## Part II

## Newpde composite $3e 2P_{3/2}$ at r=r<sub>H</sub> state

## **Review Of Part I**

Newpde from Occam's razor *optimized* **postulate1**:

z=zz is the algebraic definition of 1 and adding (small) constant C is trivial in z+C=zz,  $\delta$ C=0 (1)

- A Substitute  $z=1+\delta z$  into equation 1 and get 2D Dirac equations for e, v (sect.1.1).
- **B** Substitute left side z into right side zz repeatedly in equation 1 and get the 2D Mandelbrot set

A,B together gives 4D Newpde. So Postulate1→Newpde

## I Summary: Algebra Details of the two substitutions (A,B):

1.1 (Simply postulate 1. But z=zz is the algebraic definition of 1 and adding constant C (orK) is trivial in:)zz=z+C,  $\delta C=0$  (1)

A Substitute  $z=1+\delta z$  into eq.1, get  $\delta z+\delta z\delta z=C$  (2) (or  $\delta(\delta z+\delta z\delta z)=0$ ) with in general (for  $\delta z=\frac{-1\pm\sqrt{1+4C}}{2}$ , C<<-<sup>1</sup>/<sub>4</sub>, so **Big**-C and so big  $\delta z$  for random time t complex solution  $\delta z=dr+idt$  (2a) Big  $\delta z$  so,  $\delta z<<\delta z\delta z$ ,  $\delta \delta z=0$  So from eq.2,2a  $\delta((\delta z+\delta z\delta z)\approx\delta(\delta z\delta z)=(dr^2-dr^2+i(drdt+dtdr))=0$  (2b)

and take real Minkowski and imaginary component Clifford algebra which for negative and positive dr,dt the *non*infinite extremum  $drdt+dtdr=0=\gamma^i dr\gamma^j dt+\gamma^j dt\gamma^i dr=(\gamma^i \gamma^j + \gamma^j \gamma^i) drdt = 0$  (3) Factor  $dr^2-dt^2=ds^2$  eg.,  $\delta[(dr+dt)(dr-dt)]=0=[[\delta(dr+dt)](dr-dt)]+[(dr+dt)[\delta(dr-dt)]]$  and solve to get: (ds = proper time invariant.)

(→± <b>e</b> )	dr+dt=ds,	$dr-dt=ds = ds_1$		for	$+ds \rightarrow$	I, IV quadrants	(4)
$(\rightarrow light cone v$	) dr+dt=ds,	dr=-dt,		for	$+ds \rightarrow$	II quadrant	(5)
	dr-dt=ds,	dr=dt,		"	"	III quadrant	(6)
(→vacuum)	dr=dt.	dr=-dt	so dt=0=dr				(7)

( $\rightarrow$ vacuum) dr=dt, dr=-dt so dt=0=dr (7) We square eq.4 ds<sub>1</sub><sup>2</sup>=(dr+dt)(dr+dt) =[dr<sup>2</sup>+dt<sup>2</sup>] +(drdt+dtdr) =ds<sup>2</sup>+ds<sub>3</sub>=ds<sub>1</sub><sup>2</sup>. Since ds<sub>3</sub> (is max or min) and ds<sup>2</sup> (from eq.4) are invariant then so is **Circle** ds<sup>2</sup>=dr<sup>2</sup>+dt<sup>2</sup> =ds<sub>1</sub><sup>2</sup>-ds<sub>3</sub>. Note this separate ds is a minimum at 45° and so Circle= $\delta z$ =dse<sup>iθ</sup>=dse<sup>i( $\Delta \theta + \theta \sigma$ )</sup>= dse<sup>i((cos\theta dr+sin\theta dt)/(ds)+\theta \sigma)</sup>,  $\theta_{\sigma}$ =45°. We define k=dr/ds,  $\omega$ =dt/ds, sin $\theta$ =r, cos $\theta$ =t. dse<sup>i45°</sup>=ds'. So take the global ordinary dr derivative

(since flat space) of 'Circle'  $\frac{\partial \left(dse^{i\left(\frac{rdr}{ds}+\frac{tdt}{ds}\right)}\right)}{\partial r} = i\frac{dr}{ds}\delta z$  so  $\frac{\partial \left(dse^{i\left(rk+wt\right)}\right)}{\partial r} = ik\delta z$ ,  $k\delta z = -i\frac{\partial \delta z}{\partial r}$  (8) Cancel that  $e^{i45^\circ}$  coefficient then multiply both sides of eq.8 by h and define  $\delta z \equiv \psi$ ,  $p \equiv hk$ . Eq.8 then implies (Hermitian) operator hk **observables** formalism(QM). See next paragraph for real k part, (also appendix B3). Eqs 4,5,6,8 and eq.3 Clifford algebra imply 2D Dirac equations for e,v. Needed Big -C to have random t and Hamiltonians (*observability*)

**B** Substitute left side z into right side zz repeatedly of eq.1 to get a Mandelbrot set iteration (for small C limit eq.1 cases  $z\approx 1,0$ ). The right side extremum Cauchy seq. iteration defines real k (dr/ds eigenvalues) in eq.8. It's left side small drdt (eq.3,  $\delta(drdt)=0$ ) extremum is the fractal Fiegenbaum pt.=C<sub>M</sub> Mandelbulb  $\delta z'$ . Note certain Cs have a higher density here so *perturbations* A&B Big  $\gamma$  (1/ $\gamma$ ) $\delta z$  boost means small C $\approx \delta z \gg \delta z \delta z$  and so no more quadratic equation but still keeps our complex  $\delta z$  and so dt. So we must Fitzgerald contract (boost= $\gamma$ )  $\delta z$  to get small C and so z=zz and the **postulate** of 1. So then boost C<sub>M</sub> as in  $\delta z'=C=C_M/\gamma=C_M/\xi=r_H$  so  $\xi_1$  (defining mass) *must be big*. For z=1 in z=1+\delta z,  $\delta z$  is small and so in C<sub>M</sub>= $\xi \delta z$ ,  $\xi_1$  *is* big and so we got

z=zz and the **postulate** of 1. But for z=0, δz is big so ξ<sub>0</sub> is small and in δC<sub>M</sub>=δξδz+ξδδz=0 then since δδz=0 then δξ small so ξ<sub>0</sub> is both stable and small (≡electron). On small scales δz >>δzδz so in eq.2 |δz|≈C=constant so can only perturb eq.4 at 45°using (dr-δz')+(dt+δz')≡dr'+dt'=ds (9) since ds=|δz|=C. Define  $\kappa_{rr}$ =(dr/dr')<sup>2</sup>= (dr/(dr-(C<sub>M</sub>/ξ<sub>1</sub>)))<sup>2</sup>=1/(1-r<sub>H</sub>/r)<sup>2</sup>=A<sub>1</sub>/(1-r<sub>H</sub>/r) +A<sub>2</sub>/(1-r<sub>H</sub>/r)<sup>2</sup> r=dr. The A<sub>1</sub> term can be split off from RN (as in classic GR) and so  $\kappa_{rr}$ ≈1/[1-((C<sub>M</sub>/ξ<sub>1</sub>)r))]. So we have: ds<sup>2</sup>=  $\kappa_{rr}$ dr'<sup>2</sup> + $\kappa_{oo}$ dt'<sup>2</sup> +..(10). Note from eq.3 dr'dt'=  $\sqrt{\kappa_{rr}}$ dr $\sqrt{\kappa_{oo}}$ dt =drdt so  $\kappa_{rr}$ =1/ $\kappa_{oo}$  so given that 2D perturbation we get curved space 4D. For 4D our eq.3 Cliffordalgebra then implies ( $\gamma^x \sqrt{\kappa_{xx}} dx + \gamma^y \sqrt{\kappa_{yy}} dy + \gamma^z \sqrt{\kappa_{zz}} dz + \gamma^t \sqrt{\kappa_{tt}} idt$ )<sup>2</sup>= $\kappa_{xx} dx^2 + \kappa_{yy} dy^2 + \kappa_{zz} dz^2 - \kappa_{tt} dt^2$ =ds<sup>2</sup>. Multiplying the bracketed term by 1/ds &δz then eq 8 implies 4D **Newpde**  $\gamma^{\mu}(\sqrt{\kappa_{\mu\mu}}) \partial \psi / \partial x_{\mu} = (\omega/c) \psi$  for e,v (11) (covariant derivative still ordinary since  $\psi$  (complex) scalar). Therefore **postulate1→Newpde** 

Small C boost gets z=zz (so postulate 1) but also gets the *numerical value* of Large  $\xi_1$ For that stable z=0 the only way to get stable large  $\xi$  (required by that small C boost) is with the Newpde composite 3e  $2P_{3/2}$  at r=r<sub>H</sub> state (partII davidmaker.com). So *stability* (dt'<sup>2</sup>=(1-r<sub>H</sub>/r)dt<sup>2</sup>) clocks stop at r=r<sub>H</sub> The two positron motion and h/2e quantization of flux BA then gives us the exact proton mass  $m_p$  as a reduced mass for the associated Hund rule  $\tau = 2S_{1/2}, 1S_{1/2} = \mu$  states (so  $\tau + \mu = \xi_1$ ,  $m_p = \xi_1/2$ ). We rewrite this in the Kerr metric formalism with the 3<sup>rd</sup> mass also reversing the pair annihilation (Thus virtual pair creation inside the r<sub>H</sub> volume given  $\sigma = \pi r_{\rm H}^2 \approx (1/20)$  barns) and reducing the inertial frame dragging due to the spin  $\frac{1}{2} \xi_1$  thereby adding a Kerr metric  $-(a/r)^2$  $1+\varepsilon+\Delta\varepsilon+2e^2/(\xi_0 r)=\kappa_{00}$  (Fiegenbaum pt. C<sub>M</sub> defines charge  $e^2$ .). Divide by  $\xi_1=1+\varepsilon$  to normalize for only free electron  $\Delta \varepsilon$  energy (Needed for the following two electron applications in eqs.15,16) asymptotic local flat space and thereby finally getting back to that initial requirement for that free particle z=1, *large*  $\xi_1$  case:  $\kappa_{oo}=1-\xi_o/(1+\epsilon)-C_M/(\xi_1 r)=1+\Delta\epsilon/(1+\epsilon)-2e^2/(2m_p r)$  (12) also giving us the *numerical value* of that large  $\xi_1$  (=2m<sub>p</sub>). With  $\tau$  normalized to  $\tau$ =1 with the Newpde ground state e mass then  $\Delta \varepsilon = m_e = .0005799$  with  $\varepsilon = \mu = .06$ . (12a)

### **Composite 3e**

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## **Toy Model Derivation**

## 7.1 Derivation of Mainstream Toy Model Composite System |1>|2>|3> Interpretation From Newpde 3e State

To explain the above 3e stability result sect.1.4 implies  $r=r_{HP}$  in dt'=0 in dt'<sup>2</sup>=  $\kappa_{oo}dt^2=(1-r_{HP}/r)dt^2$  (in the new pde, nucleon radius) and so dt'=0 so that clocks stop. So we have complete proton stability in the new pde given eq. 11  $2P_{3/2}$  at  $r=r_{HP}$  fills first (see review section just above). Eq.11,  $2P_{3/2}$  is trifolium shaped  $\psi^*\psi$  so the electron spends 1/3 time in each lobe (fractional charge), lobes can't leave (asymptotic freedom), P wave scattering (jets), 6 P states (6flavors udscbt) explaining the major properties of quarks and so explaining the strong interaction. Note that  $ds_1^2 \neq 0$  with  $dt \approx 0$  so  $\sqrt{2}ds^2 = dr^2 + dt^2$  implies, after eq.3.6 operator formalism derivatives are put in (figure 4 2<sup>nd</sup> diagram from right), the Klein Gordon equation spin 0 mesons, the carriers of the strong force. Also our Mandlebulb analysis implies that the proton mass is 937Mev (appendix C) implying ultrarelativistic internal electron motion is needed to get this mass. Given this ultrarelativistic electron (needed to obtain the much larger proton mass hard shell result P. Alberto, R. Lisboa, M. Malheiro and A. S. de Castro, Phys.Rev. .... 58 (1998) R628) at r<sub>H</sub> the field lines must be Fitzgerald contracted to a flat "plate" and thus be high density, field lines thus *explaining the strength of the 'strong' force*. Note here we have three ultrarelativistic 2P<sub>3/2</sub> electrons (2 positrons and one electron) needed to create the proton (which as we noted above is much heavier) at r=r<sub>H</sub>. Given a central (negative) electron the two outer positron  $2P_{3/2}$  state plates (at 120°) only intersect at the center and so don't see each other at all and so don't repel each other, *explaining why the proton is still a bound state even though the* two positive positrons are both inside  $r_H$  (along with that electron). Possible (but low probability) positron-electron annihilation inside r<sub>H</sub> also implies, given the momentum transfer to the *third* particle with strong field plate and large mass (quadruply differential cross-section), that the resulting gamma ray will be short lived and pair creation will occur almost immediately within r<sub>H</sub>, since  $\sigma = 1/20$  barn $\approx \pi r_{\rm H}^2$  the cross-sectional equatorial area of the proton, guaranteeing pair creation occurs) replacing the previous pair immediately.



2P3/2 for r=rH New pde solution

Fig1

So this is a *virtual* annihilation-creation process inside  $r_H$ , implying that this two positronsingle electron state is stable (yet another reason for baryon stability). See eq.11 also. We rigorously derive the low mass (<3Gev) hyperon eigenvalues using the Frobenious series solutions to eq.11 near r=r<sub>H</sub> (from 3<sup>rd</sup>pt, r≈r<sub>H</sub>) in Ch.8-Ch.11.

## 7.2 Single Electron Probability Trifolium Statistics Inside r<sub>H</sub>

A single electron in the trifolium implies that on average each of the 3 trifolium lobes has (1/3)e charge (hence the origin of hyperon fractional charge of the lobes). This allows for a toy model in which we give these  $\psi^*\psi$  2P<sub>3/2</sub> at r=r<sub>H</sub> lobes (not particles) names (quarks, the toys.).

## Clebsch Gordon Coefficients For Newpde state: 2P<sub>3/2</sub> at r=r<sub>H</sub>

It is well known that (and also implied by the new pde eq.11 with  $C_M$ ) for the composite system of two electrons  $|1\rangle|2\rangle$  you get, from the analysis of the invariance of the resulting Casimir operator J<sup>2</sup>, the resulting state  $|J_A, J_B, J, M\rangle$  with combined operator  $J_A+J_B=J$ . This is our para state t and 3 ortho states below. For the third spin  $\frac{1}{2}$  particle of far lower energy (the central electron, object B) we have  $|1\rangle|2\rangle|3\rangle$  and so the Clebsch Gordon coefficients imply the decomposition  $(2\otimes 2)\otimes 2=(3\otimes 2)\oplus(1\otimes 2)=4\oplus 2\oplus 2$  so that **three spin 1/2 particles** group together into **four spin** 3/2 and only two spin  $\frac{1}{2}$ s, **6 states** altogether, (the splitting u,d,s,c,b,t). Note then the **majority**  $2P_{3/2}$  (trifolium core) states. Recall also that the  $2P_{3/2}$  solution to new pde at r=r<sub>H</sub> gives a trifolium shape, and  $2P_{3/2}$  fills first. This fills in the broken degeneracy ortho states at the end of part II. The states close to the proton mass are filled in by the Frobenius solution below.

## Eq.2 Single Electron Probability (trifolium statistics):

Consistent with the toy model and also the electron or positron moving between lobes, this time using integer charge distributed over all three  $2P_{3/2}$  lobes at r=r<sub>H</sub>, just randomly put the lobe charges (lobe,lobe) on top of one another Monte Carlo style to determine the probability of a given charge in each lobe. For two positrons [(+1/3,+1/3,+1/3)+(+1/3,+1/3,+1/3)] and one electron (-1/3,-1/3,-1/3) (2P<sub>3/2</sub>) the probability of seeing a +(2/3)e lobe is twice that of seeing a - (1/3)e lobe so ((2/3,2/3,-1/3 or uud proton) eg.,proton, C and b are the <sup>1</sup>/<sub>2</sub> state components of (2 $\otimes$ 2) $\otimes$ 2. For (2P<sub>3/2</sub>) two positrons [(+1/3,+1/3,+1/3)+(+1/3,+1/3)], an electron (-1/3,-1/3,-1/3) and an outlier electron (2P<sub>1/2</sub>) (-1/3,-1/3,-/3) the probability of seeing a -(1/3)e lobe is twice as high as a +(2/3)e lobe so (-1/3,-1/3,2/3) or ddu Neutron).

It is well known that (and also implied by the new pde) for the composite system of two electrons  $|1\rangle|2\rangle$  you get, from the analysis of the invariance of the resulting Casimir operator  $J^2$ , the resulting state  $|J_A, J_B, J, M\rangle$  with combined operator  $J_A+J_B=J$ . Using the resulting Clebsch Gordon coefficients we find the decomposition  $2\otimes 2=3\oplus 1$ , m=1,0,-1 ortho triplet state and singlet para state, which indeed are well known.(eg., Zeeman or Paschen Back line splitting, Ch.8.).

But for a third spin 1/2 particle we have |1>|2>|3> and so the Clebsch Gordon coefficients imply the decomposition  $(2\otimes 2)\otimes 2=(3\otimes 2)\oplus(1\otimes 2)=4\oplus 2\oplus 2$  so that **three spin 1/2 particles** group together into **four spin 3/2** and only two spin  $\frac{1}{2}$  s, **6 states** altogether. Note then the majority  $2P_{3/2}$  (trifolium core) states.

We could now quit and use the mainstream quark (our 2P lobes) applications but that theory is inadequate (eg., neutron-proton binding energy, sect.10.7) and instead will proceed to directly solve equation 2 using the Frobenius series method.

Arfken, *Mathematical Methods of Physics*, 3<sup>rd</sup> ed. Page 454 Enge, Harold, *Introduction To Nuclear Physics*, 1966, Addison Wesley, page. 45

## Newpde Composite **3e** 2P<sub>3/2</sub> at r=r<sub>H</sub> Model

Note that we derived  $\gamma$ ,Zo,W+,W- in Appendix A Of Part I So we derived the E and B fields. **7.3 Eq.11 B Field Flux Quantization In This Enclosed Current Loop** Note if a charged particle moves in loop in a field free region that surrounds another region, there is trapped magnetic flux  $\phi$  in that region. Also we can include minimal interaction E&M momentum/h =k->k+eA/h =eBr/h for uniform B field. If  $\psi$  phase is a unique function on the loop then phase kr= (eBr/h)r=(eBrr/h)= e(Barea)/h=e\Phi/h=n2\pi.  $\Phi$ =4.13X10<sup>-15</sup> for integer spin. Then upon completing a closed loop the particle's wave function will acquire an additional phase factor  $exp\left(\frac{ie\Phi}{h}\right)$ . But the wave function must be single valued at any point in space. This can be accomplished if the magnetic flux  $\Phi$  is quantized:  $e\Phi/h=\pi n$ , n=0±1,±2,±3, so  $\Phi_0$ =h/(2e). From NIST: 2.067833848X10<sup>-15</sup>Wb = $\Phi_0$ . half integer spin<sup>1</sup>/<sub>2</sub>. Integer spin 2 $\Phi_0$ =2(h/2e)=h/e for the two positrons.

#### 7.4 Ultrarelativistic Rotator. Side View

On the side of the rotator  $r_H = R_e \int_0^{\pi/2} \left( \sqrt{1 - c^2 \sin^2 \theta / c^2} \right) \sin \theta d\theta = R_e / 2 = 2.8 \times 10^{-15} \text{m} = r_H \quad (7.1)$ 

for the ortho 2P state observer (i.e,  $2P_{3/2}$ ,  $2P_{\frac{1}{2}}$ ) in the horizontal plane and  $r'_{H}/2=r''$ . We must repeat this integration on the end para states, the radius is shrunk by  $\tau+2(\epsilon+\Delta\epsilon)$  and so is nearly a point source  $S_{\frac{1}{2}}$  state (for the observer above the circle as for the deuterium central electron sect.10.7). We next show that the jump from ortho to para must then correspond to the jump from  $\epsilon$  to  $\tau$  fractal quantum state given  $\tau$  is separable and so a orthogonal state transition.  $2r_{H}=2(2.81X10^{-15}=2e^{2}/(m_{e}c^{2})$ , **Side view**  $\frac{1}{2}(2r_{H})=r_{H}$ , **Top view**  $2r_{H}$ .

#### Inserting Ortho States By Hand Into Flux $\Phi = B_{\perp}A$

Recall from above  $r_H$ = 2.81406X10<sup>-15</sup>m, c=299792458m/s, e=1.602X10<sup>-19</sup>C,  $\mu_o$ =4 $\pi$ X10<sup>-7</sup>T-m/A,  $u_B$ =9.274X10<sup>-24</sup>J/T,

To insert ortho states into the flux by hand assume a circle with another perpendicular coil moving around it with the observer above the circle. The ortho state is put in by hand by dilating the current i coil by 2X as in the *side* view, hence the  $\perp$  in section 7.4. The circle, since it is seen above, is then dilated by the full  $1/\gamma$  Fitzgerald contraction so  $2r_H \rightarrow 2r_H/\gamma$  in BA= $\Phi$ =B( $\pi$ (2 $r_H$ )<sup>2</sup>).

CM= $e^2$  in  $r_H=2e^2/(m_Lc^2)$ . Eq.1.2  $\delta C_M=0$  implies  $\delta r_H=0$  so that  $\delta(2\pi r_H^2)=0$  so minimum area  $r_H$  is given by one of the positrons carrying all the energy  $2\gamma$  seen only in a vertical view (section 7.4) like in deuterium as in section 10.7.

## B Inside 2r<sub>H</sub>

So *top* view of B field and flux (so  $2r_H$ ) and flux *side* view (so  $2r_H$  from eq.7.1). So for top view of B field.

 $2\pi(2r_H)/\Delta t=c$  so (top)  $\Delta t=(2\pi 2r_H)/c=2\pi(2X2.8154X10^{-15})/299792458=1.18X10^{-22}$  sec. Current only observed along 2/3 length for  $2P_{3/2}$  and the time and space  $\gamma$  s cancel. So for the 2e charges  $i=(2/3)2q\gamma/\gamma t=2X(2/3)1.60218X10^{-19}/1.18X10^{-22}=1811$  Amps. So from Ampere's law (top):  $B_1=\gamma\mu_0i/(2(2r_H))=\gamma 4\pi X10^{-7}(1811)/(2X2X2.8154X10^{-15})=\gamma 2.02X10^{11}T=B$  at the center of the big circle. So what are these  $\gamma$  s? It depends on your frame of reference (side or top view). (In section 7.5 below we show they are  $\gamma=\epsilon/\Delta\epsilon$  (side) and  $\gamma=917$  (top).). The average B field inside the 360° loop is just  $B_{av}=2.02X10^{11}T$ .

## 7.5 Consequences Of B Inside $r_{\rm H}$

## **Explains Magnetostars**

So magnetostars are then merely packed neutrons in Ising model  $B_{av}=10^{11}T S_{||}$  viewing alignment caused by a supernova blast with no ortho state  $S_{\perp}$  Meisner effect. But the B field plates in ordinary neutron stars that have slowed down neutron rotation are merely spread out by  $\gamma$  so  $B_{av}/\gamma=B_N\approx10^7T$ . The rapid energy transition  $\int (B^2/2u_o)dV$  between these magnetostar and ordinary neutron star states is that FRB (Fast Radio Burst) energy.

**Boosted B Flux**  $\gamma \Phi = \Phi'$  (sect. 7.0, constant  $\xi_1$ ) Inside side view Area =  $\pi r_{\rm H}^2$ 

 $\Phi = B_{\parallel}A \rightarrow \gamma \Phi = \Phi' = (\gamma BA) = \gamma B(\pi r_{H}^{2}) = (917.2)X2.02X10^{11}X\pi((2.81406X10^{-15})^{2}) = 0$ 

2X2.3X10<sup>-15</sup>Wb $\approx \Phi_{\perp} \approx \Phi_0 = 2Xh/2e$  Boson(2e)  $\Phi_0 = \text{NIST}: 2X2.067833848X10^{-15}Wb = \Phi_0$ So the flux quantization is responsible for the  $\gamma = 917$  and so the proton mass  $2\gamma = (2)917$  (7.4.1) **Para State** S|| The next B flux quantum number (4Xh/2e) is then the next excited state so  $\gamma B = B_{\parallel}$  with then  $\gamma = 917$ .

 $\gamma$ 4B= (917)(4)2.02X10<sup>11</sup>=7.5X10<sup>14</sup>T=B<sub>||</sub>. The 4 is from the Paschen Back Para alignment. The flux quantization also makes para an excited state.

So all views side view (so  $r_H$  for flux loop and B field.). So shrink  $2r_H$  to  $r_H$  (as in sect.7.4 just above due to relativistic observation  $\gamma$  of central electron and increase B by  $\gamma$  in the plates. (side)  $\Delta t=2\pi r_H/c=5.9X10^{-23}$ ,  $i=2(2/3)1.60218X10^{-19}/5.9X10^{-23}=3620A=1$ ; (side)  $\mu_0i/2r_H=8.04X10^{11}T$ . For S\*S=1

 $\gamma$ 8.04X10<sup>11</sup>X=7.5X10<sup>14</sup>T; example: u<sub>B</sub>B4= 9.274X10<sup>-24</sup> 4X7.5X10<sup>14</sup>=2.77X10<sup>-8</sup>J=173GeV in section 7.6 Paschen Back calculation.

# **Ortho State Eq.9.22 Zero Point Energy \epsilon Implies Meisner Effect Nonzero Ortho States** $m=1,0,-1,\gamma_{\epsilon}=\epsilon/\Delta\epsilon$

The magnetic field in one of these protons is about  $10^{11}$ T, so large that any spatially oscillating charge is going to be forced to induce a counter current that tries to cancel the change in flux produced by the charge motion (Faraday's law) relative to the proton. The Frobenius method applied to the new pde has this zero point energy solution eq.9.22 SP hybrid state of the proton whose oscillation provides a Cooper pair oscillation counter current in that huge  $10^{11}$ T field that cancels it out. So at close range there are many pions  $\epsilon/\Delta\epsilon$ . (up to 7) and more distant, where

the B field drops you only need one: Hence the multi pion interactions observed near the proton and resulting Yukawa force.

The ortho state with B orthogonal to A would not exist without this zero point SP state motion since the (SP hybrid so) induced P state spinor has a horizontal component so has dot product with horizontal  $S_2$  nonzero spinor(for ortho).

These two B fields (B<sub>1</sub> and B<sub>1</sub>) are put into Paschen Back (eg.,m<sub>s</sub>c<sup>2</sup>=u<sub>B</sub> $\gamma_e$ B(1+0+0+0))



## **2** body positron-positron dynamics so ortho(s.c.b) and para(t) states on the $r=r_H$ shell near $\theta=90^\circ$ . Paschen Back dynamics

Summary Of Para and Ortho States: two  $\gamma$ s: top view Para  $\gamma$ =917, Ortho side view  $\gamma = \epsilon/\Delta \epsilon$ S<sub>1</sub> SP hybrid zero point energy eq.9.22: P state so S<sub>1</sub>•S<sub>2</sub>  $\neq$ 0, ortho:  $\gamma_{\epsilon m=\pm 1,0}$ B=( $\epsilon/\Delta \epsilon$ )B=B<sub>1</sub>. =(.06/.00058)2.02X10<sup>11</sup>=2.27X10<sup>13</sup>T=B<sub>1</sub> (eg., m<sub>s</sub>=u<sub>B</sub>B<sub>1</sub>(1+0+0+0))

 $\begin{array}{ll} S_{\parallel} \ \ Flux \ quantization \ \gamma \ top \ S_{\parallel} \ view \ both \ positrons \ so \ S_1 \bullet S_2 \neq 0, \ para: \ \gamma_{L=0}=\gamma B=B_{\parallel} \ . \\ = 917(4)2.02X10^{11}=7.5X10^{14}T=B_{\parallel} \qquad (eg., \ m_t=u_BB_{\parallel}(1+1+1+1)) \end{array}$ 

7.6 The two rapidly moving positrons:  $2 \otimes 2=3 \oplus 1$  List of 1 Para SB<sub>||</sub> and 3 Ortho SB<sub>\</sub> States Here Thomas LS LST= -(L<sub>BL</sub>\*(S<sub>AL</sub> or S<sub>CL</sub>))K±gs perturbation is subtracted off the Paschen Back energy for both the SB<sub>\</sub> and SB<sub>||</sub> cases.

**SB**<sub>I</sub>=State t **non** LS coupling **para** singlet state (the 1 single state in the  $3\oplus 1$  decomposition)  $B_{\parallel}=7.448X10^{15}T$ ,= $B=\gamma(2.02X10^{11})$ **0**° rн Pauli Principle. LEM B<sub>||</sub>u<sub>B</sub>(mLA+mSA+mLC+mSC)=PE LST PE-LST name S L=0 1 + 1 + 1 + 1 even stable Singlet 173 0 173 t para SB<sub>1</sub> State B total triplet b,c,s ground state u/d LS coupling triplet ortho state. LS coupling  $B_{\perp} \approx 4.043 \times 10^{12} \text{T}$  ( $\epsilon/\Delta \epsilon$ )2.02X10<sup>11</sup>=2.27X10<sup>13</sup>T **90°**  $S_B = \pm 1$  $r_{snf} 2P_{\frac{1}{2}} = at r = r_{H}$ . Here single  $PE = \frac{1}{2} PE = \frac{1}{2} D \text{ bond} = \frac{1}{2}$ ; D=2. See sect. 10.7  $B_{\perp}u_{B}(mLA+mSA+mLC+mSC)=PE$ LST ≈LST-PE name Pauli Principle m=1 1 + 1 + 1 + 15790 1.5+1(2) $\Xi_{\rm b}$ ortho

ortho	$\Xi_{\rm s}$	1.5-0 (2)	2471	1 + 1 + 0 + 0	m=0
ortho	$\Xi_{\rm s}$	1.5-1 (2)	1314	1 + 1 + 0 + 0	m=-1
$2P_{3/2} \& 2P_{1/2}$	$P_{u/d}$	1	=(1)938	Ground State P <sub>u/d</sub> =	

So a total of 4 states for two positrons (3ortho, 1para). 6  $2P_{3/2}$  states if you include the central electron. Since the proton is the core object for these states we can use the Frobenius solution Ch.9 perturbations below for these r>r<sub>H</sub> deviations from the spin 1 flux quantization  $2\phi_E=2h/2e$  above sect.10.13 and  $\Xi$ .. We get four multiplets of the three  $\Xi$  one P. Get ud. (Chapter 8). The above are also boson energy transitions analogous to the principle quantum number photon transition emissions of the hydrogen atom.

## **Other Ortho Consequences**

We can reverse engineer this process by modeling a large decrease in the resulting strong magnetic field:

Neutron  $2P_{\frac{1}{2}}$  1-r<sub>H</sub>/r<sub>H</sub> for charge 0 (case 2 Ch.8) is homeomorphically mapped into 1- $\varepsilon$  with added outside particle KE Meisner effect additional outside charge (reducing that r<sub>H</sub>/r<sub>H</sub> charge so preserving angular momentum a (and so KMQ) in the Kerr metric term (a/r)<sup>2</sup>. Note the negative sign still indicates inside multibody charge is still 0.

Proton  $2P_{3/2}$  1+r<sub>H</sub>/r<sub>H</sub> (case 1 in Ch.8) is then homeomorphically mapped into 1+ $\epsilon$  with added particle KE. The positive sign indicates nonzero internal multibody charge. See eq.B4. For  $2P_{3/2}$   $\kappa_{00}$ =(1-2 $\epsilon$ )- $\Delta\epsilon$ -[(C<sub>M</sub>/m<sub>e</sub>)r]. The starting point of PartII. (B3) vector  $\omega_p$ . So ultrarelativistic Thomas precession =n $3m_p$  =LS energy  $\omega$  is subtracted off from Paschen Back energy  $\omega_p$ . It adds to 0 in the ground state. Iso L<sub>EM</sub>=Nr<sub>H</sub> $\pm\omega$ /c=rX[(EXH)/c]( $\pi$ r<sub>H</sub><sup>2</sup>T)

angular momentum also cancels some of the total angular momentum of objects A,C and B. For each ortho state we apply the Frobenius solution perturbation(Ch.9). The next ortho value is

 $m_p=2$  (for s and later  $m_p=4$  for c,  $m_p=6$  for the b state) for the next ortho state.

## Calculation Of $\xi_1$

We use the equation 1.2.7 energy normalization (m<sub>c</sub>=1) for two reduced mass 2P3/2 ultrarelativistic positrons at r=r<sub>H</sub> with ansatz  $\xi \rightarrow x_2$ , in  $\xi_0 \rightarrow 1$  in  $\xi_1 = \xi_3 + \xi_2 + \xi_0$ . So  $E = C_M \xi_1^2 / \xi \sqrt{2} \rightarrow \frac{1}{2} \xi_1^2 / \xi_1 = \frac{1}{2} / (x2 + (2+\Delta)x + (1+\Delta)) =$  (partial fractions) =  $\frac{1}{2} ((1(-1/\Delta)/(x+1)) + ((1(1/\Delta)/(x+1+\Delta)) = positron1 + positron2$ . So for x→0 then  $\Delta = 1/3684$  from the boosted magnetic flux calculation  $2\gamma = 3684$ .

# 2AI+2AI+2AI Paschen Back



## **Fig.8 Chapter 8 Frobenius Series Perturbation To Each Paschen Back State Introduction**

Here we start with the ground state magnetic flux energy(u/d) set  $m_b=1$ , move on to the three orthos (s,c,b) with larger  $m_b s$  ( $\Xi$ ) and finally to the very high para (t). We are actually perturbing the motions at  $r_H$  by these r in equation 9.5 and so are taking into account the constituents of the proton in this way.

Also there are then 6 magnetic flux quantization  $2P_{3/2}$  states. Each flux quantization level has its own m<sub>p</sub> and associated Frobenius solution. So we have ground state m<sub>p</sub>=1, and excited states: m<sub>p</sub>=1.5= $\Xi_s$ , and also  $\Xi_c$ ,  $\Xi_b$ , each having it's own Frobenius solution sets.

## 8.1 Solution to eq.2 Using Separability: Gyromagnetic Ratios And Low Energy Particles (energy<3GeV) Derived For ground state Magnetic Flux r≈r<sub>H</sub> Application: Gyromagnetic Ratios

After separation of variables the "r" component of equation 9 can be rewritten as:

$$\left[\left(\frac{dt}{ds}\sqrt{g_{oo}}\,\boldsymbol{m}_{p}\right)+\boldsymbol{m}_{p}\right]F-\hbar c\left(\sqrt{g_{rr}}\,\frac{d}{dr}+\frac{j+3/2}{r}\right)f=0$$
(8.1)

$$\left[\left(\frac{dt}{ds}\sqrt{g_{00}}m_p\right) - m_p\right]f + \hbar c \left(\sqrt{g_{rr}}\frac{d}{dr} - \frac{j-1/2}{r}\right)F = 0$$
(8.2)

Because the  $\kappa_{00} = 1 - r_H/r$  is point source the object B ambient metric is local and so the vacuum is not infinite density (see also sect 6.11) as in the QED ambient metric which is homogenous. Comparing the flat space-time Dirac equation to equations 8.1 and 8.2:

$$(dt/ds)\sqrt{\kappa_{oo}} = (1/\kappa_{00})\sqrt{\kappa_{oo}} = (1/\sqrt{\kappa_{oo}}) = \text{Energy} = E$$
 (8.2a)

Using the above Dirac equation it is easiest to find the gyromagnetic ratios gy for the spin polarized F=0 case. Recall the usual calculation of rate of the change of spin S gives  $dS/dt \propto m \propto gyJ$  from the Heisenberg equations of motion. We note that  $1/\sqrt{g_{rr}}$  rescales dr in  $\left(\sqrt{g_{rr}} \frac{d}{dr} + \frac{J+3/2}{r}\right)f$  in equation 8.1. Thus to have the same rescaling of r in the second term we must multiply the second term denominator (i.e., r) and numerator (i.e., J+3/2) each by  $1/\sqrt{g_{rr}}$  and set the numerator equal to 3/2+J(gy), where gy is now the gyromagnetic ratio. This makes

our equation 8.1 compatible with the standard Dirac equation allowing us to substitute the gy

Thus again:

$$[1/\sqrt{g_{rr}}](3/2+J)=3/2+Jgy$$
, Therefore for  $J=\frac{1}{2}$  we have:  
 $[1/\sqrt{g_{rr}}](3/2+\frac{1}{2})=3/2+\frac{1}{2}gy=3/2+\frac{1}{2}(1+\Delta gy)$  (8.3)

Then we solve equation 8.3 for gy and substitute it into the above dS/dt equation.

into the standard dS/dt $\propto$ m $\propto$ gyJ to find the correction to dS/dt.

**S** States: Recall  $\varepsilon$  and  $\Delta \varepsilon$  and S states from eq. 6.4.13. These are zero point energy states (eq.9.22) that must also be the source of the Meisner effect canceling of those large B fields. Noting in equation 6.4.13 we get the gyromagnetic ratio of the electron with  $g_{rr}=1/(1+\Delta \varepsilon/(1+\varepsilon))$  and  $\varepsilon=0$  for electron. Thus solve equation 8.3 for  $\sqrt{g_{rr}}=\sqrt{(1+\Delta \varepsilon/(1+\varepsilon))}=\sqrt{(1+\Delta \varepsilon/(1+0))}=\sqrt{(1+\Delta \varepsilon/(1+0))}=\sqrt{($ 

 $[1/\sqrt{(1+.0005799)}](3/2 + \frac{1}{2}) = 3/2 + \frac{1}{2}(1+\Delta gy)$ . Solving for  $\Delta gy$  gives anomalous gyromagnetic ratio correction of the electron  $\Delta gy=.00116$ 

Going to higher energies (so  $\epsilon \neq 0$  in equation 8.3) we get the anomalous gyromagnetic ratio correction of the muon. From the momentum representation of eq.8.1,8.2:

**2P**<sub>3/2</sub> **states:** Recall the 2P3/2 states from chapter 3. Note also that k can be positive or negative since  $4\pi k=Z_{00}$  in our Lagrangian with a positive k meaning at least one charge is not canceled. Therefore  $1/g_{rr} = 1 \pm k/r + \epsilon$  (using our Frobenius solution expansion near  $r \approx r_H$  of eq.9.5 below multiply through by  $(1+\epsilon/4)((1+\epsilon+...) \approx 1+.08=1+\epsilon)$  so a pion mass is then added to the protons) from the  $\pm$ nature of  $Z_{00}$ . Therefore we have two cases from equation B3 at the boundary r=k

CASE 1	$1/g_{rr} = 1 + k/k + \epsilon$	charge 1	(core case)
CASE 2	$1/g_{rr} = 1 - k/k + \epsilon$	charge 0	(use m from case 1)

**Note:**  $\epsilon$  (9.22) is required because it is the zpe here (like  $\frac{1}{2}\omega^{1/2}$  is the zpe of 1D SHM) external to the 3e region. So through the Faraday's law Meisner effect pops up to cancel that huge  $10^{14}$ T internal B field, hence the origin of the mesonic field. So the  $\epsilon$  in case 1 and case II is the artifact of that large internal B field of section 8.1.

Also the effect of a zero charge is to make metric component  $g_{oo}$  (=1/g<sub>rr</sub>) contribution zero in case 2. Thus the effect of *non*zero charge is to increase the dimensionality by adding a metric component in eq.2. This provides the reason that Kaluza Klein theory (adding a 5<sup>th</sup> dimension) is so successful at injecting E&M into general relativity. But Kaluza Klein theory is not required here because finite C<sub>M</sub> in eq.1.11 is really responsible for charge and E&M. 2D is sufficient as

we showed in Chapter 1, eq.1.5. The extra 2D degree of freedom is associated with that extra real term  $\delta\delta z$  in the amazing equation 1.6.

CASE 1: Plus +k, therefore is the proton + charge component.  $1/g_{rr} = 1+k/k + \epsilon = 2 + \epsilon$ . Thus from equation 8.1, 8.2  $\sqrt{2 + \epsilon} (1.5 + .5) = 1.5 + .5(gy)$ , gy=2.8 (8.4)

The gyromagnetic ratio of the proton (therefore that above  $r \approx k$  stability was indeed proton stability as we concluded) mass=m<sub>p</sub>. dt/ds $\sqrt{g_{oo}} = 1/\sqrt{g_{oo}} = E = m_p$ 

CASE 2: negative k, thus charge cancels, zero charge:

 $1/g_{rr} = 1-k/k + \epsilon = \epsilon$  Therefore from equation 8.3 and case  $1 1/g_{rr} = 1+k/k+\epsilon$ 

 $\sqrt{\varepsilon} (1.5+.5)=1.5+.5(gy), gy=-1.9,$  (8.5)

the **gyromagnetic ratio of the neutron** with the other charged and neutral hyperon magnetic moments scaled using their masses by these values respectively.

### Chapter 9

## **The 3e Energies For particle energy <3GeV Derived Using Frobenius Series Solution** (at first Paschen Back m<sub>p</sub> level) Perturbation)

#### 9.1 Series Solutions y Ansatz Near r≈r<sub>H</sub>

 $m_p=1$  here.  $m_p$  determined from Paschen Back energy level (next Paschen Back level  $m_p=2$  for s) Recall equations 8.1, 8.2:

$$\begin{bmatrix} \left(\frac{dt}{ds}\sqrt{g_{oo}} \,\boldsymbol{m}_{p}\right) + \boldsymbol{m}_{p} \end{bmatrix} F - \hbar c \left(\sqrt{g_{rr}} \,\frac{d}{dr} + \frac{j+3/2}{r}\right) f = 0$$
$$\begin{bmatrix} \left(\frac{dt}{ds}\sqrt{g_{00}} \,\boldsymbol{m}_{p}\right) - \boldsymbol{m}_{p} \end{bmatrix} f + \hbar c \left(\sqrt{g_{rr}} \,\frac{d}{dr} - \frac{j-1/2}{r}\right) F = 0$$

Recall from the previous section  $g_{oo}=1-k/r-(\epsilon+\Delta\epsilon)$ . Also **recall our Dirac doublet** (equation 2AI ) must have a left handed zero mass component will be called case 1 and case 3 respectively below. Also we need the equivalent of the singlet equation 2 is our below case 2. Also in equation 2 at  $r=r_H$  the eigenvalue is  $\Delta\epsilon+\epsilon+1=2m_p$  for that principle quantum which then must be the same for the  $2P_{3/2}$  state. Here we write out the left handed Dirac Doublet Eq.2 in the general representation of the Dirac matrices. Also recall from chapter 8 that the  $2P_{3/2}$  state (and its sp<sup>2</sup> hybrid) for this new electron Dirac equation gives a azimuthal trifolium, 3 lobe shape and thus a  $\lambda/3$  spherical harmonic wavelength so that for covalent bonding r' $\approx r_H/3$  in  $\kappa_{oo}=1-r'/r$ . This  $\lambda/3$  also is used to calculate P wave scattering (called "jets" by quark people.)

To use the f & F components of the equation 8.1, 8.2 Dirac equation we write the Dirac equation for free particle motion along the symmetry axis z (r=ratio of momentum to energy) to find the chirality of the components in the general representation of section 1.6. We then compare this z motion free particle Dirac equation eigenfunction structure with radial component structure to arrive at a sense of which components of the radial equation are left handed and which aren't. This step is a little more complicated here because we are not using the chiral representation of the Dirac matrices, but the standard representation instead. In any case given that the electron is positive energy, then (as we see below) for the positron -E gives left handed f and F implying that this object *must* have a positive charge since this left handedness(doublet, Ch.3) results from the fractalness (There is a corresponding argument for G and g). The proton indeed is positive charged. So:

$$\begin{bmatrix} -\left(\frac{dt}{ds}\sqrt{g_{oo}} \ m_p\right) + m_p \end{bmatrix} g - \hbar c \left(\sqrt{g_{rr}} \ \frac{d}{dr} + \frac{j+3/2}{r}\right) G = 0 \rightarrow \mu c^2 u_1 + cpu_3 - Epu_1$$

$$\begin{bmatrix} -\left(\frac{dt}{ds}\sqrt{g_{00}} \ m_p\right) - m_p \end{bmatrix} G + \hbar c \left(\sqrt{g_{rr}} \ \frac{d}{dr} - \frac{j-1/2}{r}\right) g = 0 \rightarrow cpu_1 - \mu c^2 u_3 - Epu_3$$

$$\begin{bmatrix} -\left(\frac{dt}{ds}\sqrt{g_{oo}} \ m_p\right) + m_p \end{bmatrix} f - \hbar c \left(\sqrt{g_{rr}} \ \frac{d}{dr} - \frac{j-1/2}{r}\right) F = 0 \rightarrow \mu c^2 u_2 - cpu_4 - Epu_2$$

$$\begin{bmatrix} -\left(\frac{dt}{ds}\sqrt{g_{00}} \ m_p\right) - m_p \end{bmatrix} F + \hbar c \left(\sqrt{g_{rr}} \ \frac{d}{dr} + \frac{j+3/2}{r}\right) f = 0 \rightarrow -cpu_2 - \mu c^2 u_4 - Epu_4$$
where to get correspondence from these two Direc equation structures we see that at +

where to get correspondence from these two Dirac equation structures we see that at +E:  $u^R=$ 

$$\begin{pmatrix} 1\\0\\r\\0 \end{pmatrix} = g, \ u^{L} = \begin{pmatrix} 0\\1\\0\\-r \end{pmatrix} = f; \ -E: \text{ No } (v^{R} = \begin{pmatrix} r\\0\\1\\0 \end{pmatrix} \text{ here}), \ v^{L} = \begin{pmatrix} 0\\r\\0\\1 \end{pmatrix} = F, \text{ Note in general (with r=0) here:}$$

$$\begin{pmatrix} ig\\if\\G\\F \end{pmatrix} = \begin{pmatrix} u^{R}\\u^{L}\\v^{R}\\v^{L} \end{pmatrix} = \Psi. \text{ So we have the solution that in the standard representation of the left handed}$$

doublet is given by F and f only for –E of the electron (here a positron needed below for + proton hadron excited states) at the horizon. Dirac matrices

So for the left handed doublet: 
$$\binom{F}{f}_{L}$$
 we have respectively  $\begin{array}{c} q\pm, m \neq 0, -E, forF\\ q=0, m=0, +E, forf \end{array}$  (9.4)

Or more succinctly equation 2 in the Dirac doublet form implies in section B2: Note our postulate implies  $C \rightarrow 0$  so we are on the dr axis thus dt'=0 so dt'<sup>2</sup>=(1-r<sub>H</sub>/r)dt<sup>2</sup> (sect.0.1 of Ch.1). Thus r=r<sub>H</sub>=k is a stable point since the clock stops since dt'=0 and the is the Meisner effect formalism for canceling out that huge B field at a distance and a;so making it so the protons mass is mp, and not much larger.

Note in sect.1.5 we can have a zero and nonzero charge in the 3<sup>rd</sup> quadrant (where dt=dr) massive Proca boson case given the possibilities in sign we have for  $\pm\epsilon'/2$  in  $((-\epsilon/2) \pm\epsilon'/2)dr \cdot ((-\epsilon/2)\pm\epsilon'/2)dt$ . In the first quadrant ds=0 again (section 1.4) so they have to add to zero.  $+dr+\epsilon/2+dt-\epsilon/2$  and  $-dr-\epsilon/2-dt+\epsilon/2$  solutions. Multiply the second equation by -1, then add the two resulting equations, then divide by 2 and get  $dr+\epsilon/4\pm\epsilon/4+dt-\epsilon/4\pm\epsilon/4$  so that  $\epsilon/2\rightarrow\epsilon/2\pm\epsilon/2$ . So we multiply each of the two ds<sup>2</sup> cases (above |dr+dt| discussion) by its own dz, each with its own  $\kappa_{rr}=1/(1-\epsilon//r) \rightarrow 1/(1-(\epsilon/2\pm\epsilon/2)/r)$  (sect.4.7) implying 2 charges  $\epsilon/2-\epsilon/2=0$ ,  $\epsilon/2+\epsilon/2=\epsilon$  and so two Proca equation massive W,Z.

See B2. .See above B2:

CASE 1	$1/g_{rr} = 1 + k/k + \varepsilon$	F	charge 1, m=1 (core case) 2P3/2
CASE 2	$1/g_{rr} = 1 - k/k + \epsilon$	F	charge 0, m from case 1) $2P1/2$
CASE 3		f	charge 0, m=0

We solve these equations only near  $r\approx_{k_{\rm H}}$  since that is where the stability is to be found (and also fortunately were these equations are *linear* differential equations). Thus our first step is to expand  $\sqrt{g_{rr}}$  about this radius and drop the higher order terms.

The Frobenius series solution method can now be used to solve equations 8.1 and 8.2 at  $r\approx r_H$ . See for example Mathematical Methods of Physics, Arfken 3<sup>rd</sup> ed. Page 454. First we solve the f in equation 8.1, plug that into equation 8.2 and then have an equation in only F. There we substitute a series solution ansatz  $F = \sum a_n r^n$  in the resulting combined equations. We can then separate out the results into coefficients of respective rn and get recursion relations that will give us series that must be terminated at some N. Note the energy Eigenvalue 'E' will be in this series as  $dt/ds\sqrt{g_{oo}}$  so we can then solve for the mass energy of these hadrons at specific J. We will need an indicial equation for the first term to start out this process. Also in this Frobenius solution method 'n' turns out to be a multiple of  $\frac{1}{2}$  and the series must start at n=-1. Finally to get the charge zero case the charged case must be done first and its constant masses used in the uncharged state calculations.

#### 9.2 CASE 1 Excited States for F, m≠0, q± 2P3/2

Again case 1 is one of the equation 8.1 possibilities. Therefore let  $R=k_H-r$ , r<<R (for stability) we can write in 8.1:

$$\sqrt{g_{rr}} = 1/\sqrt{1 + k_H/R + \varepsilon} \approx \frac{\sqrt{R}}{\sqrt{R + k_H + R\varepsilon}} =$$
(9.1)

$$\frac{\sqrt{k_H - r}}{\sqrt{k_H - r + k_H + (k_H - r)\epsilon}} =$$
(9.2)

$$\frac{\sqrt{k_H - r}}{\sqrt{k_H (2 + \varepsilon) - r(1 + \varepsilon)}} =$$
(9.3)

$$\frac{(1-\varepsilon/4)}{\sqrt{2}} \left( \frac{1-\frac{r}{2k_{H}} + \frac{r^{2}}{8k_{H}^{2}} + ..}{1-\frac{r}{4k_{H}} + \frac{r^{2}}{16k_{H}^{2}} + ..}}{1-\frac{r}{4k_{H}} + \frac{3r^{2}}{32k_{H}^{2}} - ..} \approx \frac{1-\frac{r}{4k_{H}}}{\sqrt{2}} \right)$$
(9.4)

Note taking the first term of this Taylor expansion of the square root makes this an approximation (<2GeV.). Note that including the above  $1\pm\epsilon/4$  the compensating  $(1\pm\epsilon/4)$  in the next r term has the effect of a multiplying the derivative terms by  $1\pm\epsilon/4$ . This rescales r to allow us to still say that the stable boundary is still at r<sub>H</sub>. Thus we could use it to also rescale t in the first term of equations 8.1 and 8.2 or note that  $(1+\epsilon/4) (1+\epsilon)=1+5/4\epsilon$  thus renormalizing  $1+\epsilon$  to  $1+4/3\epsilon=1+\epsilon$ ' everywhere. Also the  $3r^2/32k_H^2$  terms must be included. We drop these perturbative terms until the end. Therefore substituting in equation 9.5 we find that equation 8.1 reads:

$$\begin{bmatrix} E+m_{p} \end{bmatrix} F -\hbar c \left( (1-r/k_{H}4) \frac{d}{\sqrt{2}dr} + \frac{j+\frac{3}{2}}{k_{H}-r} \right) f = 0 =$$
(9.6)  
$$\begin{bmatrix} E+m_{p} \end{bmatrix} F -\hbar c \left( (1-r/k_{H}4) \frac{d}{\sqrt{2}dr} + (1+r/k_{H})(j+\frac{3}{2})/k_{H} \right) f = 0; \text{ also:}$$
$$\begin{bmatrix} E-m_{p} \end{bmatrix} f +\hbar c \left( \sqrt{g_{rr}} \frac{d}{dr} - \frac{j-\frac{1}{2}}{r} \right) F = 0$$
$$\begin{bmatrix} E-m_{p} \end{bmatrix} f +\hbar c \left( (1-r/k_{H}4) \frac{d}{\sqrt{2}dr} - (1+r/k_{H}) \left( j-\frac{1}{2} \right)/k_{H} \right) F = 0$$
Therefore  
$$f = -\hbar c \frac{\hbar c}{\left[ E-m_{p} \right]} \left( (1-r/k_{H}4) \frac{d}{\sqrt{2}dr} - (1+r/k_{H}) \left( j-\frac{1}{2} \right)/k_{H} \right) F \text{ substituting into}$$

$$\left[E + m_{p}\right]F - \hbar c \left(\left(1 - r/k_{H} 4\right)\frac{d}{\sqrt{2}dr} + \frac{j + \frac{3}{2}}{k_{H} - r}\right)f = 0$$
(9.7)

We find solving for f and substituting back in:

$$\begin{split} & \left[E + m_{p}\right]F - \hbar c \bigg( (1 - r/k_{H} 4) \frac{d}{\sqrt{2}dr} + (1 + r/k_{H})(j + 1.5)/k_{H} \bigg) \bullet \\ & \frac{\hbar c}{E - m_{p}} \bigg( - (1 - r/k_{H} 4) \frac{d}{\sqrt{2}dr} + (1 + r/k_{H})(j - \frac{1}{2})/k_{H} \bigg)F = [E + m_{p}]F + \\ & \frac{(\hbar c)^{2}}{(E - m_{p})\sqrt{2}} \bigg( - (1 - r/4k_{H}) \frac{d}{k_{H} 4\sqrt{2}dr} + (1 - r/4k_{H})^{2} \frac{d^{2}}{\sqrt{2}dr^{2}} - (1 - r/4k_{H})(j - \frac{1}{2})/k_{H}^{2} \bigg)F \\ & + \frac{(\hbar c)^{2}}{E - m_{p}} \bigg( (1 + 3r/k_{H} 4)(j + 1.5) \frac{d}{\sqrt{2}k_{H}dr} - (1 + r/k_{H})^{2}(j + 1.5)(j - \frac{1}{2})/k_{H}^{2} \bigg)F = \\ & \bigg( [E + m_{p}] + \bigg[ \frac{(\hbar c)^{2}}{(E - m_{p})} \bigg( - \bigg( j + \frac{3}{2} \bigg)(j - \frac{1}{2})/k_{H}^{2} + (j - \frac{1}{2})/\sqrt{2}k_{H}^{2} \bigg) \bigg] \bigg)F + \\ & \frac{(\hbar c)^{2}}{(E - m_{p})\sqrt{2}} \bigg( - 2\sqrt{2}\bigg( j + \frac{3}{2} \bigg)(j - \frac{1}{2})/k_{H}^{3} + (j - \frac{1}{2})/4k_{H}^{3} \bigg)rF + \\ & \frac{(\hbar c)^{2}}{(E - m_{p})\sqrt{2}} \bigg( \frac{-1}{k_{H} 4\sqrt{2}} + \bigg( j + \frac{3}{2} \bigg) \frac{1}{k_{H}} \bigg) \frac{dF}{dr} + \\ & \frac{(\hbar c)^{2}}{(E - m_{p})\sqrt{2}} \bigg( \frac{1}{k_{H}^{2} 16\sqrt{2}} + \bigg( j + \frac{3}{2} \bigg) \frac{3}{4k_{H}^{2}} \bigg)r \frac{dF}{dr} + \\ & \frac{(\hbar c)^{2}}{(E - m_{p})\sqrt{2}} \bigg( \frac{1}{\sqrt{2}} \bigg) \frac{d^{2}F}{dr^{2}} + \\ & \frac{(\hbar c)^{2}}{(E - m_{p})\sqrt{2}} \bigg( \frac{-1}{2\sqrt{2}k_{H}} \bigg)r \frac{d^{2}F}{dr^{2}} + \\ & \frac{(\hbar c)^{2}}{(E - m_{p})\sqrt{2}} \bigg( \frac{-1}{2\sqrt{2}k_{H}} \bigg)r \frac{d^{2}F}{dr^{2}} \end{split}$$

Here  $r=2k_H$  is a regular singular point. Next substitute in  $F=\sum_n a_n r^n$  with again half integer n allowed as well:

$$\sum_{M}^{N} \left[ \left[ E + m_{p} \right] + \left[ \frac{(\hbar c)^{2}}{(E - m_{p})} \left( - \left( j + \frac{3}{2} \right) (j - \frac{1}{2}) / k_{H}^{2} - (j - \frac{1}{2}) / \sqrt{2} k_{H}^{2} \right) \right] \right] a_{n} r^{n} + \qquad (9.8)$$

$$\sum_{M}^{N} \frac{(\hbar c)^{2}}{(E - m_{p})\sqrt{2}} \left( -2\sqrt{2} \left( j + \frac{3}{2} \right) (j - \frac{1}{2}) / k_{H}^{3} + (j - \frac{1}{2}) / 4k_{H}^{3} \right) a_{n-1} r^{n} +$$
(9.9)

$$\sum_{M}^{N} \frac{(\hbar c)^{2}}{(E - m_{p})\sqrt{2}} \left( \frac{-1}{k_{H} 4\sqrt{2}} + \left( j + \frac{3}{2} \right) \frac{1}{k_{H}} \right) (n+1)a_{n+1}r^{n} +$$
(9.10)

$$\sum_{M}^{N} \frac{(\hbar c)^{2}}{(E - m_{p})\sqrt{2}} \left( \frac{1}{k_{H}^{2} 16\sqrt{2}} + \left(j + \frac{3}{2}\right) \frac{3}{4k_{H}^{2}} \right) n a_{n} r^{n} +$$
(9.11)

$$\sum_{M}^{N} \frac{(\hbar c)^{2}}{(E - m_{p})\sqrt{2}} \left(\frac{1}{\sqrt{2}}\right) (n+2)(n+1)a_{n+2}r^{n} +$$
(9.12)

$$\sum_{M}^{N} \frac{(\hbar c)^{2}}{(E - m_{p})\sqrt{2}} \left(\frac{-1}{2\sqrt{2}k_{H}}\right) (n+1) n a_{n+1} r^{n} = 0$$
(9.13)

Note from equation 9.12 that this series diverges. To terminate the series we now take 9.8 and 9.11 together and 9.10 and 9.13 together (since they have the same  $a_n$ ). For example combining the equation 9.8 and 9.11 terms

$$\left( \left[ E + m_p \right] + \left[ \frac{(\hbar c)^2}{(E - m_p)} \left( - \left( j + \frac{3}{2} \right) (j - \frac{1}{2}) / k_H^2 - (j - \frac{1}{2}) / \sqrt{2} k_H^2 \right) \right] \right) + \frac{(\hbar c)^2}{(E - m_p)\sqrt{2}} \left( \frac{1}{k_H^2 16\sqrt{2}} + \left( j + \frac{3}{2} \right) \frac{3}{4k_H^2} \right) n,$$

Replacing the normalization  $m_p \rightarrow m_p(1 \pm \epsilon)$  (from section 4.8):

$$\left(E^{2}-m_{p}^{2}\right)+\left(\left[\left(-\left(j+\frac{3}{2}\right)(j-\frac{1}{2})/k_{H}^{2}-(j-\frac{1}{2})/\sqrt{2}k_{H}^{2}\right)\right]\right)-\frac{1}{\sqrt{2}}\left(\frac{1}{k_{H}^{2}16\sqrt{2}}+\left(j+\frac{3}{2}\right)\frac{3}{4k_{H}^{2}}\right)N=0$$

Therefore after rearranging:

$$E = \sqrt{m_p^2 + \frac{1}{k_H^2} \left(j^2 + 1.7071j - 1.10355 - (j.5303 + .8269)N\right)}, \qquad (9.14)$$

We have for a general Laurent series ansatz:

 $..+a_{-1}r^{-1}+a_{-1/2}r^{-1/2}+a_{o}r^{o}+a_{1/2}r^{1/2}+a_{1}r^{1}+..=F$ 

Note also that equations 9.8-9.13 imply that the coefficients of a given  $r^n$  are independent. Thus adding together the coefficients of  $r^n$  for equations 9.8–9.13 at a given n:

 $9.9(j_{-1/2})a_{n-1} + (9.8+9.11)a_n + (9.10+9.13)(n+1)a_{n+1} + 9.12(n+2)(n+1)a_{n+2} = 0$  (9.15)

#### **Method of Solving Equation 9.15**

For the outside observer an F=0 finite boundary condition at infinity applies for flat vacuum value n=0, j=1/2 and for r<sup>o</sup>, r<sup>-1/2</sup>, r<sup>-1</sup> and for complete vacuum for N=0, J=0. Here then the generalized Laurent series  $..+a_{-1}r^{-1}+a_{-1/2}r^{-1/2}+a_{o}r^{o}+a_{1/2}r^{1/2}+a_{1}r^{1}+..=F$  reduces to  $..+a_{-1}r^{-1}+a_{-1/2}r^{-1/2}+a_{o}r^{o}=F$ . Thus either set 9.9(j-1/2)a<sub>n-1</sub>=0 or (9.10+9.13)(n+1)a<sub>n-1</sub>+9.12(n+1)(n+1)a\_{n+2}=0 separately in eq.9.15 or set both equal to zero:

#### $J = \frac{1}{2}$ , sets eq.9.9=0

1) N=-1, in equation 9.14 gives mass eigenvalue for  $\Xi$ 

Exact solution for all possible a<sub>n</sub>, sets none of them to zero.

2) N=0, in equation 9.14 gives mass eigenvalue for *nucleon*. dr<sup>o</sup>/dr=0 so all derivative of F terms are then zero and this solution applies inside as well.

N=0 flat J=0 allowed flat vacuum gives  $\pi^{\pm}$  and with free e, j=  $\frac{1}{2}$  muon.

3) N= -1/2, in equation 9.14 gives mass eigenvalue of two  $\Sigma$  s since a plus *and* minus square root of r.

These 9.9=0 cases have case 2 zero charge representations as well.

#### N=-1, Principle QM number Also a<sub>-2</sub> =0

- 1) J=0, in equation 9.14 gives mass eigenvalue for K
- 2) J=1, gives deuteron mass eigenvalue (bonding) given N=0,J=0 fills first (i.e., pion). Thereafter use nuclear shell model-Schrodinger equation many body techniques with these nonrelativistic lobes with this (bound state) force acting like a outer layer surface tension, finite height square well potential . Get a aufbau principle that then gives the D,F,G,..nuclear shell model states. Alternatively can fill that first S state in with free  $1S_{\frac{1}{2}}$  (next state to filled state) and we have j=3/2  $\Omega^{-}$  filling in some (i.e., uds) of the  $2P_{3/2}$  states (see Ch.9) and thereby also deriving from first principles Gell Man's 1963 eight fold way for hyperon eigenvalue classification (to finish that effort need case II zero charge and case III  $\Lambda_0$  as well).  $M_p$  is replaced by 2 in c hyperons, by 4 for b hyperons as indicated in fig. 16-1 for how to fill in the cbt 2P harmonic states given the requirement to use r<sup>2</sup> then.

Also, to include higher order r expansion term effects in equation 9.5 we must include those perturbative  $1+\epsilon/4$  and  $3r^2/32k_H^2$  contributions which gives a n(n-1)/6.4 added to the "n" term component inside the radical of equation 9.14.

In our new pde  $\delta J$ =o through LS spin-orbit coupling so the three spin<sup>1</sup>/<sub>2</sub>s and the L=1 add to a minimum. 1-<sup>1</sup>/<sub>2</sub>-<sup>1</sup>/<sub>2</sub>+<sup>1</sup>/<sub>2</sub> =<sup>1</sup>/<sub>2</sub> =S for the proton with possible Pauli principle non S=<sup>1</sup>/<sub>2</sub> possibilities for larger mass eigenvalue.

#### **Details of Above Solutions for Case 1**

Thus besides the ground state (N=0  $F_{groundstate} = \sum a_n r^n = a_0 r^{-0} = F_0$  proton) we have the two solutions:

 $F_{N=-1} = \sum a_n r^n = a_{-1}r^{-1} = F_1$ ,  $j = \frac{1}{2}$ , 0.,  $F_{N=-\frac{1}{2}} = \sum a_n r^n = a_{-\frac{1}{2}} r^{-\frac{1}{2}} = F_2$ . For  $j = \frac{1}{2}$ . 0. Note the energy eigenvalues (E) can be found from the solution to equation 9.14 and  $k_H = 1$  with E=1=938 MeV. Thus

N= 0, j=  $\frac{1}{2}$  then 9.14 gives +Nucleon (ground state) mass eigenvalue. Note that for the N=0, (with J=  $\frac{1}{2}$  and also J= 0 in section 9.5) ground state that the charge density is uniform (i.e.,  $\rho = K \propto r^0$ ) for r<k.

N=- $\frac{1}{2}$ , j= $\frac{1}{2}$  two valued because of the *two* square root solutions. Equation 9.14 then gives  $\Sigma \pm$  (charged sigma particle) 1184Mev particles, F<sub>2</sub> eigenfunction(s). Actual 1189Mev

N=-1, j=  $\frac{1}{2}$  gives *one* charged  $\Xi$  particle. Therefore the energy from equation 9.14 is 1327 Mev (actual 1321), F<sub>1</sub> eigenfunction.

Case 2 and case 3 give the neutral hyperons and  $\Lambda_0$  respectively (see case 3 below).

### 9.5 Nucleon Wavefunction: J=1, q≠0, N=-1 of Case 1

Here we recall case 1, section 9.3 above and compute energy eigenvalues for J=0 and J=1.again using equation 9.14 in case 1. J=0

N=-1, j=0 E=490 MeV from equation 9.14 case 1.  $K^{\pm}$ . Substitute into strangeness equation 9.34 case 1 we obtain strangeness =1.

N=0, j=0 then from equation 9.14 E=139.7 MeV

(9.22)

case (note again m=1+ $\epsilon$ =1.061 in 9.14 for outside). This is the nontrivial F zero point energy (and so has a fundamental SP hybrid state harmonic) for r<k= $\epsilon$  at r=r<sub>H</sub>. since the square root in equation 20.1 becomes imaginary then. Thus the mass of  $\pi^{\pm}$  is now the vacuum (e.g., note F $\propto$ r<sup>o</sup> for N=0 here)  $\epsilon$ 'at r≈k explaining why this fundamental harmonic result for  $\pi$  is used in all the successful nuclear force theories such as in the Skyrmion Lagrangian for example. Note that:

 $m_{\pi\pm}=139$ Mev=1.3(105.6MeV) =1.3 $\epsilon$ =.08= $\epsilon$ '

N=-1,J=1 case 1. Recall for J=1 we have  $\psi \propto rsin\theta \propto Y^{1}_{1}(\theta,\phi)$  double lobe  $\psi^{*}\psi$  along the z axis: From equation 9.14 we find with these inputs that E=1867Mev (9.23) implying that (because E~2m<sub>p</sub> and J=1) this eigenstate is responsible for the spin 1 deuteron (state). The L=1, 2P state solution(s) are symmetric and so of the form  $(1/\sqrt{2})(\psi_{1}\psi_{2} + \psi_{1}\psi_{2}) = \psi_{s}$  and have positive parity even if the 2P  $\psi_{1}$  and  $\psi_{2}$  each has negative parity. The Deuteron thus has + parity (Enge, 1966).

Recall if we include the background metric in eq.6.4.11  $\kappa_{oo}=1+r_H/r+2\epsilon^2+\Delta\epsilon$  and

 $\kappa_{rr}=1/(1+r_H/r+\Delta\epsilon)$ . So rescaling r $\rightarrow$ r- $\epsilon$ ' =r' for r near r<sub>H</sub> allows us to use our above solutions again. So in equation 8.1  $1/\sqrt{\kappa_{rr}}=\psi=1/\sqrt{(1+r_H/r'+\Delta\epsilon)}\psi\approx1/\sqrt{(1+r_H/r)}\psi+(\epsilon/2)\psi$ . Note if we again rescale our numerator J=1 $\rightarrow$ 1+( $\epsilon$ '/2)2 so that we have perturbed our Y<sub>1</sub> spherical harmonic with a ( $\epsilon/2$ )Y<sub>2</sub> giving a measure of the oblate, non spherical structure (e.g quadrupolar  $\psi_D$  and higher.  $\epsilon'/2 \approx .04$  from 9.22 therefore the nonspherical component of  $\psi$  is approximately 4% of the total  $\psi$  and is often called the tensor component of the Deuteron eigenstate (Enge, 1966). This simplest multiparticle state represents the *deuteron* state and this is then the explanation for the deuteron tensor component of the nuclear force.

Also the energy of the Deuteron is given just outside the  $r_H$  boundary (so  $\epsilon' \rightarrow i\epsilon$  in 6.4.11) by  $E_D=Rel\ 1876/\sqrt{\kappa_{oo}}=Rel\ 1876/\sqrt{(1+i\epsilon')+..=1876(1-i\epsilon'/2+(3/8)(i\epsilon')^2+..)}$ . So the added real term due to the  $\epsilon'$  is equal to  $1876(3/8)\epsilon'^2=1876(3/8)(.08)^2=4MeV$ . In free space  $\epsilon'=0$  and just outside the nucleus it gives this contribution to the Deuteron energy. Thus this  $(3/8)\epsilon'^2$  is the binding energy of the Deuteron.

Note from the equation 9.15 discussion for N *not* -1 we can only use J=1/2 and J=3/2 thus are restricted for two particles to S and P states (i.e.  $\frac{1}{2} + \frac{1}{2} = 1$ ) which then gives us the hyperons. For N=-1 we can use other J and can thereby construct large nuclei.

The multinucleon nuclei really are the solutions of the indicial equations of 9.15.

Recall in the shell model a hard shell nuclear outer wall is assumed with free space oscillations allowed inside this shell. The solutions to the Schrodinger equation are then spherical Bessel functions with corrections for spin orbit interaction, finite well height and tapered wells (Herald Enge, Introduction To Nuclear Physics, P.145). In any case an infinite mean free path for these oscillations is assumed to exist inside this shell. So how can there be an infinite mean free path inside this extremely high mass density region?

In that regard the above 2, J=1, N=-1 2P deuteron state can also be viewed as yet another Bogoliubov pairing interaction (such as in the SC section 4.5) giving this infinite mean free path of the electron pairs comprising a pion acting as a Cooper pair, just as in SC In the context of the section 4.5 pairing interaction model  $A(dv/dt)/v^2$  is no longer as small but dv/dt becomes very large to due to the ultrarelativistic motion of the electrons inside the nucleons. In any case this infinite mean free path for these oscillations (recall Cooper pairs have an infinite mean free path) is thereby explained here as a new type of superconductivity.

#### **Spin Orbit Interaction In Shell Model**

Recall the derivation of the shell model from first principles in section 6.12. If equal numbers of Neutrons and Protons gyromagnetic ratios then  $gy_P-gy_N = 2.7-1.9 = .8$ .

Since more neutrons in heavier elements: (1/1.1)(.8)=.7.

 $R=r_H \equiv \frac{1}{2}$  Fermi measured from singularity at  $1-\frac{1}{2} = \frac{1}{2}$ .

From 2P<sub>3/2</sub> at r=r<sub>H</sub> Fitzgerald contraction discussion in section 2.2:  $r \rightarrow R = \frac{1}{2}(1-\frac{1}{2}) = \frac{1}{4}$  Fermi =  $R_V(r-r_H)$  so  $R_v(r-r_H) \rightarrow Kr$ . From Ch1, sect 4.16 V=1/(r-r\_H). Spin orbit interaction=  $a_o^2(1/r)(\partial V/\partial r)(s \bullet L) =$ 

$$a_0^2 \frac{1}{R_V(r-r_H)} \frac{\partial V}{\partial (R_V(r-r_H))} (s \bullet L) = \frac{.7}{R_V(r-r_H)} \left( \frac{-1}{(R_V(r-r_H))^2} \right) (s \bullet L) = .7(4^3) (s \bullet L) \frac{1}{r} \frac{\partial V}{\partial r} = .7(64) (s \bullet L) \frac{1}{r} \frac{\partial V}{\partial r} = a_0^2 \frac{1}{r} \frac{\partial V}{\partial r} (s \bullet L) = .45 \text{*E\&M spin orbit interaction.}$$

Thus the  $a_0$ =1Fermi. Thus the nuclear spin-orbit interaction is much larger than the E&M spin orbit interaction because the nucleons are much closer to  $r_H$  than to r=0 and the Fitzgerald contraction of the nucleon  $2P_{3/2}$  state is on the order of  $\frac{1}{2}$ .

At close range there are higher energies available so the 4mev (=be) in equation 9.3 (if we include  $r^2$  contributions) becomes the binding energy for the deuteron in  $g_{oo}=1-k/r+be$  in 8.1 particles,  $F_2$  eigenfunction(s). Actual 1189Mev

N=-1, j=  $\frac{1}{2}$  gives *one* charged  $\Xi$  particle. Therefore the energy from equation 9.14 is 1327 Mev (actual 1321), F<sub>1</sub> eigenfunction  $\equiv \Xi$ s the fundamental structure for m=1.5. So we reapply the analysis all over again for mp->1.5 insteard of 1.

Case 2 and case 3 give the neutral hyperons and  $\Lambda_o$  respectively (see main Frobenius series solution paper).

The multinucleon nuclei are the solutions of the indicial equations of 9.15.

Recall in the shell model a hard shell nuclear outer wall is assumed with free space oscillations allowed inside this shell. The solutions to the Schrodinger equation are then spherical Bessel functions with corrections for spin orbit interaction, finite well height and tapered wells (Herald Enge, Introduction To Nuclear Physics, P.145). In any case an infinite mean free path for these oscillations is assumed to exist inside this shell. So how can there be an infinite mean free path inside this extremely high mass density region?

In that regard the above 2, J=1, N=-1 2P deuteron state can also be viewed as yet another Bogoliubov pairing interaction (such as in the SC section 4.4) giving this infinite mean free path of the electron pairs comprising a pion acting as a Cooper pair, just as in SC In the context of the section 4.5 pairing interaction model  $A(dv/dt)/v^2$  is no longer as small but dv/dt becomes very large to due to the ultrarelativistic motion of the electrons inside the nucleons. In any case this infinite mean free path for these oscillations (recall Cooper pairs have an infinite mean free path) is thereby explained here as a new type of superconductivity.

#### **Particle Lifetimes**

Recall from section 1.1:  $\kappa_{oo}=1$ - $r_H/r$  so  $r_r\kappa_{oo}=r_H$  analogous to  $dr_rct\kappa_{oo}=ds$  so  $r_H=ds\equiv|dZ|$ . From section 6.7 there are three Dirac equation contributions with one being the ultrarelativistic  $m_v$  contribution. For that contribution we put Dirac  $\alpha s$  into dr+idt=dZ the free space Dirac equation. Dividing by ds gives mass on the right side in that Dirac equation. Because the motion of the  $m_v = 1eV$  (Ch.3) particle is ultrarelativistic in these hadrons we apply figure 1-1 dr=dt so  $\theta=45^\circ$  and so  $dZ/ds = e^{i\pi/4}dr/ds$  for the ultrarelativistic  $m_v$  (on earth contribution of Ch.3). Note that  $(e^{i\pi/4})^2=i$ . We add another contribution (for spin  $\frac{1}{2}$ , N=-1) to get zero charge case II below. For added  $2P_{\frac{1}{2}}$  (K, $\pm\pi$  mesons) there are 3e in  $r_H$  below (sect.10.3). Thus we obtain:

**hyperons**, **Kaons** and  $\pm \pi$ :  $e^{i\pi/4}2e^2/m_vc^2 = e^{i\pi/4}r'_H = R_H$ Recall that domain r=r<sub>H</sub> was the most stable, the proton state. This stability condition can be restated in terms of excess energy above the proton rest mass. Next substitute this m and ultrarelativistic m<sub>v</sub> in the r<sub>H</sub> in equation 9.14 with this r'<sub>H</sub> in the relativistic solution of equation 2 described in Ch.1,sect.1.

$$E = \sqrt{m_p^2 + \frac{1}{R_H^2} (j^2 + 1.7071j - 1.10355 - (j.5303 + .8269)N)}$$
  

$$\approx m_p \left(1 + \frac{(e^{i\pi/4})^2 (j^2 + 1.7071j - 1.10355 - (j.5303 + .8269)N)}{2m_p^2 r'_H^2}\right)$$

Add to above to 9.14 result to get for the total energy:

$$m_{p}\left(1 + \left(\left(\frac{e^{i\pi/4}}{r'_{H}}\right)^{2} + \left(\frac{1}{r_{H}}\right)^{2}\right) \frac{(j^{2} + 1.7071j - 1.10355 - (j.5303 + .8269)N)}{2m_{p}^{2}}\right)$$

Plug  $(hc/e^2)^2 = (1/\alpha)^2$  back in eq.8.1 and normalize  $m_v c^2$  to 1/hz with 1/h. Next plug into the time propagator  $e^{iHt}$  and get for the r'<sub>H</sub> (decay) term:

$$= \exp i \left[ \left( \left( m_p c^2 / h \right) + \left( e^{i\pi/4} \right)^2 \left( m_v c^2 / h \right) \frac{m_v}{m_p} \left( j^2 + 1.7071 j - 1.10355 - \left( j.5303 + .8269 \right) N \right) \right] \left( \frac{\hbar c}{2e^2} \right)^2 \right] t$$
  
$$= \exp i \left[ \left( \left( m_p c^2 / h \right) + i \left( m_v c^2 / h \right) \frac{m_v}{m_p} \frac{i \left( j^2 + 1.7071 j - 1.10355 - \left( j.5303 + .8269 \right) N \right)}{(2\alpha)^2} \right] \right] t \qquad (9.23)$$
  
$$= \exp i \left[ \left( (m_v c^2 / h) + i A \right) t \text{ giving hyperon. Kaon. } \pm \pi \text{ decay times.}$$

= exp  $i(((m_p c^2 / h) + i\Delta))t$  giving hyperon, Kaon,  $\pm \pi$  decay times. The second term  $\Delta$  is also the excess mass above the proton mass.

For neutrons (939Mev) the excess mass above the proton mass (938Mev) is  $m_p/1000$  and  $R_H \rightarrow 1000R_H$ ,  $\Delta \rightarrow \Delta'$ 

$$E^{2} = m_{p}^{2} + \frac{1}{1000^{2}} \frac{1}{(1000R_{H})^{2}} (j^{2} + 1.7071j - 1.10355 - (j.5303 + .8269)N)$$

gives the neutron decay time.

For  $m_{\mu}$  muons  $j=\frac{1}{2}$ , N=0 and the excess mass is  $m_p/8.87 \equiv m_{\mu}$ .

$$E^{2} = m_{p}^{2} + \frac{1}{8.87^{2}} \frac{1}{(8.87R_{H})^{2}} (j^{2} + 1.7071j - 1.10355 - (j.5303 + .8269)N)$$

gives time for muon  $m_{\mu}$  decay.

For  $\pi^{o}$  decay time  $m_{\nu} \rightarrow m_{e}$  (E&M decay) along with 8.87 $\rightarrow$ 7= $m_{p}/m_{\pi}o$  in the above equation.

For resonances  $m_v \rightarrow m_e$  (E&M decay) in 9.23 gives time of decay.

Note the second term here contains a ii=-1 and so it is a exponential decay term  $e^{-Et}$  with .693/E=t the "half life".

Thus we get  $\pi^{0}$ ,  $\pm \pi$ , K mesons and hyperon, muon, neutron, resonance half lives from (these modifications of) equation 9.23.

#### 9.7 CASE 2 Excited State F, charge=0. 2P1/2

Recall from 9.4 that case 1 implies  $E_q \rightarrow m$  in case 2 (in 9.4). Also

 $1/g_{rr} \approx 1-k_H/k_H+\epsilon = \epsilon$  for - $\epsilon$ . Net charge=Zero. Thus let R=k\_H+r, r<<R, r'=k\_H\epsilon+r

$$\begin{bmatrix} \left(\frac{dt}{ds}\sqrt{g_{oo}}m\right) + m \end{bmatrix} F - \hbar c \left(\sqrt{g_{rr}}\frac{d}{dr} + \frac{j+3/2}{r}\right) f = 0$$

$$\begin{bmatrix} \left(\frac{dt}{ds}\sqrt{g_{00}}m\right) - m \end{bmatrix} f + \hbar c \left(\sqrt{g_{rr}}\frac{d}{dr} - \frac{j-1/2}{r}\right) F = 0$$

$$\sqrt{g_{rr}} = 1/\sqrt{1 - k_H/R + \varepsilon} \approx \frac{\sqrt{R}}{\sqrt{R - k_H + R\varepsilon}} = \frac{\sqrt{k_H + r}}{\sqrt{k_H + r - k_H + (k_H + r)\varepsilon}} \approx \frac{\sqrt{k_H + r}}{\sqrt{k_H \varepsilon + r(1 + \varepsilon)}} \approx \frac{\sqrt{k_H \varepsilon}}{\sqrt{k_H \varepsilon + r}} = \frac{\sqrt{k_H + r}}{\sqrt{r}}$$

Also  $(dt/ds)\sqrt{g_{oo}} \rightarrow E$  in the Dirac equation 18.1. Therefore equation 19.1 reads: r'=k\_{\rm H}\epsilon+r

$$[E+m]F - \hbar c \left( \sqrt{\frac{k_H}{r'}} \frac{d}{dr} + \frac{j + \frac{3}{2}}{k_H + r} \right) f =$$
$$[E+m]F - \hbar c \left( \sqrt{k_H/r'} \frac{d}{dr} + \frac{j + \frac{3}{2}}{k_H + r} \right) f = 0$$

$$[E+m]F - \hbar c \left( \sqrt{k_H/r'} \frac{d}{dr'} + \left(1 - \frac{r}{k_H}\right) \frac{j + \frac{3}{2}}{k_H} \right) f = 0 \text{ and}$$

$$\begin{split} & [E-m]f + \hbar c \left( \sqrt{\frac{k_H}{r'}} \frac{d}{dr} - \left(1 - \frac{r}{k_H}\right) \frac{j - \frac{1}{2}}{k_H} \right) F = 0 \qquad \text{Thus} \\ & \cdot \\ & f = -\frac{\hbar c}{|E-m|} \left( \sqrt{k_H/r'} \frac{d}{dr'} - \left(1 - \frac{r}{k_H}\right) \frac{j - \frac{1}{2}}{k_H} \right) F \qquad \text{Therefore} \\ & [E+m]F - \hbar c \left( \sqrt{k_H/r'} \frac{d}{dr'} + \left(1 - \frac{r}{k_H}\right) \frac{j + \frac{3}{2}}{k_H} \right) f = 0 \quad \text{Using} \qquad \mathbf{r} = \mathbf{r'} + \mathbf{k}_H \varepsilon \\ & [E+m]F - \hbar c \left( \sqrt{k_H/r'} \frac{d}{dr'} + \left(1 - \left((r' - \varepsilon k_H)/k_H\right)\right) \frac{j + \frac{3}{2}}{k_H} \right) \frac{-\hbar c}{|E-m|} \left( \sqrt{k_H/r'} \frac{d}{dr'} - \left(1 - \left((r' - \varepsilon k_H)/k_H\right)\right) \frac{j - \frac{1}{2}}{k_H} \right) F = 0 \end{split}$$

Multiplying both sides by |E-m<sub>p</sub>| we obtain:

$$\frac{E^{2} - m^{2}}{(\hbar c)^{2}}F + \left(\sqrt{k_{H}/r'}\frac{d}{dr'} + \left(1 - \frac{(r' - k_{H}\varepsilon)}{k_{H}}\right)\frac{j + \frac{3}{2}}{k_{H}}\right)\left(\sqrt{k_{H}/r'}\frac{d}{dr'} - \left(1 - \frac{r' - k_{H}\varepsilon}{k_{H}}\right)\frac{j - \frac{1}{2}}{k_{H}}\right)F = 0$$

$$\frac{E^{2} - m^{2}}{(\hbar c)^{2}} - \left(\left(1 + 2\varepsilon\right)\frac{(J + 3/2)(J - 1/2)}{k_{H}^{2}}\right)F + \sqrt{k_{H}/r'}\frac{J - 1/2}{k_{H}^{2}} + \left(1 + \varepsilon\right)\frac{(J + 3/2)}{k_{H}}\sqrt{\frac{k_{H}}{r'}}\frac{d}{dr'} + \left(\frac{k_{H}}{r'}\frac{d^{2}}{dr^{2}} - \frac{1}{2}\frac{k_{H}}{r^{4/2}}\frac{d}{dr'}\right)F = 0.$$
 Multiplying both sides by r'<sup>2</sup> we obtain:
$$\left(\left[\frac{E^{2} - m^{2}}{(\hbar c)^{2}}r'^{2}\right] - \left(1 + 2\varepsilon\right)(J + 3/2)(J - 1/2)\frac{r'^{2}}{k_{H}^{2}}\right)F + \left(1 + \varepsilon\right)\frac{(J + 3/2)}{\sqrt{k_{H}}}r^{3/2}\frac{d}{dr'} + \left[\left(r'k_{H}\frac{d^{2}}{dr'^{2}}\right) - \frac{1}{2}k_{H}\frac{d}{dr'} - \frac{\sqrt{k_{H}}r'^{3/2}}{k_{H}^{2}}(J - 1/2)\right]F =$$

Defining r'=r<sup>2</sup> and doing the derivatives in the new variable:  $\frac{dF}{dr^2} = \frac{dF}{dr}\frac{dr}{dr^2} = \frac{1}{2r}\frac{dF}{dr} \quad \text{and}$   $\frac{d^2F}{dr^2} = \frac{1}{2r}\frac{d}{dr}\left(\frac{1}{2r}\frac{dF}{dr}\right) = \frac{1}{2r}\left(-\frac{1}{2r^2}\frac{dF}{dr}\right) + \frac{1}{2r}\frac{1}{2r}\frac{d^2F}{dr^2} =$   $-\frac{1}{4r^3}\frac{dF}{dr} + \frac{1}{4r^2}\frac{d^2F}{dr^2}$ Substituting these expressions for the derivatives in:

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$$\begin{split} & \left[\frac{E^{2}-m^{2}}{(\hbar c)^{2}}r^{4}-(1+\varepsilon)(J+3/2)(J-1/2)\frac{r^{4}}{k_{H}^{2}}\right]F^{+} \\ & \left[\frac{k_{H}}{4}\frac{d^{2}}{dr'^{2}}-\frac{1}{4r}k_{H}\frac{d}{dr'}\right]F^{-}\frac{\sqrt{k_{H}}r'^{3}\left(J-1/2\right)}{k_{H}^{2}}F^{+}\frac{r^{3}(1+\varepsilon)(J+3/2)}{k_{H}}\frac{\sqrt{k_{H}}}{2r}\frac{d}{dr'}F^{+} \\ & -\frac{k_{H}}{4r}\frac{d}{dr}F^{=} \\ & \left[\left(\frac{E^{2}-m^{2}}{(\hbar c)^{2}}\right)-(1+2\varepsilon)\frac{(J+3/2)(J-1/2)}{k_{H}^{2}}\right]\sum a_{n}r^{n+4}+\frac{k_{H}}{4}\sum(n-1)na_{n}r^{n-2} \\ & \frac{k_{H}}{4}\sum na_{n}r^{n-2}-\frac{\sqrt{k_{H}}}{k_{H}^{2}}\left(J-\frac{1}{2}\right)\sum a_{n}r^{n+3}+(1+\varepsilon)\frac{(J+3/2)}{2\sqrt{k_{H}}}\sum na_{n}r^{n+1}-\frac{k_{H}}{4}\sum na_{n}r^{n-2} \\ & \left(\frac{E^{2}-m^{2}}{(\hbar c)^{2}}-(1+2\varepsilon)\frac{(J+3/2)(J-1/2)}{k_{H}^{2}}\right)\sum a_{n-4}r^{n}+\frac{k_{H}}{4}\sum(n+1)(n+2)a_{n+2}r^{n} \\ & \left(\frac{k_{H}}{4}\sum(n+2)a_{n+2}r^{n}-\frac{\sqrt{k_{H}}}{k_{H}^{2}}\left(J-\frac{1}{2}\right)\sum a_{n-3}r^{n} +(1+\varepsilon)\frac{(J+3/2)}{2\sqrt{k_{H}}}\sum(n+1)a_{n-2}r^{n} \\ & \frac{k_{H}}{4}\sum(n+2)a_{n+2}r^{n}=0. \end{split}$$

Combining terms noting simplification due to combining the  $a_{n+2}$  terms

$$\left(\frac{E^2 - m^2}{(\hbar c)^2} - (1 + 2\varepsilon)\frac{(J + 3/2)(J - 1/2)}{k_H^2}\right) \sum a_{n-4}r^n + \frac{k_H}{4} \sum (n-1)(n+1)a_{n+2}r^n + \frac{\sqrt{k_n}}{k_H^3}(J - 1/2) \sum a_{n-3}r^n + (1 + \varepsilon)\frac{J + 3/2}{2\sqrt{k_H}} \sum (n-1)a_{n-1}r^n = 0$$

Next we write the individual eigenfunctions as:

$$\left(\frac{E^2 - m^2}{(\hbar c)^2} - (1 + 2\varepsilon)\frac{(J + 3/2)(J - 1/2)}{k_H^2}\right) \sum a_{n-4}r^n = 0.$$

Thus since these series terms add to zero:

$$E = \sqrt{m^{2} + (\hbar c)^{2} (1 + 2\varepsilon) \frac{(J + 3/2)(J - 1/2)}{k_{H}^{2}}}$$
(9.24)  

$$(1 + \varepsilon) \frac{J + 3/2}{2\sqrt{k_{H}}} \sum (n-1)a_{n-2}r^{n} = 0. \text{ Here } r^{2} = r^{2} \text{ so } r^{-2/2} = r^{-1} = \sqrt{r}^{-2}$$
  

$$\frac{k_{H}}{4} \sum (n^{2} - 1)a_{n+2}r^{n} = 0$$
(9.25)  

$$-\frac{\sqrt{k_{H}}}{k_{H}^{2}} (J - 1/2) \sum a_{n-3}r^{n} = 0$$
(9.26)

J=1/2 with N=1 solves the indicial equation implied by 9.24-9.26. Recall from 9.4 that m=proton in this case (case 2). The energy in 9.24 is then that of a neutral particle (q=0) with the mass of

the neutron so  $E = E_q = m = m_N$ . See equation 9.23b for neutron lifetime and 2P3/2 for neutron spherical harmonic state, section 10.3) But in case 2 and equation 9.23 then the previously derived charged spin  $\frac{1}{2}$  hadrons  $m_{\Sigma}$ ,  $m_{\Xi}$  can also be put back into the Dirac equations for 'm'(instead of the proton). Thus the charged,  $m_{\Sigma}$ ,  $m_{\Xi}$  from equation 9.14 can be put into the "m" in 9.24 which gives the **neutral**  $E=m=m_N$ ,  $m_{\Xi}$ .  $m_{\Sigma}$  has a N=1/2 and so does not satisfy the above equations and so does not exhibit a stable neutral  $\Sigma$ . Recall the  $\Omega^-$  (which is J=3/2) is not J=  $\frac{1}{2}$  so doesn't have a neutral counterpart as does the proton and these other J=  $\frac{1}{2}$  hyperons. Recall the iterated Dirac equation is the Klein Gordon (in  $\chi$  with J=0) equation eigenstate transitions.

#### J=0, q=0 Case 2

Recall J=0 is allowed in every case.

m=1 proton, j=0 in equation 9.24 means K Long. Equation 9.23 gives K long mass eigenvalue: 1+(0+3/2)(0-1/2)/1=1/4. Thus  $\sqrt{.25} = .5$ . Thus .5X938X1.06=497 MeV=K<sub>long</sub>. Note case 2 is zero charge and note also from section 9.8 that the Strangeness=2 $|\sqrt{.5}|=2*.707 \approx 1$  as in strangeness equation 9.34 below.

m≈1 for Neutron then in 9.24 we have K short, if m=m<sub>±</sub> and J=0 then D<sup>o</sup> Long. If m=m<sub>±</sub> j=0, and neutral then 9.24 gives D<sup>o</sup> Short.

#### 9.8 CASE 3 m=0, so $\psi_L$ , f state, charge=0 (lower case of equation 9.5).

In case 3 there is no central force therefore N=0 and  $j=\frac{1}{2}$  in f. This is the m=0 left handed doublet case of Chapter 3. Let R=k<sub>H</sub>-r, r<<R for stability we can write:

$$\begin{split} \sqrt{g_{rr}} &= 1/\sqrt{1+k_H/R+\varepsilon} \approx \frac{\sqrt{R}}{\sqrt{R+k_H+R\varepsilon}} = \\ &\frac{\sqrt{k_H-r}}{\sqrt{k_H-r}} = \frac{\sqrt{k_H-r}}{\sqrt{k_H(2+\varepsilon)-r(1+\varepsilon)}} = \frac{(1-\varepsilon/2)}{\sqrt{2}} \left( \frac{1-\frac{r}{2k_H} + \frac{r^2}{8k_H^2} + ..}{1-\frac{r}{4k_H} + \frac{r^2}{16k_H^2} + ..} \right) = \\ &((1-\varepsilon/2)/\sqrt{2}) \left( 1-\frac{r}{4k_H} + \frac{3r^2}{32k_H^2} - ..) \approx \frac{1-\frac{r}{4k_H}}{\sqrt{2}} \end{split}$$

Therefore equation 9.1 reads:

$$\begin{bmatrix} E + m_p \end{bmatrix} F - \hbar c \left( (1 - r/k_H 4) \frac{d}{\sqrt{2}dr} + \frac{j + \frac{3}{2}}{k_H - r} \right) f = 0$$
  
$$\begin{bmatrix} E + m_p \end{bmatrix} F - \hbar c \left( (1 - r/k_H 4) \frac{d}{\sqrt{2}dr} + (1 + r/k_H)(j + \frac{3}{2})/k_H \right) f = 0;$$
  
$$\begin{bmatrix} E - m_p \end{bmatrix} f + \hbar c \left( \sqrt{g_{rr}} \frac{d}{dr} - \frac{j - \frac{1}{2}}{r} \right) F = 0$$
(9.27)

From the above equation 9.27 if (and  $j=\frac{1}{2}$ ) m<sub>p</sub>=0 then

$$\begin{bmatrix} E - \boldsymbol{m}_p \end{bmatrix} f + \hbar c \left( \left( 1 - r/k_H 4 \right) \frac{d}{\sqrt{2}dr} - \left( 1 + r/k_H \right) \left( j - \frac{1}{2} \right) / k_H \right) F = 0$$
  
Therefore (with j= ½) from equation 9.27 for small r. In any case:  
$$F = \hbar c \frac{\hbar c}{\left[ E + \boldsymbol{m}_p \right]} \left( \left( 1 - r/k_H 4 \right) \frac{d}{\sqrt{2}dr} + \left( 1 + r/k_H \right) \left( j + \frac{3}{2} \right) / k_H \right) f$$
$$\begin{bmatrix} E + \boldsymbol{m}_p \end{bmatrix} F - \hbar c \left( \left( 1 - r/k_H 4 \right) \frac{d}{\sqrt{2}dr} + \frac{j + \frac{3}{2}}{k_H - r} \right) f = 0$$

Solving for f and substituting back in 9.27

$$\begin{split} & \left[E - m_{p}\right]f + \hbar c \left(\left(1 - r/k_{H} 4\right) \frac{d}{\sqrt{2}dr} - \left(1 + r/k_{H}\right)(j - .5)/k_{H}\right) \bullet \\ & \frac{\hbar c}{E + m_{p}} \left(\left(1 - r/k_{H} 4\right) \frac{d}{\sqrt{2}dr} + \left(1 + r/k_{H}\right)(j + \frac{3}{2})/k_{H}\right)F^{2} \left[E - m_{p}\right]f^{+} \\ & \left(\frac{(\hbar c)^{2}}{(E + m_{p})\sqrt{2}} \left(-\left(1 - r/4k_{H}\right) \frac{d}{k_{H} 4\sqrt{2}dr} + \left(1 - r/4k_{H}\right)^{2} \frac{d^{2}}{\sqrt{2}dr^{2}} + \left(1 - r/4k_{H}\right)(j + \frac{3}{2})/k_{H}^{2}\right)f^{+} \\ & + \frac{(\hbar c)^{2}}{E + m_{p}} \left(-\left(1 + 3r/k_{H} 4\right)(j - .5) \frac{d}{\sqrt{2}k_{H} dr} - \left(1 + r/k_{H}\right)^{2}(j + 1.5)(j - \frac{1}{2})/k_{H}^{2}\right)f^{2} \\ & \left(\left[E - m_{p}\right] + \left[\frac{(\hbar c)^{2}}{(E + m_{p})} \left(-\left(j + \frac{3}{2}\right)(j - \frac{1}{2})/k_{H}^{2} + (j + \frac{3}{2})/\sqrt{2}k_{H}^{2}\right)\right]\right)f^{+} \\ & \left(\frac{(\hbar c)^{2}}{(E + m_{p})\sqrt{2}} \left(2\sqrt{2}\left(j + \frac{3}{2}\right)(j - \frac{1}{2})/k_{H}^{3} - (j + \frac{3}{2})/4k_{H}^{3}\right)rf^{+} \\ & \left(\frac{(\hbar c)^{2}}{(E + m_{p})\sqrt{2}} \left(-\frac{1}{k_{H} 4\sqrt{2}} + \left(j - \frac{1}{2}\right)\frac{1}{k_{H}}\right)\frac{df}{dr} + \\ & \left(\frac{(\hbar c)^{2}}{(E + m_{p})\sqrt{2}} \left(\frac{1}{\sqrt{2}}\right)\frac{d^{2}f}{dr^{2}} + \\ & \left(\frac{(\hbar c)^{2}}{(E + m_{p})\sqrt{2}} \left(\frac{1}{\sqrt{2}}\right)\frac{d^{2}f}{dr^{2}} + \\ & \left(\frac{(\hbar c)^{2}}{(E + m_{p})\sqrt{2}} \left(\frac{-1}{2\sqrt{2}k_{H}}\right)r\frac{d^{2}f}{dr^{2}} \\ & \text{Next substitute in } F = \sum_{n}a_{n}r^{n} \\ & \frac{N}{M} \left[\left[E - m_{p}\right] + \left[\frac{(\hbar c)^{2}}{(E + m_{p})} \left(-\left(j + \frac{3}{2}\right)(j - \frac{1}{2}\right)/k_{H}^{2} + (j + \frac{3}{2})/\sqrt{2}k_{H}^{2}\right)\right]\right)a_{n}r^{n} + \\ \end{array}$$

$$\sum_{M}^{N} \frac{(\hbar c)^{2}}{(E+m_{p})\sqrt{2}} \left( 2\sqrt{2} \left( j + \frac{3}{2} \right) (j - \frac{1}{2}) / k_{H}^{3} - (j + \frac{3}{2}) / 4k_{H}^{3} \right) a_{n-1} r^{n} + \qquad (9.28)$$

$$\sum_{M}^{N} \frac{(\hbar c)^{2}}{(E+m_{p})\sqrt{2}} \left( -\frac{1}{k_{H}} \frac{4\sqrt{2}}{\sqrt{2}} + \left( j - \frac{1}{2} \right) \frac{1}{k_{H}} \right) (n+1) a_{n+1} r^{n} + \qquad (9.29)$$

$$\sum_{M}^{N} \frac{(\hbar c)^{2}}{(E+m_{p})\sqrt{2}} \left(\frac{1}{k_{H}^{2}16\sqrt{2}} - \left(j - \frac{1}{2}\right)\frac{3}{4k_{H}^{2}}\right) na_{n}r^{n} +$$
(9.30)

$$\sum_{M}^{N} \frac{(\hbar c)^{2}}{(E+m_{p})\sqrt{2}} \left(\frac{1}{\sqrt{2}}\right) (n+2)(n+1)a_{n+2}r^{n} +$$
(9.31)

$$\sum_{M}^{N} \frac{(\hbar c)^{2}}{(E + m_{p})\sqrt{2}} \left(\frac{-1}{2\sqrt{2}k_{H}}\right) (n+1)na_{n+1}r^{n} = 0$$
(9.32)

We now take 9.27 and 9.30 together and 9.29 and 9.32 together (since they have the same a<sub>n</sub>). Thus there are 4 independent series (with 9.28 and 10.31) here. The equation 9.27 and 9.30 nth terms give:

$$\left[ \left[ E + m_p \right] + \left[ \frac{(\hbar c)^2}{(E - m_p)} \left( - \left( j + \frac{3}{2} \right) (j - \frac{1}{2}) / k_H^2 + (j + \frac{3}{2}) / \sqrt{2} k_H^2 \right) \right] \right] + \frac{(\hbar c)^2}{(E - m_p)\sqrt{2}} \left( \frac{1}{k_H^2 16\sqrt{2}} - \left( j - \frac{1}{2} \right) \frac{3}{4k_H^2} \right) n,$$

At some value of n=N we have for a solution

$$\left(E^{2}-m_{p}^{2}\right)+\left(\left[\left(-\left(j+\frac{3}{2}\right)(j-\frac{1}{2})/k_{H}^{2}+(j+\frac{3}{2})/\sqrt{2}k_{H}^{2}\right)\right]\right)+\frac{1}{\sqrt{2}}\left(\frac{1}{k_{H}^{2}16\sqrt{2}}-\left(j-\frac{1}{2}\right)\frac{3}{4k_{H}^{2}}\right)N=0$$

therefore rearranging:

$$E = \sqrt{m_p^2 + \frac{1}{k_H^2} \left( (j+3/2)(j-1/2) + .7071j + 1.0607 + (.0156 - (j-.5).5303)N \right)}$$
(9.33)

Recall from the equation 9.4 'f' case that we have m<sub>p</sub>=m=0, and zero charge therefore no central force thus N=0 in f $\propto$ r<sup>o</sup> in equation 8.1. Therefore since there is small r and dr<sup>0</sup>/dt=d1/dt=0 in the equations just above equation 9.27 along with 9.33 then the 9.27-9.32 equations add to zero and thus are solved. Also the j=3/2 (so L=1) case is not allowed since that requires a central force to give  $L \neq 0$ ,  $j = \frac{1}{2}$  and of course j=0 is allowed here. Thus

N= 0, j=  $\frac{1}{2}$ , m=0 then from 9.33 we have E=1115.8 Mev  $\Lambda_0$ 

N=0, j=0,  $\eta$  mass and also gives m=.56 (with m=0) in 9.33 used in gyromagnetic ratio calculation for f. Recall  $\epsilon$ =.08 (with m=0) for F in 9.14. This is the nontrivial f state zero point energy for r<k since  $\Psi = \psi + \chi$  from our observability definition. Note Kaons then give no strange bound states because this mass is real (in contrast to the imaginary pion mass in 9.22).

#### 9.9 Strangeness

Recall that in 9.14 (which applies to Case 1 and Case 2) the energy is  $E^2=m_o^2+(j^2+1.7071j-1.10355-(j (.53033))+.7642))N/k_H^2$ . Now  $m_o^2$  and E is conserved ( $m_o$  is a constant) here and thus it appears that energy conservation implies that the square root of  $j^2+1.7071j-1.10355$ )-(j.5303+.7642)N =S must be conserved. Therefore  $E^2=m^2+S^2$  then and "S" is conserved for the charged core states and thus for the neutrals given that in section 9.8 that  $E_q \rightarrow m$  then (for f state m=0 we also have S≈E for  $\Lambda$ ). We could also write  $E^2=m^2+C^2$  for the next 2P state eigenstates (call C charm if you want) which would also have their own associated production (since <|> not zero). Thus, as an example, normalizing to a factor of 2X:

 $\begin{aligned} &2XSQR[(.5(.5303)+.7642)(0)]=0=S_{nucleon}, &2XSQR|[(.5(.5303)+.7642)(-1)]\approx 2=S_{\Xi}, \\ &2XSQR|[.5(.5303+.7642)(-1/2)]|\approx 1=S_{\Sigma}, &2XSQR|[(1.5^2+1.7071(1.5)-1.10355-(1.5(.5303+.7642)))(-1)]|\approx 3=S_{\Omega}. \end{aligned} \tag{9.34}$ 

Strangeness is only an approximate conservation law in the examples in 9.34 but there is enough conservation at least for the "associated production" and we have not yet included the weak interaction here. This is a **direct derivation of strangeness**, instead of just having postulated it as it is in the standard model and QCD. Strangeness isn't strange anymore.

Charm, bottom, top: In chapter 9 equation assuming hard spherical shell. We obtain other (less stable, resonances) particle groups using equation 9.5 by taking the quadratic approximation of  $g_{rr}$  (i.e., include the (3/32)(r/k<sub>H</sub>)<sup>2</sup> term in 9.5) Using 10.8 instead of just the linear approximation we used above. Recall that the perturbative (3/32)(r/k<sub>H</sub>)<sup>2</sup> term had to be included since it gave a  $\approx$ 20Mev correction to the hyperon masses.

## C Meson Mass Derivation From Potential Of Chapter 10 And The New Pde eq.9

## C Spherically Symmetric Wave Function Required

PROGRAMFracsN

DOUBLE PRECISION A,B,C,D,E,F,H,I,I1,J,KK DOUBLE PRECISION K1,K2,K3,K4,N1,N2,N3,N4,R,W,X,Y,Z DOUBLE PRECISION Y1,E1,E2,MM1,MM2,MM3,EE,JJ integer N,M,M1 DIMENSION EE(400)

- C Variational principle on E with respect to I and Y1,
- C RungeKutte on D equation 8.1. Y=2 width Deuteron
- C pion oscillation resonance modeled between 0 and Y=2. H=0.001

mH=2 !harmonic number for oscillation inside Y=2.

- C mN=1 gives pion 0 and K+-,mN=2 gives pi+- and Ko resonance ep=0.08\*mH !pion 1st and 2nd harmonic resonance added to Y1 W=1.0+ep !pion mass added to nucleon. J=0.0 !spin 0 mesons X=0.0001 !mass energy increments I1=100000000.0
  - A=0.0
  - B=0.0
  - C=0.0
  - E=0.0

KK=78.8 !gives MeV energy units JJ=J\*1. Y1=2.0+ep !pion increases Y1. 50 D=.0000001 I1=0.0F=.0000001 Y=Y160 R=Y V=1.0/(1.0+ep-R) !chapter 14 potential for spin 0 E1=EK1=((W-E-V)\*F)+(((J-0.5)/R)\*D)N1 = ((E+W+V)\*D) - (((J+1.5)/R)\*F)R = R + (0.5 \* H)V=1.0/(1.0+ep-R)K2=((W-E-V)\*(F+(0.5\*H\*N1)))+(((J-0.5)/R)\*(D+(0.5\*H\*K1)))N2=((E+W+V)\*(D+(0.5\*H\*K1)))-(((J+1.5)/R)\*(F+(0.5\*H\*N1)))K3=((W-E-V)\*(F+(0.5\*H\*N2)))+(((J-0.5)/R)\*(F+(0.5\*H\*K2)))N3=((E+W+V)\*(D+(0.5\*H\*K2)))-(((J+1.5)/R)\*(F+(0.5\*H\*N2)))R = R + (.5\*H)V=1.0/(1.0+ep-R)K4=((W-E-V)\*(F+(H\*N3)))+(((J-0.5)/R)\*(D+(H\*K3))) N4=((E+W+V)\*(D+(H\*K3)))-(((J+1.5)/R)\*(F+(H\*N3)))E=E1F=F+((H/6.0)\*(N1+(2.0\*N2)+(2.0\*N3)+N4))D=D+((H/6.0)\*(K1+(2.0\*K2)+(2.0\*K3)+K4)) $I = (F^*F) + (D^*D)$ 100 I1=I1+(I\*(R+(0.5\*H))\*(R+(0.5\*H)))) IF((abs(R-1.0-ep)).LT.(0.9\*H))THEN Y = Y - (2.0 \* H)GOTO 60 **ENDIF** Y=Y-H IF(Y.LT.0.0)THEN **GOTO 200 ENDIF** GOTO 60 200 E = E + XC=I1IF(B.LT.A)THEN **GOTO 310 ENDIF GOTO 312** 310 IF(C.GT.B)THEN **ENDIF** 312 IF(B.GT.A)THEN

GOTO 315 ENDIF GOTO 320 315 IF(C.LT.B)THEN print \* 'F=' (F-X)\*KK ' I=' I ' '

- print \*,'E=',(E-X)\*KK,' J=',J,' max I' ENDIF
- 320 IF(E.GT.8.0)THEN GOTO 349 ENDIF A=B B=C
- 330 GOTO 50
- 349 print\*,'program finished'
- 350 stop End

C **Results for** spin 0,L=0 are

C For mN=1 get 135MeV  $\pi^o$  and 493K^{\pm} for resonance with 1 meson.

C For mN=2 get 139Mev  $\pi^{\pm}$  and 497Mev K<sup>o</sup> for resonance with two mesons in ordinary nuclear matter nucleus would split before K energy created. In a neutron star however K s could be created.

This fortran computer program only requires a few seconds to run on a PC. On the other hand lattice gauge theory programs (assuming a SU(3) lattice) require massive computing power and really do not duplicate high energy liquid state strong interactions anyway.

Here the pion is a  $r=2R_H$  proton with no net rotation and the central electron in a m=0 state so net spin =0 .  $\sigma$ >>1/20Barn so annhilation occurs outside  $r_H$  and the pion decays.

## **Chapter 10**

# r≈r<sub>H</sub> Application: 2P3/2 Half Integer Spherical Harmonics Solutions. This is a continuation of Chapter 9

# 10.2 Overview of $2P_{3/2}$ Solutions to Equation 9 (the New Dirac equation) at $r \approx r_H$ in the Context of the Equivalence Principle (single charge e) Implication

Allowing this single charge 'e' to move near and inside that stable singularity radius  $r \approx 2q^2/mc^2$ in the  $\sqrt{g_{ii}}$  in this new Dirac equation (equation 2) as we see below makes the motion relativistic but stable requiring all the Dirac equation spherical harmonic solutions, not just the ones allowed by the Schrodinger equation. Also the next order of approximation above the hard shell for our  $g_{00}$  horizon  $r_H = 2e^2/m_ec^2$  is the harmonic oscillator  $V \alpha r^{+1}$  giving the SU(3) SYMMETRY of the three dimensional harmonic oscillator. The +1 in the exponent of V (instead of the inverse law-1)also reverses the sign the exchange square on integral  $\pm \int \psi^*_{111}(r') \psi^*_{lmn}(r'') V(r',r'') \psi_{lmn}(r') \psi_{111}(r'') d\tau = J$  designating the symmetric and antisymmetric states), making here then the J=3/2 state m=-3/2 and 3/2

(i.e.,  $\psi = \mathcal{Y}^{3/2}_{3/2}(\theta, \phi) + \mathcal{Y}^{-3/2}_{3/2}(\theta, \phi) = 2P_{3/2}$  eigenspinor) the first ground state that varies with azimuthal angle(baryons) above the already filled 1S (in analogy with helium) on the energy ladder instead of the expected  $\frac{1}{2}$  and  $-\frac{1}{2}$  (these  $\frac{1}{2}$  s by the way give  $2P_{1/2}$  in the  $\psi^*\psi$  of the next higher P orbital slots) that vary with azimuthal angle (baryons).

Also recall the identity  $(exp(i\phi) + exp(-i\phi))/2 = cos\phi$ . The  $\mathcal{Y}^{3/2}_{3/2}$  orbital is a  $exp(i3/2\phi)$  and  $Y^{-3/2}_{3/2}$  orbital is  $exp(-i3/2\phi)$  and thus from the identity the summed state is  $cos(3/2\phi)$  with probability density  $\psi^*\psi = cos^2(3/2\phi)$ , the trifolium three lobed shape. Thus there are TWO +e s giving a net charge of +2/3e in each lobe because the +electron charge 'e' is in each orbital lobe on the average only 1/3 the time (FRACTIONAL CHARGE) giving the many scattering properties (such as jets) associated with the angular distribution of multiple fractional charges interior to this horizon. The lobe 'structure' can't leave (ASSYMPTOTIC FREEDOM) as in the Schrödinger equation case or move so is NONRELATIVISTIC in contrast to its rapidly moving m<sub>e</sub> constituent.

Finally we solve the problem with the new pde using a computer program, set the boundary conditions as if the Deuteron was a square well. See end of chapter 9 for the fortran program. In any case we can build the hyperons and mesons with integer charges e, don't need the fractional charges.

## **10.3 Trifolium Diagram**



 $\begin{array}{ll} 2\mathsf{P}_{3/2} \text{ solutions to Dirac equation.} & \sqrt{\kappa_{\mu}}\mu^{\gamma\mu}\frac{\partial\Psi}{\partial x^{\mu}}-\omega\Psi=0 & \kappa_{00}=1-\frac{r_{H}}{r} & \text{Stability at}\\ \text{Ultrarelativistic LS coupling fills } 2\mathsf{P}_{3/2} & \text{first} & \kappa_{00}=\frac{1}{\kappa_{rr}} & \text{Stability at}\\ \text{Exchange integral implies (after filled 1S +2e) next bound state}= \end{array}$ 

$$\Psi = \mathscr{Y}_{3/2}^{3/2} + \mathscr{Y}_{3/2}^{-3/2} = \frac{e^{i\frac{3}{2}\phi} + e^{-i\frac{3}{2}\phi}}{2} = \cos(\frac{3}{2}\phi)$$

$$\psi^*\psi = \cos^2(\frac{3}{2}\phi) = \bigcirc$$
 trifolium

Electron charge e spends 1/3 of its time in each lobe making each lobe (1/3)e charged on average. For two such electrons it is (2/3)e.

Start with 1S state filled singlet = $(1/2 - 1/2) = Y_0^0$ . Add  $2P_{3/2}$  and electron -e for bound state e.g., proton. Add to this proton  $2P_{3/2}$  state -e and +e in filled  $2P_{1/2}$  and a third -e with spin down to get neutron(spin1/2) Singlet =  $\pi^0$  with one of the electrons  $\Delta \varepsilon$  in the first electron excited state  $\varepsilon$ (muon). Add -e and +e in filled  $2P_{1/2}$  and a third +e with down spin get  $\pi^+$ (spin0)

Figure 10-1 Trifolium diagram

2P<sub>3/2</sub> fills first in Aubau principle for ultrarelativistic hard shell (Alfredo 1998).



Possible SHM interaction between these lobes gives excited states. LS coupling Lande' g-factor structure gives minimal LS energy for smallest L So net spin 1/2 states preferred.

Fig2 Note we get a similar shape to the trifolium with lattice QCD theory.





Fig.3 Second diagram is the E field of a ultrarelativistic charged particle ("Electronic Motion", McGraw Hill, Harman)

Thus the neutron charge configuration allows for the creation of both the W and the  $H_v$  and the proton charge configuration does not allow this.



ultrarelativistic electron motion gives correct mass. plate field from object A does not intersect particle C so bound state.

Fig.4

### 10.5 Two (e+e+e at r≈r<sub>H</sub>) Objects: The Deuteron

The only stable lepton state in this theory so far are 2 (eq.2 in 4D) and e+e+e at r=r<sub>H</sub> with the latter being internally unstable but regenerative through that particle (re)creation step above (fig6). Also if a second e+e+e object is nearby a  $2P_{1/2}$  state (sect.1.6, eq.2) the second proton picks up the electron making this state also regenerative and stable (Deuteron). According to sec.7.3 the neutron in deuteron decays but its electron is then transferred to the other proton and it then days creating a net oscillation. For the electron at r<sub>H</sub> between the protons in the deuteron the electric potential energy PEe is still  $2ke^2/r_H$  because r'<sub>H</sub> for the electron is  $\sim 10^{-18}$  m. But time dilation still makes r<sub>H</sub> (=2X10<sup>-15</sup>m) clock stop dt'<sup>2</sup>=(1-r\_H/r)dt<sup>2</sup>.So the z=0 stable orbit is still at r<sub>H</sub>

Matrix

### **Schrodinger Equation For Many Deuterons**

Equation 9 iterated for Bosons is the Klein Gordon equation. Here we derive the SHM Schrodinger equation version for the deuteron. From the energy component of polarized representation of equation 8.1, 8.2: and using iterated (as in bosonic, here deuterons)  $E^2=p^2c^2+m^2_oc^4=$ 

$$E^{2} = \left(\frac{1}{\sqrt{\kappa_{oo}}}\right)^{2} = \frac{1}{1 - \frac{r_{H}}{r}} = \frac{r}{r - r_{H}} = \frac{r}{r - r_{H}} - \frac{r - r_{H}}{r - r_{H}} + 1 = \frac{r_{H}}{r - r_{H}} + 1 = V + 1$$
(10.8.1)

Note the resemblance of  $E^2=p^2c^2+m^2{}_{o}c^4$  to the Schrodinger equation if the  $E^2=[r_H/(r-r_H)]+1$ ,  $p^2c^2 \rightarrow k^2 \ 2^{nd}$  derivative of equation 10.8.1. r'\_H=[ $^3\sqrt{((6Z/4\pi))}r_H$  in eq.10.8.1. Note for low energy large r m<sub>e</sub>c<sup>2</sup>(V)≈PE=ordinary E&M potential energy. For high energy small r in contrast then V+1≈0=E<sup>2</sup>.

From the new pde model  $r=r_H=2.817X10^{-15}$  (DavidMaker.com, partI, eq.1.2.7)

So given the Schrodinger equation we can use the well known SHM and rotation mathematics. The nucleus is somewhere you would expect the equipartition of energy to hold (recall it holds quite well in a hot gas).. So as usual, given those ultrarelativistic positrons in each proton, as is well known PEe provides the highest escape speed in  $KE = \frac{1}{2}mv^2$  which in turn provides the  $\frac{1}{2}kx^2$  potential energy for the spring SHM model since the SHM also provides the same high speeds which you cannot exceed. Recall this same max  $\frac{1}{2}kx^2$  gives us  $\frac{1}{2}\omega_0(\frac{1}{2})$  in the SHM quantum model for the oscillation (model energy).

 $PE=2e^{2}/r_{H}=2(9X10^{9})(1.602X10^{-19})^{2}/2.8X10^{-15}=1.7X10^{-13}J.$ 

1.7X10<sup>-13</sup>/1.602X10<sup>-19</sup>=1.06X10<sup>6</sup>eV=1.06Mev. Binding energy=1.1Mev

From the equipartition of energy that free space 3D PE equals oscillation energy SHM equals the rotational energy LS so in the deuteron you have these three sources of binding energy. The single neutron binding energy is =PE=NE

So from the Equipartition of energy:

PE=SHM=LS

PE+SHM+LS=BindingEnergyD

MeasuredBindingEnergy=BindingEnergyD-NE

Then subtract the neutron binding energy because it doesn't count in the binding energy

### **Equipartition of Energy**

The central electron in the deuteron is also doing 1D simple harmonic motion! The  $(1/2)kr_{H}^{2}$  SHM PEosc = $2e^{2}/r_{H}$  gives (the same escape velocity v in  $(1/2)mv^{2}$ ) = $\hbar\omega_{0}\frac{1}{2}=2e^{2}/r$  (=1.1Mev). This electrostatic PE= $2ke^{2}/r_{H}=\frac{1}{2}kr_{H}^{2}$  is  $\hbar\omega_{0}\frac{1}{2}$ , the zero point energy oscillation. So PErot= PE= $2e^{2}/r_{H}=2(9X10^{9})(1.602X10^{-19})^{2}/2.8X10^{-15}=1.7X10^{-13}J$ .

1.7X10<sup>-13</sup>/1.602X10<sup>-19</sup>=1.06X10<sup>6</sup>eV=1.06=1.1Mev=KEOsc

The rotation degree of freedom energy  $(J(J+1)/(8\pi^2 I) = KErot$  also contributes = energy from the equipartition of energy so  $KE_{1D} = KErot = KEosc = 1.1$ Mev. So

 $KE_{1D}+KEosc+KErot=1.1+1.1+1.1=3.1$ Mev from the *equipartition of energy* in Deuterium. Subtract off the neutron PEe and get [(1+1+1)-1]1.1=2.2Mev binding energy of the Deuteron. The nucleus is made of Deuterons. The odd neutron just adds an excess PEe=PE<sub>3D</sub>.

We can continue this use of the equipartition of energy for larger and large nuclei which makes it just M(1.1+1.1+1.1)=energy for each deuteron in the nucleus where M ps a integer.

Note we did not have the absurd assumptions of 1D and motion through a vacuum as in the shell model. The 1D arises from the electron motion between the two protons which also moves in a vacuum between the two protons. So we recreated the shell model without all its silly assumptions and know why it works!

**10.6** Recall for J=1 we have  $\psi \propto rsin\theta \propto Y^{1}_{1}(\theta,\phi)$  double lobe  $\psi^{*}\psi$  along the z axis: From equation 9.14 we find with these inputs that E=1867Mev (9.23)

The deuteron. The only stable lepton state in this theory so far are 2e and 3e at r=r<sub>H</sub> with the latter being internally unstable but regenerative through that particle (re)creation step above (fig6). Also if a second 3e object is nearby a  $2P_{1/2}$  state (sect.1.6, eq.1.2.7) the second proton picks up the electron making this state also regenerative and stable (Deuteron). The two protons are stuck at the line of sight horizon  $r_H \approx 2.88 \times 10^{-15}$  m. Because of the ultrarelativistic motion in the transverse direction there is that Fitzgerald contraction down to nearly a point (lepton,  $S_{1/2}$  states) allowing Maxwell's equations to again be used and also showing we can use +e+e-e=+e charge for the center of each proton. A Gaussian pillbox can then be defined obtaining the usual ke<sup>2</sup>/r potential energy for each central positive e. Define r=2.88/ $\sqrt{2}$ . So the average distance to the central Deuteron from all 3 electrons is  $(r\sqrt{2}+r\sqrt{2}+r)/3=(2\sqrt{2}r+r)/3=2.6=r_{\rm H}$ '. Tabulated D charge radius=2.128X10<sup>-15</sup> m $\approx$ r\_H. Note the middle electron is also at  $2r_{\rm H}/2=r_{\rm H}$  since it has side to side wobble. From the equipartition of energy PE<sub>e</sub>=SHM so potential energy is PE=SHM +PE<sub>e</sub> so =2PE<sub>e</sub> and PE is provided by both protons on the central electron potential energy  $e^2/r_{\rm H}+e^2/r_{\rm H}$  between the electron and 2 protons is  $PE_e=2e^2/r_{\rm H}$ . NB= neutron binding energy.



#### Rotational energy at r<sub>H</sub>

The equation used here is the iterated **new pde** eq.2 (2nd derivative with eq.10.8.1) with SHM approximation of bond vibration spring potential energy with  $PE=2ke^2/r_H=\frac{1}{2}kx_o^2=p^2/2m=\frac{1}{2}\omega_o/2$  with 3D spherical harmonics with rotational energy  $[\frac{h^2}{2}(2I)](L(L+1))=[\frac{h^2}{2}(2mr^2)](L(L+1))$ 

#### **Comparison with QCD SU(3)**

The 3D components of the SHM tensor  $A_{ij}=(1/(2m))(p_ip_j+m^2\omega^2x_ix_j)$  and components of L satisfy Poisson bracket relations SU(3). So by including as a perturbation the rotation, the 3D SHM version gives SU(3) symmetry in the S derived from those  $A_{ij}$  (Herbert Goldstein, 'Classical Mechanics' 2<sup>nd</sup> edition, pp.425) which holds in both the classical and QM case then (So we have just explained the origin of the adhoc QCD gauged alternative to this nuclear physics.)

#### **Multiple Deuterons**

So there are then simple rules to follow for finding nuclear binding energy. They come out of the new pde and use the conservation of energy. Use equipartition of energy.

1) Number of bonds gives number of PE+LS (1PE=1.11Mev). LS=1 for PE=1.

2) LS=PE except if extra electron

3)  $Mh\omega(N+D/2)$  equally distributed among all bonds. N is SHM quantum number, D dimension. Equipartition of energy between Eosc, PE and LS.except (when extra electron).

4) subtract neutron number since binding energy does not include them.

1.11XN=	Mev
---------	-----

**10.7 Details of Finding Binding Energies: Deuterium**. P-e-P Note PE plus LC. Subtract 1neutron+1p  $E=E_{osc}+PE+LS=1+1+1=3$ . 3-1=2, 1.11X2=2.2 MeV **Neutron** is then  $\frac{1}{2}$  Deuteron (since we then dropped  $2P_{3/2}$ ).:3/2=1.5 MeV Multiple nucleons constrain inside each Deuteron to PE+SHM, so no LS. Excited state=SHM#=2 So Bond#XSHM#X(SHM+PE)+LS=BE If straight line then LS=0 He3. Straight line Bond#XSHM#X(SHM+PE)=2X2X2/3=8/3 P-n-P 3(8/3)+1/3=8, 8-1=7, 7X1.11=7.6 Mev **Tritium.** Not straight line so Bond#XSHM#X(SHM+PE)+LS= **BE**. n- n Extra electron so 3(8/3)+3/3=9, 9-1=8, 8X1.11=8.6 Mev. Р **He4** For rotation SHM=rot=3/3, 1D rot for each deuterium dumbbell so 2/3+1/3=3/3=16X2X(2/3)+3/3=24/3+1=8+1=9 9+9+9=27. P - N 27-2=25 25X1.11=27.75, 28/4=7 P-N, magic# Р Lithium. Note 2 more bonds so +2PE1. Subtract 3n and 3P 10+10+10=30. 30-3=27, 27X1.11=**30** 30/6=55 other 4nucleon 2P isotopes and isotones.eg., **Carbon** is 3 He4 -2P so 28+1=29, 3X29=87, 87/12=7.3. Rest of even #nuclei multiples of He4 and D, odd number add 1. Instead of N in  $\frac{h}{\omega}(N+\frac{1}{2})$  going up as in the shell model the N in  $\frac{Nh}{\omega}(1+\frac{1}{2})$  and  $\frac{Nh}{\omega}(2+\frac{2}{2})$ 

goes up (as well as PE and L). So we still have the same spherical harmonics (S.P.D.F.G...) states and the binding energy per nucleon is approximately constant from carbon onward.

## Neutron

If the deuteron broke up most of the energy is be in the single proton 1.3MeVand the remaining 1/3 would be in the 2P<sub>1/2</sub> state .78MeV. Add .511 for the electron and get 1.3Mev, the excess mass of a neutron above a proton.

**Radius Of Proton**: Average both the  $\cos^2$  edge ((1/2)r) and axial width (near zero)

 $(2.81713/2+0)/\sqrt{2}=.98F=r_{p}$ . But the spectroscopic and scattering methods that are used to measure  $r_p$  do not assume a singularity physics at 2.8F/2= $r_{Hp}/2$  so allow for the slightly larger radius in practice proton radius=  $r_p = 1.1F$ . The shell model sect.6.12 also requires this average radius and so has that inner and outer metric (sect.6.11, 3.1) as in sect 3.1 just like all solutions of eq. 2AI.

1.1.1 deuteron 3,3,3 He3 3,4,3 Tritium. (extra e) 9,9,9 He4 10,10,10 Lithium 30,30,30 C

We actually understand the nucleus of the atom from first principles this way. For history of the alternative Shell model, also see study by A.E.S. Green in 1956).

## 10.8 High (>100Gev) Energy Solutions

Note at high energy the electrons in the  $2P_{3/2}$  lobes (e.g., udd) would appear stationary, not averaged blob (density distributions). We are back to having single 'e' (not fractional "e") scatterers again. Thus at very high energies (>100GeV) single *e* (not fractional charge) should once again dominate scattering and we should no longer see these "jets" (which in the above context is mere P wave scattering) caused by higher probability emission in these trifolium lobe directions. Also note that  $r_H$  in  $\kappa_{00}$  is a hard shell and therefore Van der Waals type liquid equation of state at >100Gev energies. Note by the way that the 6<sup>th</sup> 2P resonance is observable at these energies.

Let  $\langle A' |$  represent the outgoing scattering wave immediately after a incident plane wave scatters off V. Let  $|A\rangle$  be the  $2P_{3/2}$  hyperon state for r=r<sub>H</sub> having the V. Thus at r=r<sub>H</sub> V itself will have the  $2P_{3/2}*2P_{3/2} = \psi^*\psi$  trifolium shape and thus commute with  $|A\rangle$  since they constitute the same structure ( $2P_{3/2}$  commutes with itself). So since V commutes with  $|A\rangle$  then  $\langle A'$  also is a  $2P_{3/2}$ state or we have  $\langle A' | V | A \rangle = 0$  and so no scattering into such states. Thus a type of 'P wave scattering' results from an incident plane wave. Thus we explain the origin of the 'jets' that are otherwise ascribed to scattering off quarks.

Note that when the mean free path d during the interaction time is very short (d<<(1/3)  $2\pi r_{\rm H}$ ) there is no more smearing between the  $2P_{3/2}$  lobes and we have scattering off of independent point particles and the  $2P_{3/2}$  state ceases to be relevant in the scattering and so the jets disappear. (jet quenching). Thus at extremely high energy the scattering is from charge e (not 1/3e) again and there are no more jets above top energy. LEP actually observed this effect just before it was shut down.

#### **10.9 Charge Independence Of The Strong Interaction**

It is well known that the strong interaction is approximately the same magnitude between Neutron-Proton, Neutron-Neutron and Proton-Proton pairs and thus is 'charge independent'. Also note our theory deals with electrons only which only has charge dependence if certain QM effects are ignored. But recall the orthogonality of S and P states as in  $\langle S|P \rangle = 0$ ,  $\langle S|S \rangle = 1$ ,  $\langle P|P \rangle = 1$  given all the superscript and subscript substates (e.g.,S and m) are the same as well in the bra and kets. The ordinary nuclear interaction here is due to a covalent bond (sharing electrons) which is also a very strong interaction (bond) at r=r<sub>H</sub> and is dependent on the spin S and m state and not so much on the sign of the charge. Thus these QM (valence, spin) effects are very strong at r=r<sub>H</sub>. Thus the charge independence of the strong interaction is really an S state *in*dependence and 2P<sub>3/2</sub> state dependence at r=r<sub>H</sub> of a 2P<sub>3/2</sub> structure interacting with an S state.



### Fig.5

There are no gauges required in this theory and the QCD SU(3) is such a gauge. We have found that hadrons are excited states composed of these half integer spherical harmonic lobes.

### **Chapter 11 Scattering Cross-Sections**

From the energy component of polarized representation of equation 8.1, 8.2: and using iterated (as in bosonic) section 9.13  $E^2=p^2c^2+m^2_{o}c^4=$ 

$$E^{2} = \left(\frac{1}{\sqrt{\kappa_{oo}}}\right)^{2} = \frac{1}{1 - \frac{r_{H}}{r}} = \frac{r}{r - r_{H}} = \frac{r}{r - r_{H}} - \frac{r - r_{H}}{r - r_{H}} + 1 = \frac{r_{H}}{r - r_{H}} + 1 = V + k$$
(10.8.1)

Note the resemblance of  $E^2=p^2c^2+m^2{}_{o}c^4$  to the Schrodinger equation if the  $E^2=k^2+1/(r-r_H)+1$  of equation 10.8.1 is substituted into it. We interpret this equation as representing a bounded volume with energy E=V+k therefore allowing us to use that V in the usual Gauge theory method and so substitute it into the ordinary Dirac equation as gauge force. term. So we use  $1/(r-r_H)$  instead of 1/r. in the Dirac equation S matrix.

We use the equation 4.1 source and proceed in the usual way of Bjorken and Drell (here  $1/r \rightarrow 1/(r-r_H/2)$  to construct the one vertex S matrix for the new Dirac equation 9. Recall the  $\frac{1}{2}$  came from the square root in equation 4.1. Thus the k in the integrand denominator is found from the result of our V=-1/(r-r\_H/2) potential in equation 10.8.1 instead of the usual Coulomb potential 1/r in the large r limit (so a free electron otherwise):

$$S_{if} \equiv iZ \frac{1}{V} \sqrt{\frac{m^2}{E_f E_i}} \overline{u}(p_f, s_f) \gamma^o u(p_i, s_i) \int_0^\infty \frac{e^{i(p_f - p_i)x}}{r - r_H} dx^4 \equiv if(u) \int_0^\infty \frac{e^{i(p_f - p_i)x}}{r - r_H} dx^4 \qquad (16.5)$$

rescaling  $r \rightarrow r' + r_H = r$  and  $t \rightarrow t' + (r_H/c) = t$  to minimize the resonance energy in  $p_f - p_i$ . We then obtain:

$$S_{if} \equiv iZ \frac{1}{V} \sqrt{\frac{m^2}{E_f E_i}} \overline{u} (p_f, s_f) \gamma^o u(p_i, s_i) e^{ir_H q} \int_0^\infty \frac{e^{i(p_f - p_i)x}}{r} dx^4 \equiv if(u) e^{iqr_H} \int_o^\infty \frac{e^{i(p_f - p_i)x}}{r} dx^4$$
For so:  

$$S_{if} \equiv if(u) \left[ \int_0^\infty \frac{e^{i(p_f - p_i)(x' - r_H/2)}}{|x'|} dx^4 - \int_0^\infty \frac{e^{i(p_f - p_i)(x' + r_H/2)}}{|x'|} dx^4 \right] \approx if(u) e^{i(r_H/2)q} \left( \int_0^\infty \frac{e^{i(p_f - p_i)(x')}}{|x'|} dx'^4 - 2\pi (r_H)^3 \right)$$

(10.8.6)

Note that  $\overline{u}(p_f, s_f)\gamma^o u(p_i, s_i) = (1 - \beta \sin^2 \frac{\theta}{2})$ =Mott scattering term with the e<sup>i(rH/2)q</sup> our

resonance term. The other left side coefficients and reciprocal  $|\boldsymbol{x}|$  part of  $S_{ij}$  comprise the well known Rutherford scattering term

 $d\sigma/d\Omega = [(Z_1Z_2e^2)/(8\pi\epsilon_omv_o^2)]^2 \csc^4(\phi/2) = 1.6X10^4(\csc(\phi/2)/v_o)^4$ . (Note that equation 10.8.6 applies to the  $2P_{1/2}$ - $2P_{3/2}$  state electron-electron interaction (i.e., neutron) below). Here  $p_f$ -p=q. Note in equation 10.8.6 the factor ie<sup>ikq</sup>=i(coskq+isinkq). Here we find the rotational resonances at the  $2P_{3/2}$  r=r<sub>H</sub> lobes associated with maximizing the imaginary part which is icoskq to obtain absorption scattering (at kq= $\pi$ ), which here will then be the masses exchanged in inverse beta decay. Also a solution to the Dirac component is always a solution to equation 14.1 (but not vice versa) if we invoke an integer spin in this resonance term. Here also the p part uses the old De Broglie wave length to connect to the p=h/ $\lambda$ . In that regard recall that hv/c= h/ $\lambda$ =p and for a DeBroglie wave fundamental harmonic resonance we have  $\lambda_{rot}=2\pi r$  for a stationary particle of **spin 1**=L (ambient E&M field source gives L=1 De Broglie).

Coulomb scattering of electrons, taking account spin-spin scattering

$$\frac{d\sigma}{d\Omega} = \frac{4Z^2 \alpha^2 m^2}{2|q|^4} Tr \gamma_o \frac{(\mathfrak{p}_i + m)}{2m} \gamma_o \frac{(p_i + m)}{2m}$$

$$\frac{d\sigma}{d\Omega} = \frac{Z^2 \alpha^2}{4p^2 \beta^2 \sin^4(\theta/2)} \left(1 - \beta^2 \sin^2 \frac{\theta}{2}\right)$$
  
Mott scattering, relativistic correction to Rutherford scattering  
Ultrarelativistic electron scattering: Electron rest mass m neglected.  
$$\frac{d\sigma}{d\Omega} = \frac{\alpha^2}{4E^2} \left(\frac{(\cos^2(\theta/2)) - (q^2/(2M^2))(\sin^2(\theta/2))}{(\sin^4(\theta/2))(1 + (2E/M))\sin^2(\theta/2)}\right)$$

$$\frac{d\sigma}{d\Omega} = \frac{\alpha^2}{4E^2} \left( \frac{(\cos^2(\theta/2)) - ((q^2/(2M^2))\sin^2(\theta/2))}{((\sin^4(\theta/2)))(1 + (2E/M)\sin^2(\theta/2))} \right)$$
(10.8.7)

m/E<<1, m<sup>2</sup>  $\rightarrow$ 0, q<sup>2</sup> =(p<sub>f</sub>-p<sub>i</sub>)<sup>2</sup>= -4EE'sin<sup>2</sup>( $\theta$ /2). Proton behaves like a heavy electron of mass M. E $\rightarrow$ 1/ $\sqrt{\kappa_{00}}$ =1/ $\sqrt{(1-\epsilon-\Delta\epsilon-r_H/r)}$ .

For forward scattering  $\theta \approx w/\beta \approx 0$  in  $\sin^2(\theta/2) \ll 1$  in the below figure 6. So  $d\sigma/d\Omega = d(1/E^2)/dE^2(1+\text{tiny})/\text{tiny} = ((d(1/t)/dt)(1+\text{tiny})/\text{tiny} = -(1/t^2)(1+\text{tiny})/\text{tiny}.$  $t=E^2 = (\text{energy transfer})^2.$ 

Ultrarelativistic dependence of 10.8.7 new pde electron differential cross-section on  $1/E^2$ . E=energy. Recall in this theory this should also be the energy dependence of ultrarelativistic proton-proton scattering since protons are made of electrons (in my work) and at very high energies(E>>150GeV) the electron cloud binding energies in the proton don't matter anymore (that Paschen back binding energy starts becoming negligible at TeV energies): we have free electrons hitting free electrons once again.

For energy transfer t on left side graph (fig.6)  $d\sigma/dt \propto d(1/E^2)/dE^2=d(1/t)/dt=-1/t^2$ . t=E<sup>2</sup>. Energy transfer  $\sqrt{t}$  is proportional to 1/p. But p<sup>2</sup> is proportional to area which is then proportional to  $1/\Delta E^2=1/t$ . So  $\sigma=$ Area  $\propto 1/\Delta E^2=1/t$ .  $\Delta\sigma/\Delta\sqrt{s}=(1000nb-10nb)/1GeV$ . But this is my equation 1 in figure 6 (also eq.10.8.7) for near forward elastic scattering. For 1GeV $\approx 1\sqrt{s}$  then this  $\Delta\sigma/(1)$  is a measure of  $d\sigma/dt$  on the left side forward scattering elastic energy transfer graph since 1<sup>2</sup>=1. But square root energy transfer  $\sqrt{t}$  in a scattering event for a beam at a specific energy (let's say at 13KeV) is also the abscissa of that big graph (on the left of the totem figure 6). So it is possible to get from total cross-section  $\sigma$  of electron scattering verse energy  $\sqrt{s}$ :  $\sigma/\sqrt{s}$ ) to  $d\sigma/dt$  vs t where t=(energy transferred)<sup>2</sup> at least at (literally) ONE (1GeV) energy transfer. The fact that LHC totem measures elastic forward scattering thereby made it possible to test this theory (eq.2, 1.11, new pde) at 13TeV (and  $\sqrt{s}\approx 2$ ), the very highest energy particles that mankind can produce. I could estimate from LHC data the asymptotically infinite beam energy transfer (curve) energy (red line) and compare it with my own  $\sigma/\sqrt{s}$  at  $\sqrt{s}\approx 2$ .From the graph of my equation 1 (10.8.7):

1000nb at  $\sqrt{s=1.5GeV}$ 

100nb at  $\sqrt{s}=2GeV$ 

10nb at  $\sqrt{s}=2.5$ GeV

But that curve of eq.1 in figure 6 is for one eq.2 electron scattering off of one equation 2 (new pde) electron. Since there are 3 such electrons in each of the two protons you must multiply by 9 to get  $d\sigma/dt$ .

On my QED graph of my eq.1 had  $\sigma \approx 100$ nb at about  $2 \approx \sqrt{s}$ . So multiply by 9 and get  $\Delta \sigma / \Delta s \approx (1000$ nb-10nb)/(( $(1.5-2.5)\sqrt{s}$ )= $10^{-3}$ mb/1GeV·But  $1^2 = 1$  and there are 3 electrons per proton so we

multiply by 2:  $9X10^{-3}$ mb/ $1^{2}$ GeV<sup>2</sup> $\approx 10^{-2}$ mb/ $1^{2}$ GeV<sup>2</sup> which is sitting in approximately the same t $\approx 2$  spot on the left side  $d\sigma/dt$  vs t graph of Fig.6.

So we proved from the data that a ultrarelativistic proton-proton scattering event (~13TeV) is equivalent to 6 free electrons scattering off each other with the electrons obeying (2AI) equation 9, the new pde. Thus the hadron theory that should be used is 2+2+2 at r=r<sub>H</sub>, not quark theory. Note also the cusp is at the proton reduced mass here. It is where (~.5GeV) binding energy must be added to break the electrons apart in the head on collision which takes away from the elastic scattering energy transfer. So we must apply this theory at much higher energies (eg., $\sqrt{s}\approx 2$ ). Quark theory (QCD) implies some kind of exponential dependence which is not seen in this scattering data.

Hard Shell Scattering Peak Of d $\sigma$ /ds Implies Protons Made of Electrons, not Quarks The electron radius at 2.8X10 <sup>-15</sup>(m<sub>e</sub>/m<sub>t</sub>)=8.1X10<sup>-19</sup>m provides the hard shell cross-section limit. For colliding beams we have an additional factor of 2 here.  $2(m_{\tau}+m_{\mu})/m_{e})m_{p}=6.91347$ TeV. There 3 are electrons in the proton so the proton energy is 3X6.913Tev=20.74TeV $\approx$ 21TeV. So the d $\sigma$ /ds should level off at proton energy 21TeV. This is in analogy with the Q= $\sigma/\pi r^{2} \propto d\sigma/ds$ (1/s $\propto \lambda$ ) r= $\lambda$  peak of Mie scattering theory.



We thereby discovered that both the dct/dt values and functional dependence imply a high energy proton acts like 3free electrons of eq.9 (in the context of collisions, not Brehmstrahlung). This implies a proton is really two positrons and an electron

#### Analogy to Mie scattering



Extinction efficiency (Q<sub>e</sub>) as a function of diffraction parameter x (=  $2\pi r/\lambda$ ). Analogy of Mie scattering with Q= $\sigma/\pi r^2$ . Here the lepton hard shell is at r=2X10<sup>-19</sup>m. Note analogy of leveling off d $\sigma/ds \lambda$ =r (i.e.,x=6) as at x=21TeV for LHC.

## Totem

LHC totem forward scattering gives elastic scattering cross-sections for high energies and so small scattering angles. We choose  $\theta=1^{\circ}$ .

$$\frac{\left(\frac{d\sigma}{d\Omega}\right)_{R}}{\left(\frac{d\sigma}{d\Omega}\right)_{R}} = (3)\frac{\alpha^{2}}{8E^{2}}\left[\frac{1+\cos^{4}(\theta/2)}{\sin^{4}(\theta/2)} - \frac{2\cos^{4}(\theta/2)}{\sin^{2}(\theta/2)} + \frac{(1+\cos^{2}\theta)}{2}\right] = 1.43X10^{-43}\left(\frac{1-.99985}{5.8X10^{-9}} - \frac{2(.99985)}{7.6X10^{-5}} + \frac{1+.9997}{2}\right) = 1.43X10^{-43}(2.626X10^{4} - 2.631X10^{4} + .99985) = (1.43X10^{-43})52.788$$

=7.55X10<sup>-42</sup> =7.55X10<sup>-14</sup> barns/steradian.

This result might be found in the totem data archives.

LHC totem forward scattering gives Coulomb scattering cross-sections for high energies for larger (but still small) scattering angles. We choose  $\theta=3^{\circ}$ . So

$$\frac{d\sigma}{d\Omega} = (3) \frac{Z^2 \alpha^2}{4p^2 \beta^2 \sin^4(\theta/2)} \left(1 - \beta^2 \sin^2 \frac{\theta}{2}\right) =$$

$$= (3) \frac{(1)5.33X10^{-5}}{4(1.04X10^{37}4.7X10^{-7})} (1 - 1^2 6.85X10^{-4}) = 8.18X10^{-38}m^2 =$$

=8.17X10<sup>-10</sup>barns/steradian. This result might be found in totem data archives data.

#### **Meson Multiplets**

"tetra quarks" are merely two mesons bound together! They can bind together more deeply if the components of the mesons themselves are bound individually to the components of the other meson giving more mass, section 8.11.

In this theory (DavidMaker.com, Ch.9-10) this is called singlet and doublet states with one bound with more binding energy than the other for those heavy upper 2P Paschen Back states. So these look like heavy and light tetraquark states but they are not, they are merely two types of meson binding states. You could predict the energies from the Paschen Back effect associated with those large plate fields, section 8.11.

## References

Pugh, Pugh, *Principles of Electricity and Magnetism*, 2<sup>nd</sup> Ed. Addison Wesley, pp.270 Bjorken and Drell, *Relativistic Quantum Mechanics*, pp.60 Sokolnikoff, Tensor Analysis, pp.304

## 11.1. W Compton Wavelength Region

Recall in appendix A m<sub>e</sub> Source Term at r=r<sub>H</sub> Inside Angle C. Analogously from **2AC** we get with the eq.3 doublet  $\varepsilon \pm \varepsilon$  the Proca equ (3) *neutrino and electron*  $\Delta \varepsilon$  at r=r<sub>H</sub>. As in sect.6.13 in  $\kappa_{00}$  we normalize out the muon  $\varepsilon$ . So we are left with the electron  $\Delta \varepsilon$  in  $\kappa_{00} = 1 - [\Delta \varepsilon / (1 \pm 2\varepsilon)] + [r_H(1 + ((\varepsilon \pm \varepsilon)/2))/r]$  from the two above rightmost (Proca) diagrams. So Source =  $E_{ZW} = \frac{1}{\sqrt{\kappa_{00}}} = \frac{1}{\sqrt{1 - \frac{\Delta \varepsilon}{1 \pm 2\varepsilon} - \frac{r_H e}{r}(1 + (\varepsilon \mp \varepsilon)/2)}} \approx \frac{1}{(1 \pm \varepsilon)\sqrt{\Delta\varepsilon}}$ , at  $r = r_{H_e}$  + is for Z

and – is for W. So W (right fig.) is a single electron  $\Delta \varepsilon + \nu$  perturbation at  $r=r_H=\lambda$  (since  $m_e$  ultra relativistic): So  $H=H_o+m_ec^2$  inside V<sub>w</sub>.  $E_w=2hf=2hc/\lambda$ ,  $(4\pi/3)\lambda^3=V_w$ . For the two leptons

$$\frac{1}{V^{1/2}} = \psi_e = \psi_3$$
,  $\frac{1}{V^{1/2}} = \psi_\nu = \psi_4$  . Fermi 4pt=

$$2G \iiint_0^{r_w} \psi_1 \psi_2 \psi_3 \psi_4 \, dV = 2G \iiint_0^{r_w} \psi_1 \psi_2 \frac{1}{V^{1/2}} \frac{1}{V^{1/2}} V$$

 $= 2 \iiint_{0}^{r_{W}} \psi_{1}\psi_{2}G \equiv \iiint_{0}^{V_{W}} \psi_{1}\psi_{2}(2m_{e}c^{2})dV_{W} = \iiint_{0}^{V_{W}} \psi_{1}(2m_{e}c^{2})\psi_{2}dV_{W}$ (A3) What is Fermi G?  $2m_{e}c^{2}(V_{W}) = .9X10^{-4}Mev-F^{3} = G_{F}$  the strength of the weak interaction.

## Derivation of the Standard Model But With No Free Parameters

Since we have now derived  $M_W$ ,  $M_Z$ , and their associated Proca equations ,  $m_\mu$ ,  $m_\tau$ ,  $m_e$ , etc., Dirac equation figure 2 part.1),  $G_F$ ,  $ke^2$ , Bu, Maxwell's equations, etc. we can now set up a Lagrangian density that implies all these results. In thisFormulation  $M_z=M_W/\cos\theta_W$ , so you find the Weinberg angle  $\theta_W$ ,  $gsin\theta_W=e$ ,  $g'\cos\theta_W=e$ ; solve for g and g', etc., We will have thereby derived the standard model from first principles (i.e., postulate1) and so it no longer contains free parameters!

Thus we have the interaction  $\varepsilon$  operating in W radius using the doublet of Ch.3.

In general then we have obtained an ortho triplet state here since we are merely writing the Clebsch Gordon coefficients for this addition of two spin  $\frac{1}{2}$  angular momentums:

 $|\frac{1}{2},\frac{1}{2},0,-1\rangle,...|\frac{1}{2},\frac{1}{2},0,0\rangle,...|\frac{1}{2},\frac{1}{2},0,1\rangle,...$  or  $+W Z_{o}, -W$ .

Anyway, this small S matrix involves the neutrino and so can allow spin 1/2 neutrino emission jumps instead of just the usual E&M spin 1 jumps. 100km/sec metric quantization translates to a neutrino rest mass of .165eV.

## 11.2 Excited Z States

## Put me in Equation 6.4.1

The beautiful thing to be noted here is that for the doublet resonance with the  $2P_{3/2}$  lobe at r=r<sub>H</sub> that minimizes energy you get the spin 1 W and Z and the value of the Fermi G! We have also shown that this doublet interaction corresponds to the exchange of massive spin 1 particles (recall spin  $\frac{1}{2}$  s forbidden by that j-1/2 factor).

## 11.3 Probability for 2P<sub>3/2</sub> Giving One Decay 1S Product at r≈r<sub>H</sub> In W Region

In equation 4.12 we note that invariance over  $2\pi$  rotations using  $(1+2\epsilon)d^2\theta$  does not occur anymore thus seemingly violating the conservation of angular momentum. To preserve the conservation of angular momentum the additional angle  $\epsilon$  must then include its own angular momentum conservation law here meaning intrinsic spin<sup>1</sup>/<sub>2</sub> angular momentum in the S state case and/or isospin conservation in the  $2P_{3/2}$  case at r=r<sub>H</sub>. In any event we must also integrate to C= $\epsilon$ . Here we do the E&M component decay given by equation 3.2.

Plug in  $S_{\frac{1}{2}} \propto e^{i\phi/2}$ ,  $\frac{1}{2}(1-\gamma^5)\psi=\chi$  into the 4pt. interaction integral. In that regard note that the expectation value of  $\gamma^5$  is proportional to  $v \propto$  Heisenberg equation of motion derivative of  $2P_{3/2} \propto e^{i(3/2)\phi}$ . We integrate <lepton|baryon> over this W exchange region where we note  $(\sim 1/100)F$  for 90Gev particle, so  $dV=((1/100)F)^3$ =VolW. Also  $ck_o=\epsilon=106$ Mev from section 2.1. From Ch.3 on the vacuum constituents e and v we note that  $\iiint d\tau=Vol, \chi$  is defined as the vacuum eigenfunction. Vacuum expectation sect.B2:  $\Sigma < |va_M\rangle |\epsilon| < vac_{M+1}| >= <|$  $\iiint \psi^* v \epsilon \chi_e dV| >= <|Pot| >= \epsilon Volume of W. Recall also that appendix A implies that the W and the Z are composites. This application of eq.2 for example applies to the <math>2P_{\frac{1}{2}} - 2P_{3/2}$  electron-electron scattering state inside the neutron  $< Proton 2P_{3/2}|Pot|Neutron 2P_{\frac{1}{2}} - 2P_{3/2}\rangle$ . Plug in  $S_{\frac{1}{2}} \propto e^{i\phi/2}$ ,  $\frac{1}{2}(1-\gamma^5)\psi=\chi$  Also we can get a weak, strangeness changing (second term below), decay from a  $2P_{\frac{1}{2}} - 2P_{\frac{3}{2}} > m_p$  to the S state branch equation. eq.2 expectation values in the 4pt..

= $\Sigma$ <lepton|vac>| $\epsilon$ |<vac|baryon>=Fermi interaction integral = $\int \psi_1 * \psi_2 \psi_3 k_0 c \chi dV$ =  $\iint \psi_1 * (\epsilon \text{VolW}) \chi dV = \iiint \psi_1 * (\epsilon \text{VolW}) \chi dV$ . Also dV=dAd $\phi$ =Kd $\phi$ .

So the square root of the probability of being in the final state is equal to the Fermi integral=  $\int \psi_1 * (Pot) \chi dV = \int \int \psi_1 * \psi_2 \psi \varepsilon \Delta V_W \chi dV =$ 

$$= \int K \left\langle e^{i\phi/2} | \left( \varepsilon \Delta V_{W} \right) | 1 - \left( \gamma^{5} i e^{i(3/2)\phi} \right) \right\rangle dt c dV = \int K \left[ \varepsilon \left( \frac{F}{100} \right)^{3} \right] \left\langle e^{i\phi/2} - \gamma^{5} i e^{i4/2\phi} \right\rangle d\phi$$

$$\varepsilon VolW = \int K \left[ \varepsilon \left( \frac{F}{100} \right)^{3} \right] \left\langle e^{i\phi/2} - \gamma^{5} i e^{i4/2\phi} \right\rangle d\phi$$

$$= KG_{F} \int \left\langle e^{i\phi/2} - \gamma^{5} i e^{i4/2\phi} \right\rangle d\phi = KG_{F} \left( \frac{2e^{i\phi}}{i} |_{0}^{2\pi+C} - \frac{2\gamma^{5} e^{i4\phi}}{i4} |_{0}^{2\pi+C} \right)$$
(10.8.7)

with  $\langle initial|c|final\rangle^2 \approx$  transition probability as in associated production with the separate 2P proton ground state transition being the identity ( $\Delta S=0$ ). Factoring out the 2 and then normalizing 1 to .97 simultaneously normalizes the 1/4 to .24 in section 3.2. With this normalization we can set  $\cos\theta_c=.97$  and  $\sin\theta_c=.24$ . Thus we can identify  $\theta_c$  with the *Cabbibo angle and we have derived its value*. We can then write in the weak current sources for hadron decay the **VA structure**:  $|\cos\theta_C-\gamma^5\sin\theta_C|$ . Thus with the above Cabbibo angle and this CP violation and higher order  $(r_H/r)^n$  terms in section 3 we have all the components of the CKM matrix. Note we have also derived the weak interaction constant G<sub>F</sub> here.

Given the role  $\varepsilon$  plays here in decay we find the expectation value of energy  $\varepsilon$  within the S matrix scattering region in chapter 10.

Recall from section 1.2 the possible mixing of real and imaginary terms in that energy coming out of that first order Taylor expansion. There we found the 1+x and 1-x solutions cancel and we could ignore the 1+1=2 term as it is still a flat metric.

Also there are still extra terms provided by the 'small' higher order  $r^2$  terms in that Taylor expansion so that "higher and lower" than the speed of light mixed condition still can exist (for  $\Delta G \neq 0$ . See end of section 4.6 and 10.8.6). In that regard note for the next higher order Taylor term at largest curvature  $d^2(1/k_{rr})/dr^2$  is large negative and  $r^2$  is positive implying a net negative term and therefore a neutral charge (see case 2, of section 19.6)! In that case the perturbative squared r term appears to overwhelm the rest since the lower order terms then cancel. Note from

the above we put these neutral conditions also into that decay since net charge is zero in the Cabbibo angle derivation. This then appears to be the beta decay condition *where the neutrino (higher than c) and the electron, (lower than c), decay from this neutral particle condition* (bottom of section 10.8). The beginning 2P<sub>3/2</sub> ground state still exists however in the respective Cabbibo angle calculation. Thus those real and imaginary terms coming out of that Taylor expansion provide the explanation for beta decay.



#### Fig.7

#### me Source Term Inside Angle

See section 0.2 and B2 So W is a single electron  $\Delta \varepsilon$ , v perturbation at r=r<sub>H</sub>: H=H<sub>o</sub>+2m<sub>e</sub>c<sup>2</sup> inside  $V_w$ .  $E_w=2hf=2hc/\lambda$ ,  $(4\pi/3)\lambda^3=V_w$ . For the two leptons  $\frac{1}{V^{-1/2}} = \psi_e = \psi_3, \frac{1}{V^{-1/2}} = \psi_v = \psi_4$ . Fermi 4pt= G  $\iiint_0^{V_W} \psi_1 \psi_2 \psi_3 \psi_4 \, dV = G \iiint_0^{V_W} \psi_1 \psi_2 \frac{1}{V^{1/2}} \frac{1}{V^{1/2}} V$  $= \iiint_0^{V_W} \psi_1 \psi_2 G \equiv \iiint_0^{V_W} \psi_1 \psi_2 (2m_ec^2) dV_W$  (A3)

What is Fermi G?  $2m_ec^2(V_W)/F^3 = .9X10^{-4}Mev-F^3 = G_F$  the strength of the weak interaction. Next we plug the respective  $\psi$ s into  $\psi_1, \psi_2$  in sect.B2. In that regard the expectation value of  $\gamma^5$  is speed and varies with  $e^{i3\phi/2}$  in the trifolium. The spin<sup>1</sup>/<sub>2</sub> decay proton  $S_{1/2} \propto e^{i\phi/2} \equiv \psi_1$ , the original  $2P_{1/2}$  particle is chiral  $\chi = \psi_2 \equiv \frac{1}{2}(1-\gamma^5)\psi = \frac{1}{2}(1-\gamma^5e^{i3\phi/2})\psi$ . Initial  $2P_{1/2}$  electron  $\psi$  is constant. Plug these terms into equation B2 =  $\iiint_0^{V_W} \psi_1 \psi_2 (2m_ec^2) dV_W =$   $\iint \psi_{S1/2} * (2m_ec^2V_W))\chi dV =$  $K \int \langle e^{\frac{i\phi}{2}} [\Delta \varepsilon V_W] (1-\gamma^5e^{i\phi\frac{3}{2}})\psi \rangle d\phi = KG_F \int \langle e^{i\phi/2} - \gamma^5ie^{i4/2\phi} \rangle d\phi = KG_F \left( \frac{2e^{i\phi}}{i} |_0^{2\pi+C} - \frac{2\gamma^5e^{i4\phi}}{i4} |_0^{2\pi+C} \right)$ 

with VA <initial|c|final><sup>2</sup>  $\approx$  transition probability as in associated production. Factoring out the 2 and then normalizing 1 to .97 simultaneously normalizes the 1/4 to .24 in Ch.3. With this normalization we can set  $\cos\theta_c$ =.97 and  $\sin\theta_c$ =.24. Thus we can identify  $\theta_c$  with the **Cabbibo** angle and we have derived its value.

 $\Gamma = [G^2/(4\pi N)] m_i^{\alpha} P^{|m|}(x)|^2, x = \Sigma m_j/m_t Eg. 1/\tau_{\mu} = [G^2 m_{\mu}^5/(192\pi^3)](1-m_e^2/m_{\mu}^2)^6.$ 

# r<r<sub>H</sub> Application: Rotational Selfsimilarity With pde Spin: CP violation 12.1 Fractal selfsimilar spin

The fractal selfsimilarity with the spin in the (new) Dirac equation 2 implies a selfsimilar cosmological ambient metric (Kerr metric) rotation as well as in section 4.1. Thus there will be  $2d_{s_{\phi}}$  rotation metric cross terms with the dt (without the square) implying time T reversal nonconservation and therefore CP *non*conservation since CPT is always conserved. We thereby derive CP nonconservation from first principles: **CP nonconservation is a direct consequence of the fractalness**. This adds another matrix element of magnitude ~1/3800 (sect.6.3) for Kaon decays thus adding off diagonal elements to the CKM matrix.

Or for Kerr rotator use

$$ds^{2} = \rho^{2} \left( \frac{dr^{2}}{\Delta} + d\theta^{2} \right) + \left( r^{2} + a^{2} \right) \sin^{2} \theta d\phi^{2} - c^{2} dt^{2} + \frac{2mr}{\rho^{2}} \left( a \sin^{2} \theta d\theta - c dt \right)^{2},$$
(13.1)

$$\rho^{2}(r,\theta) \equiv r^{2} + a^{2}\cos^{2}\theta; \qquad \Delta(r) \equiv r^{2} - 2mr + a^{2}, \text{ or}$$
$$ds^{2} = dr^{2} + dt^{2} + 2dtdr + ..$$

In a polarized state ( $\theta = 0^{\circ}, 180^{\circ}$ ) in 25.3, 25.25 the off diagonal elements are proportional to  $\phi = (\phi + c)e^{-C}$ . Thus if the charge e is conjugated (C, e changes sign), if dr changes sign (P, parity changes sign) and dt is reversed (t reversal) then the ds quantity on the left side of equation 1.6 is invariant. But if dr (P) changes sign by itself, or even e and P together (CP) change sign then ds is not invariant and this explains, in terms of our fractal picture, why CP and P are not conserved generally. P becomes maximally nonconserved in weak decays as we saw in above. The degree to which this nonconservation occurs depends on the "a" (in eq.3.2.1) transfer <final lal initial> (equation 3.2) which itself depends on the how much momentum and energy is transferred from the S<sub>M+2</sub> to the S<sub>M+1</sub> fractal scales as we saw in this section. Recall chapter 5 alternative derivation of that new (dirac) equation pde (eq.2) **linearization of the Klein Gordon** equation(c=1,  $\hbar$ =1, m=1, eq.2):

$$\left(-\alpha_{1}i\frac{\partial}{\partial x_{1}}-\alpha_{2}i\frac{\partial}{\partial x_{2}}-\alpha_{3}i\frac{\partial}{\partial x_{3}}+\beta\right)\left(-\alpha_{1}i\frac{\partial}{\partial x_{1}}-\alpha_{2}i\frac{\partial}{\partial x_{2}}-\alpha_{3}i\frac{\partial}{\partial x_{3}}+\beta\right)=$$
(13.2)  
$$-\alpha^{2}\frac{\partial^{2}}{\partial x_{2}}-\alpha^{2}\frac{\partial^{2}}{\partial x_{3}}+\beta^{2}+2\sum_{i=1}^{3}\sum_{j=1}^{3}\alpha_{j}\alpha_{j}\frac{\partial}{\partial x_{j}}-\alpha_{3}i\frac{\partial}{\partial x_{3}}+\beta=$$
(13.2)

 $-\alpha_{1}^{2} \frac{1}{\partial x_{1}^{2}} - \alpha_{2}^{2} \frac{1}{\partial x_{2}^{2}} - \alpha_{3}^{2} \frac{1}{\partial x_{3}^{2}} + \beta^{2} + 2\sum_{\mu=1}^{2} \sum_{\nu=1}^{n} \alpha_{\mu} \alpha_{\nu} \frac{1}{\partial x_{\mu}} + 2\sum_{\mu=1}^{n} \beta i \alpha_{\mu} = -\frac{1}{\partial t^{2}}.$  This equals  $= c^{2} p_{1}^{2} + c^{2} p_{2}^{2} + c^{2} p_{3}^{2} + m^{2} c^{4} = E^{2} \text{ if the off diagonal elements zero which is the condition used}$ in the standard Dirac equation derivation of the  $\alpha$  s and  $\beta$ . Note that the off diagonal elements  $2\sum_{\mu=1}^{3} \sum_{\nu=1}^{3} \alpha_{\mu} \alpha_{\nu} \frac{\partial}{\partial x_{\mu}} \frac{\partial}{\partial x_{\nu}} + 2\sum_{\mu=1}^{3} \beta i \alpha_{\mu} \text{ are equivalent to the off diagonal elements in equation 5.1}$ (and are corrections to 5.2 in fact) so *are not* zero for parity and CP NONconservation in this context (in a rotating universe). So in the context of the Dirac equation the CP violation term  $e_{s}(dr)dt \rightarrow (dr/ds)(dt/ds) \rightarrow pE\chi \quad (after division by ds^{2}).$  Thus CP violation goes up as the square (pE) of the energy (so should be larger in bottom factories). The section 13.2 below

Cabbibo angle calculation (not rotation related however) is an example of how this method can give the values of the other terms in the CKM matrix. They arise from calculation of  $\langle Z \rangle$  between higher order m harmonics.

This section is important in that we see that CP violation is explainable and calculable in terms of perturbative effects on the ambient metric (and therefore the Dirac equation) of a rotating universe with nearly complete inertial frame dragging (eq. A6 in the E&M form), CP violation doesn't need yet more postulates as is the case with the GSW model. In fact the whole CKM matrix is explainable here as a consequence of this perturbation.

Note the orientation relative to the cosmological spin axis is important in CP violation. Integration of the data over a 3 month time (at time intervals separated by a sidereal day) is going to yield different CP violation parameters than if integration is done over a year.

## Miscellaneous 12.2 GIM Derivation

Recall in the GIM (Glashow, Iliopoulos, Maiami) hypothesis that u,d were a pair of left handed Fermion states as in V-A.  $d'=d\cos\theta_c+s\sin\theta_c$ ,  $s'=-d\sin\theta_c+s\cos\theta_c$  where  $\theta_c$  is the Cabbibo angle. Thus u,d are paired, s,c are paired, b,t are paired and we have the V-A transitions. Here we identify the new pde 2P state for  $r=r_H$  has  $P_x$ ,  $P_y$  and  $P_z$  states which split in energy due to that Paschen Back effect given those ultrarelativistic plates, into paired spin up and spin down states  $(P_x, P_x')$ ,  $(P_y, P_y')$ ,  $(P_z, P_z')$  analogous to the GIM (u,d), (s,c), (b,t). Here the spin orbit interaction (LS) coupling energy term is much stronger than the SS coupling term. So we have pairs of states J,M,M'> with  $P_x$  and  $P_y$  being orthogonal, except for those weak interaction V-A terms. The ds<sup>2</sup> to ds' transition is through the V-A term. Recall equation the 16.7  $|\int \chi_f^* G \chi_o dV|^2 =$ =transition probability of a ds<sup>2</sup> to a ds. of eq.B2 ( $\chi$ =.5(1- $\gamma$ <sup>5</sup>) $\psi$  with  $\psi$  in ds<sup>2</sup>,  $\chi$  in ds) for V-A Cabibbo angle transitions (transitions inside P<sub>X</sub> separately from P<sub>Y</sub> and separately from P<sub>z</sub>) where  $|1|-|\gamma^5|$  replaces the  $\cos\theta_c + \sin\theta_c$ . So in analogy for transitions between P<sub>X</sub> and P<sub>Y</sub>  $P_x$ 'cos $\theta$ c+ $P_y$ sin $\theta_c$ =Px',  $P_y$ '=- $P_x$ 'sin $\theta_c$ + $P_y$ 'cos $\theta_c$ . This is a first principles understanding of GIM thereby allowing us to derive the electroweak cross-sections (WS). Recall that dz=-1, 0 solution to eq.2 for C=0 implies dr<0 at least for small C. (low noise). because -1 is on the real r axis.

## 12.3 Normed Division Algebra, Octonians, E8XE8 and SU(3)XSU(2)XU(1) Basis Change

Note from the above that the new pde fractal theory generated the electron 2AI with mass, the near zero mass left handed neutrino 1.12. Recall also from above the  $\kappa_{oo} = \sqrt{(1-\epsilon - r_H/r)}$ . The W was generated from a nonzero ambient metric  $\epsilon$  in that S matrix derivation part of the metric coefficient  $\kappa_{\mu\nu}$ . Interestingly that Normed Division Algebra (NDAR) on the real numbers (as in:  $||Z_1||^*||Z_2||=||Z|||$ ) from equation 1 implies that octonians (and thereby the *largest normal Lie group* E8XE8) are also allowed. Recall we have that SU(2) Lie group rotation for the 0° extrema imbedded in a E8XE8 rotation since one of its subgroups being SU(3)XSU(2)XU(1). This is the only subgroup we can use because it is the one that only contains that SU(2).

## 7.3 Eigenstates

Recall the  $m_{\tau}=1$  was separable at 45° from the rest (of the eq.1.15 diagonal states) since ds is constant there for small rotations. So  $ds_{\tau}$  can be normalized.

The B field rotations are here reciprocals of the rotations in the Mandelbulbs since  $\kappa_{oo}=1-r_H/r$  and so  $r_H/r \rightarrow \xi dr$  for B field motion given smaller radius r means high energy  $\xi dr$ . So the *ortho* state is the smaller  $\epsilon$  Mandelbulb eigenstate and the *para* state is the larger  $\tau$  limacon Mandelbulb eigenstate.

## Meta Theory Of Couplings In SM

From the 1.15 diagonal on the Mandelbrot set:  $m=m_L=1+\epsilon+i\Delta\epsilon=m_\tau+m_\mu+m_e$ .  $mv^2/r_H=qvB(1)$   $\pi r_H^2B=\Phi o$  (2) v=c. Solve Eq.1 and eq.2 for q:  $q=mc\pi r_H/\Phi_o$ . (3) The effect of the *E field lines coming* together by Fitzgerald contraction  $\gamma$  imply a force increase that can be realized by invoking an *effective* charge increase  $e \rightarrow q'$ .or in V/dr'=Electric field with r=r\_H in dr'^2=\kappa\_{rr}dr^2=[1/(1-r\_H/r)]dr^2 (The charge e itself really does not change.). For  $m=m_{\mu}=\epsilon$  2P<sub>3/2</sub> state (So ultrarelativistic so E field line contraction.). From equation 3  $\mathbf{q}'=\mathbf{e}, \mathbf{H}\rightarrow \mathbf{e}^2 \ \mathbf{E} \& \mathbf{M}$  for the Nth fractal scale, Gravity for the N+1th fractal scale. For  $\mathbf{m}=\mathbf{m}_{\tau} + \mathbf{m}_{\varepsilon}$  as meson.  $2P_{3/2}$  state (so ultrarelativistic). From equation 3:  $\mathbf{q}'=46\mathbf{e}, \quad \mathbf{H}\rightarrow (46\mathbf{e})^2 \ \text{Strong Force.}$ 

For  $\mathbf{m}=\mathbf{m}_{e}=i\Delta\varepsilon + v$ , v small,  $2P_{\frac{1}{2}}$ ,  $dr'^{2}=\kappa_{rr}dr^{2}=1/(1-r_{H}/r)dr^{2}$ . So V/dr'=E small. From equation 3: q'=ie/200, H $\rightarrow$ iq'V so  $\psi(t)=e^{iiq'Vt}\psi(t_{o})=e^{-q'Vt}\psi(t_{o})$ .

exponential decay with a force  $q'^2=40000X$  *smaller* than the E&M. Weak interaction. dr' large allowing large uncertainty principle dr' for small nonrelativistic mass m<sub>e</sub> in  $(dr' \cdot m_ec) \geq h/2$ . This occurs for small externally observed dr and m<sub>e</sub>c in the 2P<sup>1/2</sup> state and 1S<sup>1/2</sup> state at r=r<sub>H</sub>. But these are decay states (PartII Sect.7.3). Given these strength and decay parameters we can alternatively integrate over the r<sub>c</sub> volume our W and Z particles to get the Fermi G 4pt coupling of weak interaction theory in the SM. W is then a virtual intermediary here. So we just derived all 4 forces from that diagonal on the Mandelbrot set.

**Calculations**: So for the Kerr mass ortho state  $(2^{nd} \text{ Mandelbulb}) (a/r)^2 = \epsilon + \Delta \epsilon$  (thus added to 1) at r=r<sub>H</sub>, for (N+1):  $m_{\epsilon}v^2/r=qvB$  so  $m_{\epsilon}v/(qB)=(1-2\epsilon)m_{\mu}c/qB=r$ . Thus  $(1-2\epsilon)m_{\mu}c/q=rB=(1-2\epsilon)(1.883X10^{-28})(299792458)/(1.6022X10^{-19})=.3525(1-2\epsilon)$  $\Phi o/(rB)=\Phi_o/[0.3525(1-2\epsilon)]=B\pi r^2/(rB)=\pi r$ .  $r_{H}=1.359X10^{-15}(2)/[(.3526)\pi(1-2\epsilon)]=2.805X10^{-15}m=e^2/m_ec^2$  for 2P states (eq.7.1).