma^i = 1$ in eq.7a, no sum on 'i' and also the lemniscate formulation(fig.7). Note this separate ds is a minimum at 45° (that

Mandelbrot set rotation) given the eq.7 constraints and so Circle $\equiv\delta z=dse^{i\theta}=dse^{i(\Delta\theta+\theta_0)}=dse^{i((\cos\theta dr+\sin\theta dt)/(ds)+\theta_0)}$, $\theta_0=45^\circ$. We define $k\equiv dr/ds$, $\omega\equiv dt/ds$, $\sin\theta\equiv r$, $\cos\theta\equiv t$. $dse^{i45^\circ}\equiv ds'$. Take

ordinary derivative dr (since flat space) of 'Circle' $\frac{\partial(dse^{i(\frac{rdr+tdt}{ds})})}{\partial r} = i \frac{dr}{ds} \delta z$ so $\frac{\partial(dse^{i(rk+wt)})}{\partial r} = ik\delta z$, $k\delta z = -i \frac{\partial\delta z}{\partial r}$ (11).

$\langle F \rangle^* = \int (F\psi)^* \psi d\tau = \int \psi^* F \psi d\tau = \langle F \rangle$ Hermitian). The observables $dr \rightarrow k \rightarrow p_r$ condition gotten from eq.11 **operator formalism** thereby converting eq.7-9 into Dirac eq. pdes (Circle solution also implies lemniscates, fig.7, so fig.1 observability). Cancel that e^{i45° coefficient ($45^\circ = \pi/4$) then multiply both sides of eq.11 by \hbar and define $\delta z \equiv \psi$, $p_r \equiv \hbar k$. Eq.11 becomes the familiar:

$$p_r \psi = -i\hbar \frac{\partial\psi}{\partial r} \quad (11)$$

z=1, z=0 steps together (on Circle with small C boost)

Postulate 1 also implies a small C in eq.1 which thereby implies a (Minkowski metric) Lorentz contraction $1/\gamma$ boosted frame of reference* in eq.3 $C=C_M/\gamma \equiv \delta z' = \Delta$ for next small smaller fractal scale $N_{ob} < 0$ so $\delta z' \ll 1$. Also recall $\delta ds^2 = 0$ in eq.5. Thus we have a angle perturbation of big dr, dt for $\theta_0 = 45^\circ$ on that above **ds Circle** and so have a slightly modified eq.7:

$$(dr - \delta z') + (dt + \delta z') \equiv dr' + dt' = ds \quad (12)$$

Since (eq.12) dr, dt is the (eq.11) observer $N_{observer} > 0$ scale then $\delta z'$ defines the $N_{ob} = 0$ object. Also the r, t axis' are the max extremum for ds^2 , and the ds^2 at 45° is the min extremum ds^2 so $\Delta\theta = \theta \text{ modulo } 45^\circ$ is pinned to an axis' so extreme $\Delta\theta \approx \pm 45^\circ = \delta z'$. So in eq.12 the 4 rotations $45^\circ + 45^\circ = 90^\circ$ define Bosons (appendix A), and $45^\circ - 45^\circ = 0$ eq.7-9 defines leptons. Again, for $N_{ob} < 0$, you also have other (smaller) fractal scale extreme $\delta z'$ (eg., tiny Fiegenbaum pts) so metric coefficients $\kappa_{rr} \equiv (dr/dr')^2 = (dr/(dr - (C_M/\xi_1)))^2 = 1/(1 - r_H/r)^2 = A_1/(1 - r_H/r) + A_2/(1 - r_H/r)^2$, $r \equiv dr$. The partial fractions A_i term can be split off from RN and so $\kappa_{rr} \approx 1/[1 - ((C_M/\xi_1)r)]$ (13)

(C_M defined to be charge, $\gamma \equiv \xi_1$ mass). So: $ds^2 = \kappa_{rr} dr'^2 + \kappa_{oo} dt'^2$ (14)

From eq.7a $dr' dt' = \sqrt{\kappa_{rr}} dr' \sqrt{\kappa_{oo}} dt' = dr dt$ so $\kappa_{rr} = 1/\kappa_{oo}$ (15)

We do a rotational dyadic coordinate transformation of $\kappa_{\mu\nu}$ to get the Kerr metric which is all we need for our GR applications. So from eqs.4,5,14,15 we found the relation between x_i, x_j pairs:

$(\sum_{i=1}^2 \gamma^i \sqrt{\kappa_{ii}} dx_i)^2 = \sum_{i=1}^2 \kappa_{ii} d^2 x_i$ (14a). So given this added 2D Δ perturbation we get curved space $2D \otimes 2D = 4D$ independent $x_1, x_2 \rightarrow x_1, x_2, x_3, x_4$. Also assuming orthogonality

$dr^2 \equiv dx_1^2 + dx_2^2 + dx_3^2$ (as $r \rightarrow \infty$ in eq.13,15) the right side of eq.14a therefore has the 2 in the sum replaced by a 4 implying the left side then has to be in eq.14a for the 2D form to be a special case (of any 2 x_i, x_j) in B2: Imposing orthogonality thereby creates 6 pairs of eqs.4&5. So eq.14a becomes: $(\gamma^x \sqrt{\kappa_{xx}} dx + \gamma^y \sqrt{\kappa_{yy}} dy + \gamma^z \sqrt{\kappa_{zz}} dz + \gamma^t \sqrt{\kappa_{tt}} dt)^2 = \kappa_{xx} dx^2 + \kappa_{yy} dy^2 + \kappa_{zz} dz^2 + \kappa_{tt} dt^2 = ds^2$.

Multiplying the bracketed term by $1/ds$ & $\delta z \equiv \psi$ so eq 11 implies 4D **Newpde** (lemniscates 2.1):

$$\gamma^\mu (\sqrt{\kappa_{\mu\mu}}) \partial\psi/\partial x_\mu = (\omega/c) \psi \quad \text{for } e, \nu, \quad \kappa_{oo} = 1 - r_H/r = 1/\kappa_{rr} \quad r_H = e^2 X 10^{40N} / m N (= -1, 0, 1, \dots) \quad (16)$$

$= C_M/\gamma$ (from sect.2) $C_M = \text{Fiegenbaum point}$. So: **postulate 1** \rightarrow **Newpde**. syllogism

*That small C boost is thereby created by Newpde $r = r_H$ $2P_{3/2}$ stable state P (so $2S_{1/2} = \tau$, $1S_{1/2} = \mu$, appC). The 4 eq.12 Newpde e, ν rotations at $r = r_H$ are the 4 W^+, γ, W^-, Z_0 SM Bosons (appendix A)

So Penrose's intuition(6) was right on! There *is* physics in the Mandelbrot set, all of it.