

Part II

Incorporating New Pde Hamiltonian Into Heisenberg Equations of Motion

Chapter 9

9.1 Review of Derivation of Heisenberg's Equations of Motion Given equation 1.9

Here we multiply section 1.3 chain rule result $p\psi = -i\hbar\partial\psi/\partial x$ by ψ^* and integrate over volume to define the expectation value

$$\int \psi^* p_x \psi dV \equiv \langle p_x \rangle = \langle p, t | p_x | p, t \rangle \text{ of } p_x. \quad (9.1)$$

In general for any QM operator A we write $\langle A \rangle = \langle a, t | A | a, t \rangle$. Let A be a constant in time (from Merzbacher, pp.597). Taking the time derivative then:

$$\begin{aligned} i\hbar \frac{d}{dt} \langle a, t | A | a, t \rangle &= i\hbar \frac{d}{dt} \langle \Psi(t), A \Psi(t) \rangle = \left(\Psi(t), A i\hbar \frac{\partial}{\partial t} \Psi(t) \right) - \left(i\hbar \frac{\partial}{\partial t} \Psi(t), A \Psi(t) \right) \\ &= (\Psi(t), A H \Psi(t)) - (\Psi(t), H A \Psi(t)) = i\hbar \frac{d}{dt} \langle A \rangle = \langle A H - H A \rangle \equiv [H, A] \end{aligned}$$

In the above equation let $A = \alpha$, from equation 1.9 Dirac equation Hamiltonian H, $[H, \alpha] = i\hbar d\alpha/dt$ (Merzbacher, pp.597).

9.2 Derivation of Velocity of Expansion for $r < r_H$ from Heisenberg's Equations of Motion and Equation 1.9

The second and first integral solutions to the Heisenberg equations of motion (i.e., above $[H, \alpha] = i\hbar d\alpha/dt$) is:

$$\begin{aligned} \mathbf{r} &= \mathbf{r}(0) + c^2 \mathbf{p}/H + (\hbar c/2iH) [e^{(i2Ht/\hbar)} - 1] (\alpha(0) - c\mathbf{p}/H). \\ \mathbf{v}(t)/c &= c\mathbf{p}/H + e^{(i2Ht/\hbar)} (\alpha(0) - c\mathbf{p}/H) \end{aligned} \quad (9.2)$$

Chapter 10

$r < r_H$ Application implies Cosmology

10.1 $r < r_H$ Application of Equation 1.9 Using Sections 3.4 and 9.2 Results

Again solve the Heisenberg equations of motion $[H, \alpha] = i\hbar d\alpha/dt$ using the Hamiltonian H of equation 1.9 and equation 9.2 get again (also Merzbacher, pp.597)

$$r = r(0) + c^2 p/H + (\hbar c/2iH)[e^{(i2Ht/\hbar)} - 1](\alpha(0) - cp/H). \quad (10.1)$$

In the QM operator formulation of R_{ij} the equation 10.1 zitterbewegung oscillation makes GR **ungauged** since the harmonic gauge in $R_{ij}=0$ can be replaced with the zitterbewegung oscillation of equation 10.1 as in section 25.7. Also taking the bilinear operator expectation values of equation 10.1 (and the rest of the $R_{ij}=0$ equations) gives us the Maxwell equations in the weak field approximation (section 25.4) with the excess above coulomb potential giving the lamb shift after energy expectation values are taken for the $\psi_{2,0,0}$ eigenstates. After taking expectation values $\psi^* \psi = e^{-i\omega t} \psi^*(r, \theta, \phi) e^{i\omega t} \psi(r, \theta, \phi) = |\psi(r, \theta, \phi)|^2$ the time dependence of the zitterbewegung goes away, becomes random for the observer, as in the standard QM (but with our Maxwell equations and Ungauged GR of section 12.3 remaining).

Note also from equations 18.1, 18.2 that $(dt/ds)\sqrt{g_{00}} = 1/\sqrt{g_{00}}$ is where the energy was in the old Dirac equation. Thus $\langle H \rangle = E = 1/\sqrt{g_{00}}$ (with $g_{00} = 1 - r_H/r$) the imaginary part of ω *outside* ($r > r_H$) becomes the real part *inside* ($r < r_H$) the horizon since the square root of a negative number (here g_{00}) is imaginary. Thus $\omega \rightarrow i\omega$ and so the above complex conjugation does not occur and so there is “observed” motion (for $r < r_H$) with $1 + \sinh \omega t$ hyperbolic sine below in equation 10.1. So taking the imaginary part of r in equation 10.1 (setting $r(0)=0, p=0$) gives us the physics inside:

$$\begin{aligned} \text{IM}[(r(0) + c^2 p/H + (\hbar c/2iH)[e^{(i2Ht/\hbar)} - 1](\alpha - cp/H)] &= \text{IM}(\hbar c/2iH)[e^{(i2Ht/\hbar)} - 1]\alpha \\ &= \text{IM}(c/i\omega)[e^{(i\omega t)} - 1]\alpha = \text{IM}[-(ic/\omega)][(\sin \omega t) - 1]\alpha = \text{IM}(ic/\omega)[1 - (\sin \omega t)]\alpha. \end{aligned}$$

Recall that $\sin i\omega t = \sinh \omega t$. Thus with above $\omega \rightarrow i\omega$ (in eq.10.1) as we go inside the horizon then this equals

$$(c/\omega)[1 + \sinh(\omega t)]\alpha$$

(the αdx is used in the eigenvalue calculation in that earlier bilinear function sense).

Define $t=0$ at $r=0$ using time and phase setting: $\omega t \rightarrow \omega t + \omega t_0$ and note that

$\sinh(-.8814) \approx -1$. Therefore we have:

$$r = (c/\omega)[1 + \sinh(\omega t - .8814)]\alpha \quad (10.2)$$

Equation 10.2 bares a remarkable resemblance to the Lemaitre cosmological model scale factor expansion which also contains a $\sinh(\omega t)$ and is identical in form to it if the negative square root of λ (in the its scale factor terms) is taken. Recall that the Lemaitre model distinguished itself by including the effects of a cosmological constant with its

accompanying (recently discovered) accelerating ($d^2r/dt^2=a$) cosmological metric expansion implied by equation 10.2. So inside the horizon the motion is observed, not random anymore. Note that the from chapter 2 that ($\omega t-.8814$) substitutes for the $2Ht/\hbar$ in $g_{00} = e^{(2Ht/\hbar)}$.

Chapter 11

11.1 $r < r_H$ Application: Expansion of Zitterbewegung Cloud From $r = r_0[1 + \sinh(\omega t - .8814)]$: Gravity

Again given the structure of the Hamiltonian of Dirac equation 1.9 we see from the Heisenberg equations of motion applied to this Dirac equation that (Merzbacher's eq. 24.27) equation 10.2

$$r = (c/\omega)[1 + \sinh(\omega t - .8814)]\alpha \quad (11.1)$$

Again equation 10.2 bears a remarkable resemblance to the Lemaitre cosmological model scale factor expansion which also contains a $\sinh(\omega t)$ and is identical in form to it if the negative square root of λ (in the its scale factor terms, see section 12.2) is taken. Recall that the Lemaitre model distinguished itself by including the effects of a cosmological constant with its accompanying (recently discovered) accelerating cosmological metric expansion. Note that the time dependence is not conjugated out for $r < r_H$ as in the zitterbewegung so this $r \approx r_0[1 + \sinh(\omega t - .8814)]$ is a real observable motion.

Chapter 12

12.1 $r < r_H$ $\sinh \omega t$ Coordinate transformation of $Z_{\mu\nu}$: Gravity Derived

In that regard the Heisenberg equations of motion give $e^{i\omega t}$ oscillation (zitterbewegung) both for velocity and position so we use the classical harmonic oscillator probability distribution in a 2D with 1D of radial center of mass of the zitterbewegung cosine oscillation lobe. So the COM (linear) is: $x_{cm} = (\sum x m) / M = \int_0^{\pi/2} x \cos x dx / 1 = 0.5708$. As a

fraction of r_m we have $.5708 / \pi / 2 = .3634 = 1 / 2.751$ (12.1)

from the $r_M = \lambda_M / 2\pi = h / (2\pi 2m_e c)$ (i.e., half Compton λ for radius increase) and from equation 10.3 and from equation 4.15 in the vacuum $\sum H = 0$ in iHt/\hbar we have $\omega t = .8814 = 0$ so that $\omega t = .8814$ (in chapter 2 the N th (smaller) fractal scale was the source u) so that using one plausible value of the proper elapsed time is $t = 14 \times 10^9$ years, then $\omega = .8814 / t = 1.99 \times 10^{-18}$ (1/sec)

And from the third term in equation 9.2 the zitterbewegung amplitude:

$$r_M = \lambda_M / 2\pi = h / (2\pi 2m_e c) = 6.63 \times 10^{-34} / (2\pi \times 9.11 \times 10^{-31} \times 3 \times 10^8) = 1.93 \times 10^{-13} \text{ m}$$

$$r_M / 2.751 m_e c = 1.93 \times 10^{-13} / (2.751 \times 9.11 \times 10^{-31} \times 3 \times 10^8) = 2.57 \times 10^8$$

Using $S_{1/2}$ state of equation 2.6. $\epsilon = .06006 = m_\mu + m_e$

$$e^2 / 2(1+\epsilon)m_p = 9 \times 10^9 (1.6 \times 10^{-19})^2 / 2(1+\epsilon) 1.67 \times 10^{-27} = 0.06515$$

Note from equation 2.6 the $2(1+\epsilon)m_p$ in the denominator for this S state.

$$\text{Also } \cosh(0) = 1 \quad (12.2)$$

To find the physical effects of the equation 10.2 expansion *we must* do a dyadic radial coordinate transformation (equation 10.2) on this single charge horizon (given numerical value of the Hubble constant H in determining its rate) in 4.3. In doing the time derivatives we take the ω as a constant in the linear t limit:

$$\frac{\partial x^\alpha}{\partial x'^\mu} \frac{\partial x^\beta}{\partial x'^\nu} Z_{\alpha\beta} = Z'_{\mu\nu} \text{ with in particular } Z_{00} \rightarrow Z'_{00} = Z_{00} + z_{00} \quad (12.3)$$

After doing this Z'_{00} calculation the resulting (small) z_{00} is set equal to the Einstein tensor gravity source ansatz $G_{00} = 8\pi G m_e / c^2$ for this *single* charge source m_e allowing us to solve for the value of the Newtonian gravitational constant G here as well. We have then derived gravity for **all** mass since this single charge m_e source composes all mass on this deepest level as we noted in the above discussion of the equivalence principle. Note Lorentz transformation similarities in section 2.3 between $r = r_0 + \Delta r$ and $ct = ct_0 + c\Delta t$ using $D\sqrt{1-v^2/c^2} \approx D(1-\Delta)$ for $v \ll c$ with just a sign difference (in $1-\Delta$, + for time) between the time interval and displacement D interval transformations. Also the t in equation 10.2 and therefore 12.3 is for a light cone coordinate system (we are traveling near the speed of light relative to $t=0$ point of origin) so $c^2 dt^2 = dr^2$ and so equation 10.2 does double duty as a $r=ct$ time x_0 coordinate. Also note we are trying to find G_{00} (our ansatz) and we have a large Z_{00} . Also with $Z_{rr} \ll Z_{00}$ we needn't incorporate Z_{rr} . Plugging values of equation 12.1 and 12.2 and the resulting equation 10.2 into equation 12.3 we have in 4.3:

$$\frac{8\pi e^2}{m_p c^2} \delta(0) = Z_{00} = R_{00} - \frac{1}{2} g_{00} R, \quad \frac{\partial x^0}{\partial X^\alpha} \frac{\partial x^0}{\partial X^\beta} Z_{\alpha\beta} = Z'_{00} = Z_{00} + z_{00} \approx \quad (12.4)$$

$$\frac{\partial \left[x^0 + \frac{r_M}{2.751} [1 + \sinh(\omega t - .8814)] \right]}{\partial x^0} \frac{\partial \left[x^0 + \frac{r_M}{2.751} [1 + \sinh(\omega t - .8814)] \right]}{\partial x^0} Z_{00} = Z'_{00} =$$

$$\left[\frac{1}{1 + \frac{r_M \omega_{M+1}}{2.751 c} \cosh(\omega t - .8814)} \right]^2 \frac{8\pi e^2}{2(1+\varepsilon)m_p c^2} \delta(0) \equiv \left(\frac{8\pi e^2}{2(1+\varepsilon)m_p c^2} + 8\pi G \left(\frac{m_e}{c^2} \right) \right) \delta(0)$$

.so setting the perturbation z_{00} element equal to the ansatz:

$$2 \left(-\frac{e^2}{2(1+\varepsilon)m_p} \right) \left(\frac{r_M}{2.751 m_e c} \right) \omega \cosh(\omega t - .8814) = .$$

$$2(0.06515)(2.57 \times 10^8)(1.99 \times 10^{-18}) = 6.67 \times 10^{-11} \text{ Nm}^2/\text{kg}^2 \equiv G \quad (12.5)$$

from plugging in all the quantities in equation 12.4. This new z_{00} term is the classical $8\pi G\rho/c^2 = G_{00}$ source for the Einstein's equations and we have then **derived gravity** and incidentally also derived the value of the Newtonian gravitational constant since from our postulate the electron mass (our "single" postulated source) is the *only* contribution to the Z_{00} term. Note Dirac equation implies +E and -E solutions for -e and +e respectively and so in equation 12.5 we have $q_1 X q_2$ and when G is put into the Force law there is an additional $m_1 X m_2$ thus the resultant force is proportional to $q_1 X q_2 X m_1 X m_2$ which is always positive and so the gravitational force is always attractive.

To summarize we have then just done a coordinate transformation to the moving frame to find the contributing fields associated with the moving frame. Analogously one does a coordinate transformation to the charge comoving frame to show that current carrying wires have a magnetic field, also a 'new' force, around them. Also note that in the second derivative of eq.11.1 $d^2\mathbf{r}/dt^2 = \omega^2 \sinh(\omega t - .8814) =$ **radial acceleration** that in the early stages of expansion was negative (since $\omega t < .8818$), and in the later stages is positive, the contemporary "observed" time. Thus in equations 10.2 and 12.5 (originating in section 1.4 and thus in equation 1.2b) **we have a simple account of the cosmological radial acceleration expansion** (discovered recently) **so we don't need any theoretical constructs such as 'dark energy' to account for it.**

Note below equation 12.5 above that $t = 14 \times 10^9$ years and use the standard method to translate this time into a Hubble constant. Thus in the standard method this time translates into light years which are $14 \times 10^9 / 3.26 = 4.264 \times 10^9$ parsecs = 4.264×10^3 megaparsecs assuming speed c the whole time. So $3 \times 10^5 \text{ km/sec} / 4.264 \times 10^3 \text{ megaparsecs} = 70.3 \text{ km/sec/megaparsec} =$ Hubble's constant for this theory.

12.2 Cosmological Constant In This Formulation

In equation 4.6 r_H/r term is small for $r \gg r_H$ (far away from one of these particles) and so is nearly flat space since ϵ and $\Delta\epsilon$ are small and nearly constant. Thus equation 4.5 can be redone in the form of a Robertson Walker homogenous and isotropic space time. Given (from Sean Carroll) the approximation of a (homogenous and isotropic) Robertson Walker form of the metric we find that:

$$\frac{a''}{a} = -\frac{4\pi G}{3}(\rho + 3p) + \frac{\Lambda}{3}$$

Λ =cosmological constant, p =pressure, ρ =density, $a = 1/(1+z)$ where z is the red shift and 'a' the scale factor. G the Newtonian gravitational constant and a'' the second time derivative here using cdt in the derivative numerator. We take pressure= $p=0$ since there is no thermodynamic pressure on the matter in this model; the matter is commoving with the expanding inertial frame to get the a'' contribution. The usual 10 times one proton per meter cubed density contribution for ρ gives it a contribution to the cosmological constant of $4.7 \times 10^{-36}/s^2$.

Since from equation 12.1 $a = a_0 \sinh \omega t$ then $a'' = (\omega^2/c^2) \sinh \omega t = a(\Lambda/3) = (\Lambda/3) \sinh \omega t$ and there results:

$$\Lambda = 3(\omega^2/c^2)$$

From section 12.1 above then $\omega = 1.99 \times 10^{-18}$ with 1 year = 3.15576×10^7 seconds, also $c = 3 \times 10^8$ m/s. So:

$\Lambda = 3(\omega^2/c^2) = 1.32 \times 10^{-52} /m^2$, which is our calculated value of the cosmological constant. Alternatively we could use $1/s^2$ units and so multiply this result by c^2 to obtain: $1.19 \times 10^{-35}/s^2$. Add to that the above matter (i.e., ρ) contributions to get $\Lambda = 1.658 \times 10^{-35}/s^2$ contribution.

12.3 Ungauged Theory

Solve the Heisenberg equations of motion $[H, r] = i\hbar dr/dt$ using the Hamiltonian H of equation 1.9 and get $r = (r(0) + c2p/H + (\hbar c/2iH)[e^{i2Ht/\hbar} - 1](\alpha - cp/H))$. (9.2). Recall the detailed calculations in chapter 10 above. In the QM operator formulation of $R_{ij}A_j|a, t\rangle$ in section 4.3 the zitterbewegung (harmonic oscillation) coordinate system of equation 1.9 has $\square^2 g_{ii} = 0$ inside the harmonic (gauge) component (see also section 25.7) of R_{ij} . The phase in each sinusoidal zitterbewegung component is unknown so we sum over all such k components to get the zitterbewegung cloud. Thus $\sum_k R_{ij}A_j|\psi\rangle = \sum_k [R_{ij}(\text{nonharmonic terms}) + \square^2 g_{ii}]A_j|\psi\rangle = \sum_k [R_{ij}(\text{nonharmonic terms}) + 0]A_j|\psi\rangle$. So $\langle \psi | \sum_k R_{ij}A_j | \psi \rangle = \langle \psi | \sum_k R_{ij}(\text{nonharmonic terms}) + 0 | \psi \rangle$ with ungauged GR harmonic term still being zero. After taking probability density expectation values $\psi^* \psi = e^{-i\omega t} \psi^*(r, \theta, \phi) e^{i\omega t} \psi(r, \theta, \phi) = |\psi(r, \theta, \phi)|^2$ the time dependence of the zitterbewegung indeed conjugates out, becomes random for the observer, as in the standard QM but with our Maxwell equations and UNgauged GR remaining in the geometry component in front of the QM operator A in section 4.3. Alternatively the lack of phase information in the zitterbewegung sine wave coming out of equation 10.1 give this random motion cloud result. However from equation 10.2 we see that for the inside observer $r < r_H$ the $\omega \rightarrow i\omega$ and so the above complex conjugation does not occur and so there is "observed" motion (for $r < r_H$) that

does move with $1+\sinh\omega t$ in equation 10.2. Note that there are then no 'gauges' in this theory in any case.

Note in chapter 1 that equations 1.5 are used in section 17.3 to derive Maxwell equations in 4 vector potential form (A_x, A_y, A_z, V) . Thus this is the physics of potentials (A, V) that cannot be added to or subtracted from which is also consistent with this being an UNgauged field theory (can't add stuff to 'A' there either). In any case for the Aharonov-Bohm effect there is a lingering nonzero vector potential 'A' outside the B flux tube causing the neutrons moving in those opposite directions around the flux tube to interfere in an anomalous way on the other side. But in the standard theory it is only B that matters and there is very little or no B field in that region. In fact Aharonov-Bohm's successful minimal interaction idea totally neglects the residual B contribution and just adds a plus or minus KAv to the wave number in the phase thereby giving this neutron interference result. Thus you cannot add nontrivial gauge partial derivatives to the 'A' they chose since the phase difference is set by the numerical value of 'A'. Anyway, in the context of this theory at least, the Aharonov Bohm effect is further evidence of the UNgauged aspect of the physics of this new pde. Note also section 21.3 on the covariant *gaugeless* quantization of the electromagnetic field. See 25.6 and 25.9 for more *ungauged* theoretical physics discussion. In any case we have then shown that gauges are not needed anywhere in theoretical physics.

References

Merzbacher, *Quantum Mechanics*, 2nd Ed, Wiley, pp.597

Chapter 13

r<r_H Application: Rotational Selfsimilarity With pde Spin: CP violation

13.1 Fractal selfsimilar spin

The fractal selfsimilarity with the spin in the (new) Dirac equation 1.9 implies a cosmological ambient metric (Kerr metric) rotation as well as in 23.1. Thus there will be $2ds_t ds_\phi$ rotation metric cross terms with the dt (without the square) implying time T reversal nonconservation and therefore CP *non*conservation since CPT is always conserved. We thereby derive CP nonconservation from first principles. This adds another matrix element of magnitude $\sim 1/3800$ (section 16.3) for Kaon decays thus adding off diagonal elements to the CKM matrix.

Or for Kerr rotator use

$$ds^2 = \rho^2 \left(\frac{dr^2}{\Delta} + d\theta^2 \right) + (r^2 + a^2) \sin^2 \theta d\phi^2 - c^2 dt^2 + \frac{2mr}{\rho^2} (a \sin^2 \theta d\theta - c dt)^2, \quad (13.1)$$

$$\rho^2(r, \theta) \equiv r^2 + a^2 \cos^2 \theta; \quad \Delta(r) \equiv r^2 - 2mr + a^2, \text{ or}$$

$$ds^2 = dr^2 + dt^2 + 2dt dr + ..$$

The k_s is the same as the $2m$ in equation 25.25. In a polarized state ($\theta = 0^\circ, 180^\circ$) in 25.3. 25.25 the off diagonal elements are proportional to $\phi = (\phi + c)e^{-c}$. Thus if the charge e is conjugated (C , e changes sign), if dr changes sign (P , parity changes sign) and dt is reversed (t reversal) then the ds quantity on the left side of equation 1.2 is invariant. But if dr (P) changes sign by itself, or even e and P together (CP) change sign then ds is not invariant and this explains, in terms of our fractal picture, why CP and P are not conserved generally. P becomes maximally nonconserved in weak decays as we saw in above. The degree to which this nonconservation occurs depends on the “ a ” (in equation 23.1) transfer $\langle \text{final} | \text{initial} \rangle$ (equation 3.2) which itself depends on the how much momentum and energy is transferred from the S_{M+2} to the S_{M+1} fractal scales as we saw in this section. Recall chapter 5 alternative derivation of that new (dirac) equation pde (eq.1.9) **linearization of the Klein Gordon equation** ($c=1, \hbar=1, m=1$, eq.5.5):

$$\left(-\alpha_1 i \frac{\partial}{\partial x_1} - \alpha_2 i \frac{\partial}{\partial x_2} - \alpha_3 i \frac{\partial}{\partial x_3} + \beta \right) \left(-\alpha_1 i \frac{\partial}{\partial x_1} - \alpha_2 i \frac{\partial}{\partial x_2} - \alpha_3 i \frac{\partial}{\partial x_3} + \beta \right) = \quad (13.2)$$

$$-\alpha_1^2 \frac{\partial^2}{\partial x_1^2} - \alpha_2^2 \frac{\partial^2}{\partial x_2^2} - \alpha_3^2 \frac{\partial^2}{\partial x_3^2} + \beta^2 + 2 \sum_{\mu=1}^3 \sum_{\nu=1}^3 \alpha_\mu \alpha_\nu \frac{\partial}{\partial x_\mu} \frac{\partial}{\partial x_\nu} + 2 \sum_{\mu=1}^3 \beta i \alpha_\mu = -\frac{\partial^2}{\partial t^2}. \text{ This}$$

equals

$= c^2 p_1^2 + c^2 p_2^2 + c^2 p_3^2 + m^2 c^4 = E^2$ if the off diagonal elements zero (with 23.3) which is the condition used in the standard Dirac equation derivation of the α s and β . Note that

from appendix A that the off diagonal elements $2 \sum_{\mu=1}^3 \sum_{\nu=1}^3 \alpha_\mu \alpha_\nu \frac{\partial}{\partial x_\mu} \frac{\partial}{\partial x_\nu} + 2 \sum_{\mu=1}^3 \beta i \alpha_\mu$ are

equivalent to the off diagonal elements in equation 5.1 (and are corrections to 5.2 in fact) so *are not* zero for parity and CP NONconservation in this context (in a rotating

universe). So in the context of the Dirac equation the CP violation term $e_s(dr)dt \rightarrow (dr/ds)(dt/ds) \rightarrow pE\chi$ (after division by ds^2). Thus CP violation goes up as the square (pE) of the energy (so should be larger in bottom factories). The section 13.2 below Cabbibo angle calculation (not rotation related however) is an example of how this method can give the values of the other terms in the CKM matrix. They arise from calculation of $\langle Z \rangle$ between higher order m harmonics.

This section is important in that we see that CP violation is explainable and calculable in terms of perturbative effects on the ambient metric (and therefore the Dirac equation) of a rotating universe with nearly complete inertial frame dragging (section 22.1 in the E&M form), CP violation doesn't need yet more postulates as is the case with the GSW model. In fact the whole CKM matrix is explainable here as a consequence of this perturbation.

Note the orientation relative to the cosmological spin axis is important in CP violation. Integration of the data over a 3 month time (at time intervals separated by a sidereal day) is going to yield different CP violation parameters than if integration is done over a year.

13.2 Probability for $2P_{3/2}$ Giving One Decay 1S Product at $r \approx r_H$

In equation 4.12 we note that invariance over 2π rotations using $(1+2\varepsilon)d^2\theta$ does not occur anymore thus seemingly violating the conservation of angular momentum. To preserve the conservation of angular momentum the additional angle ε must then include its own angular momentum conservation law here meaning intrinsic spin $1/2$ angular momentum in the S state case and/or isospin conservation in the $2P_{3/2}$ case at $r=r_H$. In any event we must also integrate to $C=\varepsilon$.

Plug in $S_{1/2} \propto e^{i\phi/2}$, $1/2(1-\gamma^5)\psi=\chi$ into equation 3.2. In that regard note that the expectation value of γ^5 is proportional to $v \propto$ Heisenberg equation of motion derivative of $2P_{3/2} \propto e^{i(3/2)\phi}$. We get the weak, strangeness changing (second term below), decay from a $2P_{3/2} > m_p$ to the S state branch equation.1.9 expectation values from equation 3.2:

$K \int \psi/k_o c \chi dx = \langle \text{lepton} | \text{baryon} \rangle =$

$$K \left\langle e^{i\phi/2} | c k_o | 1 - (\gamma^5 i e^{i(3/2)\phi}) \right\rangle = K c k_o \left\langle e^{i\phi/2} - \gamma^5 i e^{i4/2\phi} \right\rangle = K \left(\frac{2e^{i\phi}}{i} \Big|_0^{2\pi+C} - \frac{2\gamma^5 e^{i4\phi}}{i4} \Big|_0^{2\pi+C} \right) \quad (13.2)$$

with $\langle \text{initial} | c | \text{final} \rangle^2 \approx$ transition probability as in associated production with the separate 2P proton ground state transition being the identity ($\Delta S=0$). Factoring out the 2 and then normalizing 1 to .97 simultaneously normalizes the 1/4 to .24 in equation 13.2. With this normalization we can set $\cos\theta_c=.97$ and $\sin\theta_c=.24$. Thus we can identify θ_c with the *Cabbibo angle and we have derived its value*. We can then write in the weak current sources for hadron decay the **VA structure**: $|\cos\theta_c - \gamma^5 \sin\theta_c|$. Thus with the above Cabbibo angle and this CP violation we have all the components of the CKM matrix.

Chapter 14

14.1 Integer Spin (Bosonic) Potential Energy Formulation

From the energy component of polarized representation of equation 18.1, 18.2: and using iterated (as in bosonic) section 19.13 $E^2 = p^2 c^2 + m^2 c^4 =$

$$E^2 \equiv \left(\sqrt{g_{00}} \frac{dt}{ds} \right)^2 = \frac{1}{1 - \frac{r_H}{r}} = \frac{r}{r - r_H} = \frac{r}{r - r_H} - \frac{r - r_H}{r - r_H} + 1 = \frac{r_H}{r - r_H} + 1 \quad (14.1)$$

$$= k + \frac{r_H}{r - r_H} = k + V$$

We use the equation 14.1 integer spin source and proceed in the usual way of Bjorken and Drell [13] to construct the one vertex S matrix for the triplet state (spin = $\frac{1}{2} + \frac{1}{2}$) Dirac equation but for integer spin sum. The results are those of a radially displaced Schrodinger equation Hydrogen atom. This V is an application to the chapter 6 alternative $\Sigma(1/r)^n$ geometric series expansion method if by mistake a gauge is added.

Also if we iterate equation 1.9 (square it) in the form of the separated 'r' dependence of 18.1 and 18.2 we also get 14.1 integer spin Proca Equation and the chapter 15 and chapter 16 results.

Chapter 15

$r > r_H$ Application: New Potential in New PDE Implies No Running Coupling Constant

15.1 No Need for a Running Coupling Constant

If the Coulomb $V = \alpha/r$ is used for the coupling instead of $\alpha/(k_H - r)$ (Equation 14.1) then we must multiply α in the Coulomb term by a floating constant (K) to make the coulomb V give the correct potential energy. Thus if an isolated electron source is used in Z_{00} we have that $(-K\alpha/r) = \alpha/(k_H - r)$ to define the running coupling constant multiplier "K". The distance k_H corresponds to about $d = 10^{-18} \text{m} = ke^2/m_e c^2$, with an interaction energy of approximately $hc/d = 2.48 \times 10^{-8} \text{joules} = 1.55 \text{TeV}$. For 80 GeV, $r \approx 20$ ($\approx 1.55 \text{TeV}/80 \text{Gev}$) times this distance in colliding electron beam experiments, so $(-K\alpha/r) = \alpha/(r_H - r) = \alpha/(r(1/20) - r) = -\alpha/(r(19/20)) = (20/19)\alpha/r = 1.05\alpha/r$ so $K = 1.05$ which corresponds to a $1/K\alpha \equiv 1/\alpha' \approx 130$ also found by QED (renormalization group) calculations of (Halzen, Quarks). Therefore we can dispense with the running coupling constants, higher order diagrams, the renormalization group, adding infinities to get finite quantities; all we need is the correct potential of equation 4.1.

Note that the $\alpha' = \alpha / (1 - [\alpha/3\pi(\ln\chi)])$ running coupling constant formula (Faddeev, 1981) doesn't work near the singularity (i.e., $\chi \approx e^{3\pi/\alpha}$) because the constant is assumed small over all scales (therefore there really is *no formula to compare* $\alpha/(r - r_H)$ to over all scales) but this formula works well near $\alpha \sim 1/137.036$ which is where we used it just above.

Chapter 16

New Potential And Therefore New S Matrix

16.1 $r > r_H$ Application: New Potential V in New PDE Implies New S Matrix with W and Z as Resonances

We use the equation 14.1 source and proceed in the usual way of Bjorken and Drell (here $1/r \rightarrow 1/(r-k/2)$) to construct the one vertex S matrix for the new Dirac equation 1.9. Recall the $1/2$ came from the square root in equation 14.1. Thus the k in the integrand denominator is found from the result of our $V = -1/(r-r_H/2)$ potential in equation 14.1 instead of the usual Coulomb potential $1/r$ in the large r limit.

$$S_{if} \equiv iZ \frac{1}{V} \sqrt{\frac{m^2}{E_f E_i}} \bar{u}(p_f, s_f) \gamma^0 u(p_i, s_i) \int_0^\infty \frac{e^{i(p_f - p_i)x}}{\sqrt{1 - r_H/r}} dx^4 \equiv if(u) \int_0^\infty \frac{\sqrt{r} e^{i(p_f - p_i)x}}{\sqrt{r - r_H}} dx^4 \quad (16.1)$$

rescaling $r \rightarrow r' + r_H = r$ and $t \rightarrow t' + (r_H/c) = t$ to minimize the resonance energy in $p_f - p_i$. We then obtain::

$$S_{if} \equiv iZ \frac{1}{V} \sqrt{\frac{m^2}{E_f E_i}} \bar{u}(p_f, s_f) \gamma^0 u(p_i, s_i) e^{i r_H q} \int_0^\infty \frac{\sqrt{r} e^{i(p_f - p_i)x}}{\sqrt{r - r_H}} dx^4 \equiv if(u) e^{i q r_H} \int_0^\infty \frac{\sqrt{r + r_H} e^{i(p_f - p_i)x}}{\sqrt{r}} dx^4$$

For from chapter 14 we have for integer spin $\frac{1}{\sqrt{1 - \frac{r_H}{r}}} \rightarrow 1 + \frac{1}{r - r_H}$ so:

$$S_{if} \equiv if(u) \left[\int_0^\infty \frac{e^{i(p_f - p_i)(x' - r_H/2)}}{|x'|} dx^4 - \int_0^\infty \frac{e^{i(p_f - p_i)(x' + r_H/2)}}{|x'|} dx^4 \right] \approx if(u) e^{i(r_H/2)q} \left(\int_0^\infty \frac{e^{i(p_f - p_i)(x')}}{|x'|} dx^4 - 2\pi(r_H)^3 \right) \quad (16.2)$$

So that the standard QED S matrix exists at large r. We note that $1/2 r_H$ makes the low energy section of the S matrix flatter than that of a straight sine with the middle at the below resonances.

Here $p_f - p = q$. Note in equation 16.2 the factor $i e^{i k q} = i(\cos k q + i \sin k q)$. Here we find the rotational resonances at the $2P_{3/2}$ $r = r_H$ lobes associated with maximizing the imaginary part which is $i \cos k q$ to obtain absorption scattering (at $k q = \pi$), which here will then be the masses exchanged in inverse beta decay. Also recall that $h v/c = h/\lambda = p$ and for a DeBroglie wave fundamental harmonic resonance we have $\lambda_{rot} = 2\pi r$ for a stationary particle of spin $1 = L$ (ambient E&M field source gives $L = 1$ De Broglie). For spin $1/2$ and spin 0 $\lambda_{rot} = 2\pi r$ is shortened by a factor of $2/\pi$ in the above geometry. Thus $K = (\pi/2)(1 - \epsilon)$ for spin $1/2$, $K = (\pi/2)(1 + \epsilon)$ zero charge positron electron for spin 0 (both with wavelength given by the diameter) and $K = 1$ for spin 1 and so $\lambda_{rot} \rightarrow (1/K)\lambda_{rot}$. Thus for this spin $1/2$ traveling wave in S_{ij} we use $\lambda = 2r$ diameter to be selfconsistent with the stationary spin 1 particle ($1 = 1/2 + 1/2$) zitterbewegung rotational circumferential motion (given $\lambda_{rot} = 2\pi r$) in this rotational resonance. Also from just above equation 4.11 (note $1 \pm \epsilon$ term there) we

have $Z_{00} = k = \alpha \hbar c (1 \pm \epsilon) / 2 (m_p c^2 / 2) \equiv k_s$, with the spin 1 coming out of the triplet state and the $1 \pm \epsilon$ from equation 4.12. To summarize, two spin $1/2$ particles in a traveling wave (of our S matrix) are physically identical to one spin 1 particle in circumferential motion. Thus we have two such spin $1/2$ particles spread out over a diameter $\lambda = 2r$ being equivalent to one spin 1 particle with $\lambda_{rot} = 2\pi r$ and so $\hbar / 2\pi r = p_{rot}$ which gives the $2P_{3/2}$ eigenfunction and its trifolium structure in which case we must be in resonance with an individual lobe. The wavelength of an individual $2P_{3/2}$ trifolium lobe is $\lambda_{lobe} = \lambda_{rot} / 3$ (i.e., true P wave scattering). Thus $p_{rot} = \hbar / \lambda = \hbar / \text{diameter} = \hbar / 2r = (\pi \hbar / 2\pi r) = \pi (\hbar / \lambda_{rot}) = (\pi p_{rot})$. with then $p_{lobe} = 3p_{rot}$. We therefore have for the equation 16.2 first resonance phase at

$$\phi = \pi = k_r q = k_s (p_{lobe} / \hbar) = k_s \pi (3p_{rot} / \hbar) = [(K \alpha \hbar c (1 \pm \epsilon) / (2m_p c^2))] ((3) \pi (M c^2 / \hbar c)) \quad (16.3)$$

Note in the Heisenberg equations of motion application that the phase in A_μ (required for the quantization, ch.20) in equation 17.2 goes as phase $= kx + \omega t$. In 19.34 for this large nonpionic M in 16.3 we have an imaginary exponent (since the argument of the square root is then negative) so oscillatory source F. Given $1/r \propto m$ the real part of the phase then $x \rightarrow x + x_0 \cos(K(J-1/2)/r)t = x + x_0 \cos(Km(J-1/2))t = \omega t + kx + x_0 (1 - (Km(J-1/2)t)^2 / 2) + \dots$ (neglecting the much smaller dt^3 term here) in $\square^2 A_\mu = 0$ (eq.17.2). The D'Alembertian second derivative of t^2 with respect to t is a constant so this gives rise to the Proca equation $\square^2 A_\mu + ((J-1/2)mk)^2 A_\mu = \text{massive spin 1 particle (the W)}$. See also section 25.4. If $J=1/2$ and/or $m=0$ we are back to the old Maxwell equation (17.2) spin 1 massless field. .

16.2 P wave scattering: $L=1 = 1/2 + 1/2$ Triplet and Singlet Solutions

For the q s that maximize the above $\text{Im } ie^{iks_q} = \cos kq$ term for scattering in S_{ij} . Note the \pm in equation 4.11.

Thus there is a plus and a minus solution. The minus applied to negative components. With $L=1$ (so $K=1$) charge $(+\epsilon)$: $M c^2 = K 2m_p c^2 / (3\alpha(1+\epsilon)) = \mathbf{M_W = 80.7 GeV}$. Alternatively with zero charge $(-\epsilon)$ in equation 4.11)

$$M c^2 = 2m_p c^2 / (3\alpha(1-\epsilon)) = 137.026 \cdot 938.27 \text{ MeV} / (1.5 \cdot (1-0.06)) = \mathbf{M_Z = 91.2 GeV}$$

Also $kq=0$ obtain photon mass=0. 'E' turns out to be the *mass energy of the W and Z* involved in beta decay. In this P wave scattering we noted above that the electron interacts with a $2P_{3/2}$ lobe which are confined to the region near r_H constituting approximately 1/18 of the total nucleon volume (this is out to $1/2$ dilation distance from r_H). G in the 4 body W interaction is then given by $E V = .511 ((1/18) F)^3 = .511 (.056)^3 = .9 \times 10^{-4} \text{ MeV} \cdot F^3 = G$ that must be used since the ψ s are normalized for V in the 4 body Fermi interaction integral $\int \psi_1 \psi_2 G \psi_3 \psi_4 dV$ we got from applying inverse separability on the Dirac left handed doublet (section 2.2) using equation 3.2. Thus we are merely creating a beta decay electron in a P wave scattering (lobe) in equation 3.2 as we must using the doublet of section 1.4.

In general then we have obtained an ortho triplet state here since we are merely writing the Clebsch Gordon coefficients for this addition of two spin $1/2$ angular momentums:

$$|1/2, 1/2, 0, -1\rangle, \dots |1/2, 1/2, 0, 0\rangle, \dots |1/2, 1/2, 0, 1\rangle, \dots \text{ or } +W Z_0, -W$$

Thus we have the **ortho** states over Z_0 , W^- and W^+ and the singlet **para** state at $L=0$, 129 GeV (given by $1-2\epsilon$ multiplication of the charged s $L=1/2$ state at $\sim 142 \text{ GeV}$ in equation eq. 16.3). If the scattering is $2P_{3/2}$ instead $m_p \rightarrow m_e$ and so $142 \text{ GeV} \rightarrow 135 \text{ MeV}$, the uncharged spin 0 pion and the same ortho-para organization applies but at these far lower

energies. Note the common denominator between the Fermi interaction term (the one with the G) and the W, Z mass scattering is the electron scattering in a 2P_{3/2} lobe: they are the same interaction analyzed two different ways. So we say that the W or Z were ‘exchanged’ in the Fermi interaction or that a left handed electron and left handed neutrino took part in that G, 4 body interaction, same thing.

The beautiful thing to be noted here is that for the doublet resonance with the 2P_{3/2} lobe at $r=r_H$ that minimizes energy you get the spin 1 W and Z and the value of the Fermi G! We have also shown that this doublet interaction corresponds to the exchange of massive spin 1 particles (recall spin 1/2 s forbidden by that j-1/2 factor).

16.3 Mass Hierarchy Problem

Note the $\sqrt{g_{tr}}$ approximation in equation 19.3 goes as $1+r/4-r^2/10+$. Thus the $r^2/10$ becomes comparable to $r/4$ when $r \approx 2$ which then corresponds to the maximum strangeness states (section 19.9) of maximum $\sim 2\text{GeV}$ (2P state $n=3$) where the $r/4$ perturbation was used. Also the S state shares an electron giving a sp^2 hybridization, thus providing a charge (nonzero Z_{0oS}) even if the state is filled. We must then find the r^2 (factorized) term that gives the $1+r/4+r^2/10$ polynomial (to move up further in energy to $n=4,5,6$) thus resulting in a $dE/r' = -kr'$ Hooke’s law $E = \frac{1}{2}kr'^2$ potential energy term with S state hybridization giving a bare S state Z_{0oS} for scattering. Thus, given these proton nonrelativistic 2P_{3/2} state objects, we can construct a (crude approximation) harmonic oscillator Schrodinger equation with here then the 5nd harmonic of twice the frequency (energy) of the next lower 2P 4th harmonic state. $\hbar\omega_0(n+1/2) = 4m_p c^2(n+1/2)$, where this $\hbar\omega_0$ is from section 16.1. Finally the last 2P state fills at the 6th particle and so no free electrons are available to scatter, they are all paired analogous to the inert noble gas states. Set $L = 1/2$ in equation 16.4 to get this mass. Thus the S state of equation 16.3 is exposed to scattering at the very high energy of 143Gev.

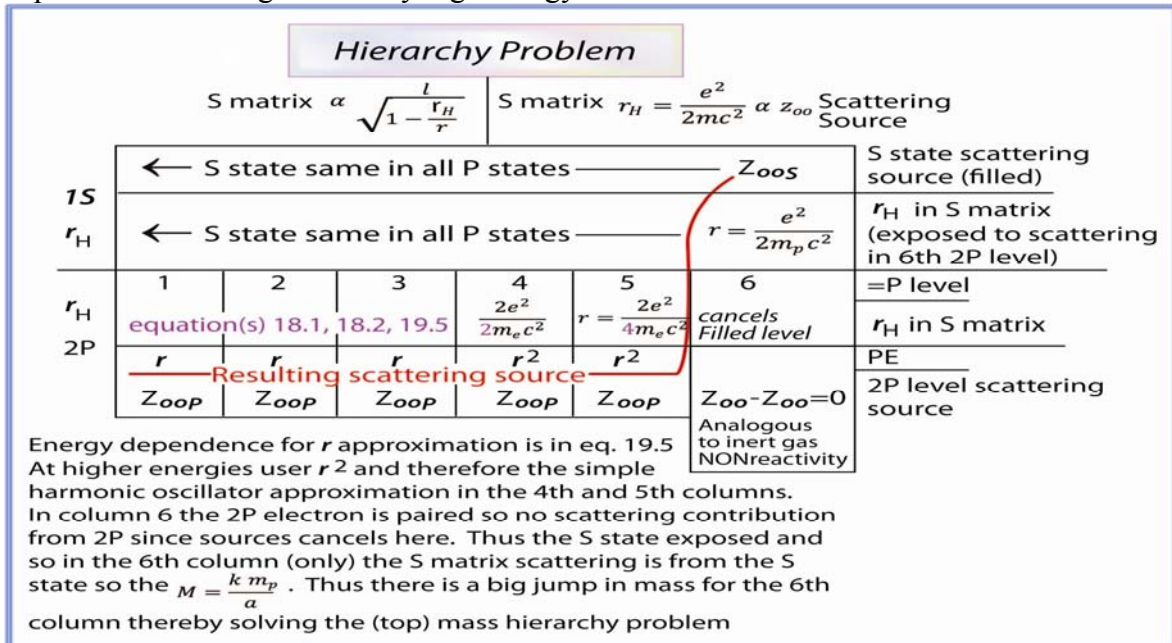


Figure 16-1 The Solution To the Mass Hierarchy Problem For the Top

Thus to chapter 3 we can add the solution to the mass hierarchy problem for the 6 th 2P state(top)

16.4 Electron E&M Component Of $\sum_M H_M=0=H_N +H_{N+1}$

Section 1.5 and in equations 11.2, 4.15 H_N is the E&M component of $\sum_M H_M=0=H_N +H_{N+1}$ and H_{N+1} is the neutrino component (recall $H_{N+1} \propto \sigma \bullet p$). Thus for the vacuum $\langle H_N+H_{N+1} \rangle = \langle \sum \hbar\omega + H_v \rangle \equiv \langle \sum \hbar\omega + \Delta\varepsilon_v \rangle$. Therefore if a set of nearby grounded conducting plates ($V=0$) were set up some of the harmonics in $\sum \hbar\omega$ (given by the distance between the plates and $V=0$ at the surfaces) between the plates would not exist since they don't satisfy the boundary conditions therefore there is less E&M $\hbar\omega$ energy density between the plates than on the outside. But the H_v neutrinos would pass right through the plates therefore there would be a higher density $\hbar\omega$ (E) outside the plates than inside and this condition would be unaffected here by the H_v neutrinos. Thus an inward force on the plates can be calculated. These excluded harmonics and the resulting inward force (dE/dx) are also the assumptions used to calculate the Casimir effect (Casimir, 1948). Note that the Casimir effect would exist in this situation even though in equation 4.15. $\sum_M H_M=0$ exactly *outside* this excluded region (far from the region between the plates) in the vacuum. Thus the vacuum has a net *zero* energy [recall $\sum_M H_M=0$]. This means that there is *no net zero point energy* in free space. There are many other reasons for assuming zero net energy for the vacuum (such as the fact that space is essentially 'flat', with the "flatness" problem solved trivially here by $\sum_M H_M=0$) so this reconciliation with the Casimir effect is important. Using the above electron $\Delta\varepsilon$ and equation 4.15 for the vacuum again:

$$\sum H_N = H_M + H_{M+1} = 0 = \Delta\varepsilon - \Delta\varepsilon_v = 0.$$

If a local perturbation, such as the sun's gravitational field is added into the electron Hamiltonian of equation 1.9 the perturbation effect on the electron mass is negligible as you could show. But since equation 4.2 (in section 4.5) is not zero anymore because the assumption of nonhomogenous and isotropic space time is no longer valid (in evaluating $R_{22} - 1/2 g_{22} R$) so ambient S_M metric has much larger contributions than the S_{M+1} ambient metric. In other words equation 4.10 must be modified to $(H_M + H_{M+1})\psi = \Delta G \psi$ so $(H_M + H'_{M+1}) = 0$ where $H'_{M+1} \equiv H_{M+1} - \Delta G$. Therefore the perturbation effect on the electron neutrino mass would be relatively large due to the small (S_{M+1}) ambient metric gravity Hamiltonian contribution H_{M+1} in equation 11.2, with again E&M field contributions in H_M . Thus in a gravitational perturbation the neutrino mass energy varies with gravitational field, with the electron neutrino mass at the sun's surface then being 28X larger than at the earth's surface due to its 28 X larger surface gravitational acceleration. Therefore near a horizon k a locally transmitted neutrino energy would be nearly all potential (energy $PE_v \propto 1/r$ or more precisely $k\Delta\varepsilon/r$ since for $r \approx k$ then $\Delta\varepsilon - \Delta\varepsilon_v = 0$ since again:

$H_M + H_{M+1} = 0 = \Delta\varepsilon - \Delta\varepsilon_v = 0$ and $|PE_v| \gg KE_v$ near the horizon. In other words $\Delta\varepsilon - (KE_v + m_v c^2) / \sqrt{g_{00}}$ (from $(dt/ds) / \sqrt{g_{00}} = E_v$) and $KE_v \gg m_v c^2$ in free space (but $|PE_v| \gg KE_v$ near the horizon). Thus in general we have: $m_e c^2 - E_v = m_e c^2 -$

$$E_{\nu_0}/\sqrt{g_{00}} = m_e c^2 - \left(\frac{E_N + m_\nu c^2}{\sqrt{1 - 2GM/rc^2}} \right) = 0 \text{ where } E_N \text{ is normalized to } KE_\nu = \text{kinetic}$$

energy = $E_N = m_e c^2$ since m_ν is so small and $r \gg 2GM/c^2$ (to cancel the energies in free space). This equals $m_e c^2 - m_e c^2(1 + GM/rc^2) - m_\nu c^2(1 + GM/rc^2) \approx \Delta\varepsilon - (PE + KE)_\nu = 0$. Further away from the horizon we have in general $\Delta\varepsilon - (PE + KE)_\nu = 0$. Therefore to calculate the neutrino *rest* mass we adopt the *rest* frame (comoving frame) in which $KE_\nu = 0$ and therefore that $H_N + H_{N+1} = \Delta\varepsilon - \Delta\varepsilon_\nu = 0 = \Delta\varepsilon - (0 + k\Delta\varepsilon/r) = 0$. For the sun $k = 2.95 \text{ km}$ and at the surface ($r = 6.96 \times 10^8 \text{ m}$) thusly $\Delta\varepsilon_\nu = k\Delta\varepsilon/r = (2,950/6.96 \times 10^8 \text{ m}) 511,000 \text{ eV} = 2.166 \text{ eV} = \Delta\varepsilon_\nu$ and at the earth's surface $\Delta\varepsilon_\nu = 2.166/28 = .077 \text{ eV}$. *Therefore the electron neutrino mass at the earth's surface is smaller than at the sun's surface and is .077 eV.* More measurements at Kamiokande might verify this result. 2.2 eV is the tritium decay upper limit of the electron neutrino with 1 eV being the average from Sudbury 2006. We found that limit to be 2.166 eV above. The acceleration gets large near $r \sim 2GM/c^2$ horizon and is the same for both the electron and neutrino implying the same force as we well as the same mass and thus implying identical particles. Thus *the neutrino is identical to an electron exactly at the horizon since the metric force at the horizon is the same as that on electron as well and so the neutrino with mass does not violate our postulate that there is a ONE nonzero rest mass particle* as in equation 4.14. But at the earth's surface the above ratio of the neutrino mass to the electron mass is $.077/511,000 = 1/6.64 \times 10^6 = 1/k = \text{ratio of charges}$. Add to equation 14.1 the interaction of this (weak) charge with the electric charge for the 2P state at r_H nucleon decay and get from 14.1

$$E = \left(1 / \sqrt{1 - \frac{2e^2}{r2m_p c^2} + \frac{2(e/k)e}{rm_e c^2} + \dots} \right) (2m_p c^2) \quad (16.3)$$

References

Bjorken and Drell, *Relativistic Quantum Mechanics*
 H.B.G. Casimir, Phys. Rev. 73(1948)

Chapter 17

17.1 Implications of $g_{00} = 1 - 2e^2/rm_e c^2 = 1 - eA_0/mc^2 v^0$ For Small Relative v Low Temperature Limit

Next we use above equation 4.1 $g_{00} = 1 - 2e^2/rm_e c^2 = 1 - eA_0/mc^2 v^0$ in the geodesics

$$\frac{d^2 x^\mu}{ds^2} = -\Gamma_{\nu\lambda}^\mu \frac{dx^\nu}{ds} \frac{dx^\lambda}{ds} \quad (17.1)$$

Or in general
$$g_{ii} \equiv \eta_{ii} + h_{ii} = 1 - \frac{eA_i(x,t)}{m_\tau c^2 v^i}, i \neq 0, \quad (17.2)$$

$A'_0 \equiv e\phi/m_\tau c^2$, $g_{00} \equiv 1 - \frac{e\phi(x,t)}{m_\tau c^2} = 1 - A'_0$, and define $g'_{\alpha\alpha} \equiv 1 - A'_\alpha/v_\alpha$, ($\alpha \neq 0$) and $g''_{\alpha\alpha} \equiv g'_{\alpha\alpha}/2$ for large and near constant v. In the weak field $g^{ii} \approx 1$. Also use the total differential $\frac{\partial g_{11}}{\partial x^\alpha} dx^\alpha = dg_{11}$ so that using the chain rule gives us

$$\frac{\partial g_{11}}{\partial x^\alpha} \frac{dx^\alpha}{dx^0} = \frac{\partial g_{11}}{\partial x^\alpha} v^\alpha = \frac{dg_{11}}{dx^0} \approx \frac{\partial g_{11}}{\partial x^0}.$$

gives a new $A(1/v^2)dv/dt$ force term added to the first order Lorentz force result in these geodesic equations (Sokolnikoff, pp.304):

$$\begin{aligned} -\frac{d^2 x^1}{ds^2} &= \Gamma_{11}^1 v_1 v_1 + \Gamma_{12}^1 v_1 v_2 + \Gamma_{13}^1 v_1 v_3 + \Gamma_{10}^1 v_0 v_1 + \Gamma_{21}^1 v_2 v_1 + \Gamma_{22}^1 v_2 v_2 + \Gamma_{23}^1 v_2 v_3 + \Gamma_{20}^1 v_2 v_0 + \\ &\Gamma_{31}^1 v_3 v_1 + \Gamma_{32}^1 v_3 v_2 + \Gamma_{33}^1 v_3 v_3 + \Gamma_{30}^1 v_3 v_0 + \Gamma_{01}^1 v_0 v_1 + \Gamma_{02}^1 v_0 v_2 + \Gamma_{03}^1 v_0 v_3 + \Gamma_{00}^1 v_0 v_0 = \\ &\frac{g^{11}}{2} \left(\frac{\partial g'_{11}}{\partial x^1} \right) v_1 + \frac{g^{11}}{2} \left(\frac{\partial g'_{11}}{\partial x^2} \right) v_2 + \frac{g^{11}}{2} \left(\frac{\partial g'_{11}}{\partial x^3} \right) v_3 + \frac{g^{11}}{2} \left(\frac{\partial g'_{11}}{\partial x^0} \right) v_0 + \\ &\frac{g^{11}}{2} \left(\frac{\partial g'_{11}}{\partial x^2} \right) v_2 - \frac{g^{11}}{2} \left(\frac{\partial g'_{22}}{\partial x^1} \right) v_2 + 0 + 0 + \frac{g^{11}}{2} \left(\frac{\partial g'_{11}}{\partial x^3} \right) v_3 + 0 - \frac{g^{11}}{2} \left(\frac{\partial g'_{33}}{\partial x^1} \right) v_3 + 0 + \\ &\frac{g^{11}}{2} \left(\frac{\partial g'_{11}}{\partial x^0} \right) v_0 + 0 + 0 - \frac{g^{11}}{2} \left(\frac{\partial g'_{00}}{\partial x^1} \right) v_0 + O\left(\frac{A_i dv}{v^2 dt}\right) = v_2 \left(\frac{\partial g''_{11}}{\partial x^2} - \frac{\partial g''_{22}}{\partial x^1} \right) + v_3 \left(\frac{\partial g''_{11}}{\partial x^3} - \frac{\partial g''_{33}}{\partial x^1} \right) + \\ &\left(\frac{\partial g''_{11}}{\partial x^\alpha} v_\alpha + \frac{\partial g''_{11}}{\partial x^0} v_0 \right) - \left(\frac{\partial g''_{00}}{\partial x^1} \right) v_0 + O\left(\frac{A_i dv}{v^2 dx}\right) \approx - \left(\frac{\partial g''_{00}}{\partial x^1} \right) v_0 + v_2 \left(\frac{\partial g''_{11}}{\partial x^2} - \frac{\partial g''_{22}}{\partial x^1} \right) + \\ &v_3 \left(\frac{\partial g''_{11}}{\partial x^3} - \frac{\partial g''_{33}}{\partial x^1} \right) + O\left(\frac{A_i dv}{v^2 dt}\right) \approx \frac{e}{m_\tau c^2} \left(-\vec{\nabla} \phi + \vec{v} X (\vec{\nabla} X \vec{A}) \right)_x + O\left(\frac{A_i dv}{v^2 dt}\right) \end{aligned}$$

Thus we have the Lorentz force equation form $\left(-\left(\frac{e}{m_\tau c^2}\right) \left(\vec{\nabla} \phi + \vec{v} X (\vec{\nabla} X \vec{A}) \right) \right)_x$ plus the derivatives of $1/v$ which

are of the form: $\mathbf{A}_i(d\mathbf{v}/dt)/v^2$. **This new term $A(1/v^2)dv/dt$ is the pairing interaction.** (17.3)

This approximation holds well for nonrelativistic and nearly constant velocities and low B fields but fails at extremely low velocities so it works when $v \gg (dv/dA)A$. This constraint also applies to this ansatz if it is put into our Maxwell equations in the next section. Given a

stiff crystal lattice structure (so dv/dt is large also implying that lattice harmonic oscillation isotope effect in which the period varies with the (isotopic) mass.) this makes the pairing interaction force $A_i(dv/dt)/v^2$. The relative velocity “ v ” will then be small in the denominator in some of the above perturbative spatial derivatives of the metric $g_{\alpha\alpha}$ (e.g., the $1/v$ derivative of $H_2 (A/v^2)(dv/dt)$). This fact is highly suggestive for the velocity component “ v ” because it implies that at cryogenic temperatures (extremely low relative velocities in normal mode antisymmetric motion) new forces (pairing interactions?) arise from the above general relativity and its spin 0 (BCS) and spin 2 states¹ (D states for CuO_4 structure). For example the mass of 4 oxygens ($4 \times 16 = 64$) is nearly the same as the mass of a Cu (64) so that the SHM dynamics symmetric mode (at the same or commensurate frequencies) would allow the conduction electrons to oscillate in neighboring lattices at a relative velocity of near zero (e.g., $v \approx 0$ in $(A/v^2)(dv/dt)$ making a large contribution to the force), thus creating a large BCS (or D state) type pairing interaction using the above mechanism. Note from the dv/dt there must be accelerated motion (here centripetal acceleration in BCS or linear SHM as in the D states) as in pair rotation but it must be of very high frequency for dv/dt (lattice vibration) to be large in the numerator also so that v , the velocity, remain small in the denominator with the phase of “ A ” such that $A(dv/dt)$ remain the same sign so the polarity giving the A is changing rapidly as well. This explains the requirement of the high frequency lattice vibrations (and also the sensitivity to valence values giving the polarity) in creating that pairing interaction force. Other attempts have been made to put in a nonlinear addition to the Maxwell equations (see section 17.3) that only is observable at low temperatures. For example there are the nonlinear Schrodinger equation theories. But in these theories these nonlinear terms were postulated (Chern Simons), not derived as they are here.

Thus we finally understand the pairing interaction in SCs, still the deepest mystery in standard SC theoretical speculations.

Most superconductivity (other than the cuprates) is for single mass m but with alternating spin states between lattice points. Thus Niobium with possible lattice of tin or Germanium (Group 16 of the periodic table) or the nictides with iron in a lattice of arsenic or antimony (group 15). Thus here v is small because the atomic weights are the same. The correct lanthanum, fluorine, doping must be applied to give the morphology required.

Also note the v in the denominator of 17.1 and 17.2 and so possible singular behavior there at the lowest temperatures. Thus at extremely low temperatures there exists some of the same type of exotic singularity behavior (e.g., chapter 19) as exhibited near r_H and so seen in high energy scattering. You are studying some of the properties of the high energy extreme then at the very lowest energy extreme!

17.3 Weak Field Approximation

We adopt an ansatz (equation 10.1) $g_{oo} = e^{iA_o}$ (or the Fourier series sum). Also recall from equation 1.7 that $g_{oo}g_{rr}$ is a constant. So we can do the equation 1.3 variation with respect to time (and not r) just as in classical Hamiltonian theory. So from equation 3.5 we have $\delta(g_{oo}g_{rr}) = 0 = (\partial(g_{oo}g_{rr})/\partial t)dt = ((\partial g_{oo}/\partial t)g_{rr})dt + ((\partial g_{rr}/\partial t)g_{oo})dt = ((g_{oo}\partial iA_o/\partial t)g_{rr})dt + (g_{rr}\partial iA_r/\partial r)(dr/dt)g_{oo}dt = (g_{oo}(\partial iA_o/\partial t)g_{rr})dt + (g_{rr}(\partial iA_r/\partial r)cg_{oo})dt$. Multiplying both sides by $1/cg_{oo}g_{rr}idt$ gives us:

$$\partial A_o/c\partial t + \partial A_r/\partial r = 0 \quad (17.1)$$

But from the Heisenberg equations of motion equation 10.1 $e^{iu} = e^{i\omega t} \rightarrow e^{i(\omega t - kr)}$ for observer motion v (with $kr \propto vt$) and so each g_{00} and g_{rr} obeys a wave equation in t $\partial^2 g_{ii}/c^2 \partial t^2 - \partial^2 g_{ii}/\partial r^2 = 0$. For example in free space: $g_{00} \partial^2 A_0/c^2 \partial t^2 - g_{00} \partial^2 A_0/\partial r^2 = 0$ and $g_{rr} \partial^2 A_r/c^2 \partial t^2 - g_{rr} \partial^2 A_r/\partial r^2 = 0$ or

$$\partial^2 A_0/c^2 \partial t^2 - \partial^2 A_0/\partial r^2 = 0 \quad \text{and} \quad \partial^2 A_r/c^2 \partial t^2 - \partial^2 A_r/\partial r^2 = 0 \quad (17.2)$$

So we have the (Lorentz gauge, (Pugh, pp.270)) weak field formulation of the free space E&M Maxwell equations in 2D and can generalize to 4D as before.

References

Pugh, Pugh, *Principles of Electricity and Magnetism*, 2nd Ed. Addison Wesley, pp.270

Bjorken and Drell, *Relativistic Quantum Mechanics*, pp.60

Sokolnikoff, *Tensor Analysis*, pp.304

17.4 Recall from section 2.2 The S State Fitzgerald Contraction

S state and equation 1.9: r_H is Lorentz contracted by just NONrotating, straight line $\sqrt{(1-v^2/c^2)}$ from the Lorentz transformation and is again given by $r \rightarrow r_H$ using equation 2.5 multiplication with $1+\epsilon$ (i.e., gives that above $N=2 m_p$ eigenvalue). So for S states and the multiple charges of the tauon, muon and electron:

$$r_S = (2e^2)/(2(1+\epsilon+\Delta\epsilon) m_p c^2) \quad (2.4)$$

Also recall the 2,0,0 state hydrogen eigenfunction

$$\psi_{2,0,0} = \frac{2}{(2a_0)^{3/2}} (1 - r/2a_0) e^{-r/2a_0}$$

Next find center of charge value of $r = \langle r \rangle =$

$$\begin{aligned} \langle r \rangle &= \int_0^\infty r \psi_{2,0,0}^* \psi_{2,0,0} r^2 dr = \frac{4}{(2a_0)^3} \int_0^\infty r (1 - r/2a_0) e^{-r/2a_0} (1 - r/2a_0) e^{-r/2a_0} r^2 dr = \\ &= \frac{4(2a_0)^4}{(2a_0)^3} \int_0^\infty r^3 (1 - 2r/2a_0 + r^2/4a_0^2) e^{-r/2a_0} dr = \frac{4(2a_0)^4}{2(2a_0)^3} \int_0^\infty \left(\frac{r^3}{8} - 2\frac{r^4}{16} + \frac{r^5}{32} \right) e^{-r/2a_0} dr = \frac{8a_0}{2} \left(\frac{6}{8} - \frac{48}{16} + \frac{120}{32} \right) = \end{aligned}$$

$6a_0$ to find the position of an electron that gives the same QM state energy if put in the classical potential energy analogous to how the role the center of mass plays in torque theory. This is a different than the position at the peak of $\psi^* \psi$, which is the Bohr radius for the 1,0,0 state. For this 2,0,0 state however the radius of the Bohr orbit is $4a_0$. $6a_0$ is measured from the Compton wavelength λ_c so $6a_0 \rightarrow 6a_0 + \lambda_c$. Recall from the position that it holds in the standard Dirac equation that the energy (in equation 2.6) goes as:

$$\frac{1}{\sqrt{\kappa_{oo}}} = \frac{1}{\sqrt{1 - \frac{r_H}{r}}} = 1 + \frac{r_H}{2r} - \frac{3}{8} \left(\frac{r_H}{r} \right)^2 - \dots \text{evaluate at } \langle r \rangle \text{ for 2 charges } (2e) \text{ to reduce to the}$$

ordinary large central mass Coulomb potential in the first order term. Again $1+\epsilon = 1.06$. From section 1 of chapter 2 set total energy equal to $[2(1+\epsilon+\Delta\epsilon)m_p c^2]$ and set k to get the

Coulomb potential energy for electron moving around heavy central mass

$$2((1 + \varepsilon + \Delta\varepsilon)m_p c^2) / \sqrt{\kappa_{oo}} = \frac{2(1.06)m_p c^2}{\sqrt{1 - \frac{r_H}{6a_o}}} = 2(1.06)m_p c^2 + \frac{4e2e}{2(6a_o)} 2(1.06)m_p c^2$$

$$E_i = -\frac{3}{8} \left(\frac{\frac{4e2e}{k2(1.06)m_p c^2}}{6a_o} \right)^2 2(1.06)m_p c^2 + \dots$$

Note that a 2 gets canceled in the Coulomb term. We are left with $4e^2/r$ in that term and the mass of a tauon, muon, and two electrons plus the 3/8 term. Next subtract off the tauon, muon, masses and their respective $e^2/r + e^2/r = 2e^2/r$ potential energies leaving two electron masses and their $2e^2/r$ contribution. Then divide the 3/8 contribution and what was leftover in the other terms by two to get the single electron mass and standard Coulomb potential contribution. The second term is the ordinary Coulomb potential at $r=6a_o$ average electron distance for $\psi_{2,0,0}$ given by $e^2/r = e^2/6a_o$ for a electron moving around heavy central mass. Thus to get the hydrogen atom 2S state electron contribution coming out of this we apply this energy relation to just the electron mass energy component in the first term. The last term then equals:

$$\left(\frac{1}{2} \right) \frac{3}{8} \left(\frac{8(8.98 \times 10^9)(1.602 \times 10^{-19})^2}{6((53 \times 10^{-10} + 2.42 \times 10^{-12})2(1.06))(1.673 \times 10^{-27}(3 \times 10^8)^2)} \right)^2 2(1.06)(1.67 \times 10^{-27})(3 \times 10^8)^2$$

$$= 1.83 \times 10^{-26} =$$

$$= hf = 6.626 \times 10^{-34} 27,400,000 \text{ so that } f=27\text{Mhz}$$

Recall also the 1000Mhz component is due to the electron zitterbewegung cloud itself taking up space which we get by adding the Compton wavelength directly into the Coulomb potential radius at $6a_o$.

Thus we account for the entire Lamb shift without evaluating any higher order diagrams as mentioned earlier in section 3.3. We don't need renormalization anymore. We also do not require renormalization and higher order diagrams for the derivation of anomalous gyromagnetic ratio as we see in the next chapter (ch.18). Even though, in principle, we do not need it in general it is better to use it since the fine structure constant α is all that is needed in that multidiagram method and α is accurately known. Recall we derived that old method in chapter 6.

17.5 S Matrix at Energies $\gg 1\text{TEV}$

As in equation 14.1 even in the ultrarelativistic limit: total energy equals

$$\frac{dt}{ds} \sqrt{g_{oo}} \rightarrow \frac{1}{\sqrt{1 - \frac{\alpha}{r}}}. \text{ But the mass term must be subtracted to get the postential. With}$$

electron rest mass neglected. In this limit V is large and can be replaced with $\frac{dt}{ds} \sqrt{g_{oo}} - 1$

$=V$ so in the one vertex S matrix we have then to evaluate the integral of $\frac{dt}{ds} \sqrt{g_{oo}} - 1$

instead of the integral of the usual coulomb V as in standard QED.. In evaluating this integral we use the integrals in spherical coordinates $dVol = r^2 \sin \theta dr d\theta d\phi$: So in the one vertex S matrix the integral coefficient goes as:

$$\int V dVol = \int \left(\frac{1}{\sqrt{1 - \frac{\alpha}{x}}} - 1 \right) x^2 dx = \int \left(\frac{x^2}{\sqrt{1 - \frac{\alpha}{x}}} - x^2 \right) dx = \int \left(\frac{x^{5/2}}{\sqrt{x - \alpha}} - x^2 \right) dx$$

So we have (recall $\int x^2 dx = x^3/3$) we

$$\int \left[\frac{x^{5/2}}{\sqrt{x - \alpha}} - x^2 \right] dx = \frac{x^{5/2} \sqrt{x - \alpha}}{3} + \frac{5\alpha}{6} \int \frac{x^{3/2}}{\sqrt{x - \alpha}} dx - \frac{x^3}{3} =$$

$$\frac{x^{5/2} \sqrt{x - \alpha}}{3} - \frac{x^3}{3} + \frac{5\alpha}{6} \left(\frac{2x^{3/2} \sqrt{x - \alpha}}{4} + \frac{3}{4} \frac{2\alpha}{\sqrt{-1}} \ln(\sqrt{x} + \sqrt{x - \alpha}) \right)$$

In the limit of high energy and closest approach impact parameter (so $x \rightarrow \alpha$) this equals

$$-\frac{\alpha^3}{3} + \left(\frac{3}{4} \frac{2\alpha}{\sqrt{-1}} \ln(\sqrt{\alpha}) \right)$$

To proceed one step further note $\frac{dt}{ds} \sqrt{g_{oo}} \rightarrow \frac{1}{\sqrt{1 - \frac{\alpha}{r}}} \rightarrow \frac{1}{\sqrt{1 - \frac{\alpha}{r} - \epsilon'}}$ with $\epsilon' = (4/3)\epsilon$.

and the pion mass calculation $=\epsilon'$ resulting in equation 19.22 center of mass energy 1.08 GeV

and so the Δ resonance pion photoproduction put in:

$$\text{(Pion photoproduction Energy)}/\text{CBR photon energy} = 2.23 \times 10^{-11} / 2 \times 10^{-22} =$$

$$\text{GZK}/\text{(Proton rest mass Energy)} = \text{GZK} / 1.5 \times 10^{-10}.$$

So $\text{GZK} = 17 = 1 \times 10^{20} \text{eV}$. Need to include pion photoproduction of the Δ energy in this S matrix for completeness.

17.6 Separation of Variables On Equation 1.9

Here we use separation of variables on

$$\Sigma_{\mu}(\sqrt{\kappa_{\mu\mu}}\gamma^{\mu}\partial\psi/\partial x_{\mu})-\omega\psi=0, \text{ with } \kappa_{00}=1/\kappa_{rr}=1-r_H/r \quad 1.9)$$

Note from chapter 5 we can begin the derivation of equation 1.9 with:

$$ds^2 = \kappa_{rr}dr^2 + r^2d^2\theta + r^2\sin^2\theta d^2\phi - \kappa_{00}c^2 dt^2$$

$\kappa_{00} = 1-r_H/r$ and $\kappa_{rr} = 1/\kappa_{rr}$. Note in the linearization no kappas will be in front of the angle terms. Thus we can show that we can separate out (i.e., use separation of variables) the angle terms and that the ψ will then be of the form $\psi=F(r,t)\Theta(\theta)\Phi(\phi)$. Thus our $2P_{3/2}$ state at r_H result is indeed a solution to equation 1.9. We also can write

$$ds^2 = \kappa_{rr}dr^2 + r^2d^2\theta + r^2\sin^2\theta d^2\phi - \kappa_{00}c^2 dt^2 = dr'^2 + r^2d^2\theta + r^2\sin^2\theta d^2\phi - c^2 dt'^2$$

When the quantum operator condition is used $p_r \rightarrow \sqrt{\kappa_{rr}} p_r = \sqrt{\kappa_{rr}} \partial/\partial r \equiv \partial/\partial r'$. Thus using r' and t' we can do the standard separation of variables for the standard Dirac equation. At the end the $\sqrt{\kappa_{rr}}$ s are put back in and we get for the radial component of the Dirac equation:

$$\left[\left(\frac{dt}{ds} \sqrt{g_{00}} m_p \right) + m_p \right] F - \hbar c \left(\sqrt{g_{rr}} \frac{d}{dr} + \frac{j+3/2}{r'} \right) f = 0$$

$$\left[\left(\frac{dt}{ds} \sqrt{g_{00}} m_p \right) - m_p \right] f + \hbar c \left(\sqrt{g_{rr}} \frac{d}{dr} - \frac{j-1/2}{r'} \right) F = 0$$

We next take into account that r' denominator in the centripetal term. We do this by finding the gyromagnetic ratios g_y for the spin polarized $F=0$ case. Recall the usual calculation of rate of the change of spin S gives $dS/dt \propto m \propto g_y J$ from the Heisenberg equations of motion. We note that $1/\sqrt{g_{rr}}$ rescales dr in $\left(\sqrt{g_{rr}} \frac{d}{dr} + \frac{J+3/2}{r} \right) f$. Thus to

have the same rescaling of r in the second term we must multiply the second term denominator (i.e., r) and numerator (i.e., $J+3/2$) each by $1/\sqrt{g_{rr}}$ and set the numerator equal to $3/2+J(g_y)$, where g_y is now the gyromagnetic ratio. This makes our equation compatible with the standard Dirac equation. The effect of this r' term then is to alter the gyromagnetic ratio from the nominal value of 2. This method is used in the next chapter.

Note the direct derivation of the low temperature results (section 17.1) and high energy results summarized in chapter 3. See figure 17.1 for how this contrasts with the mainstream approach.

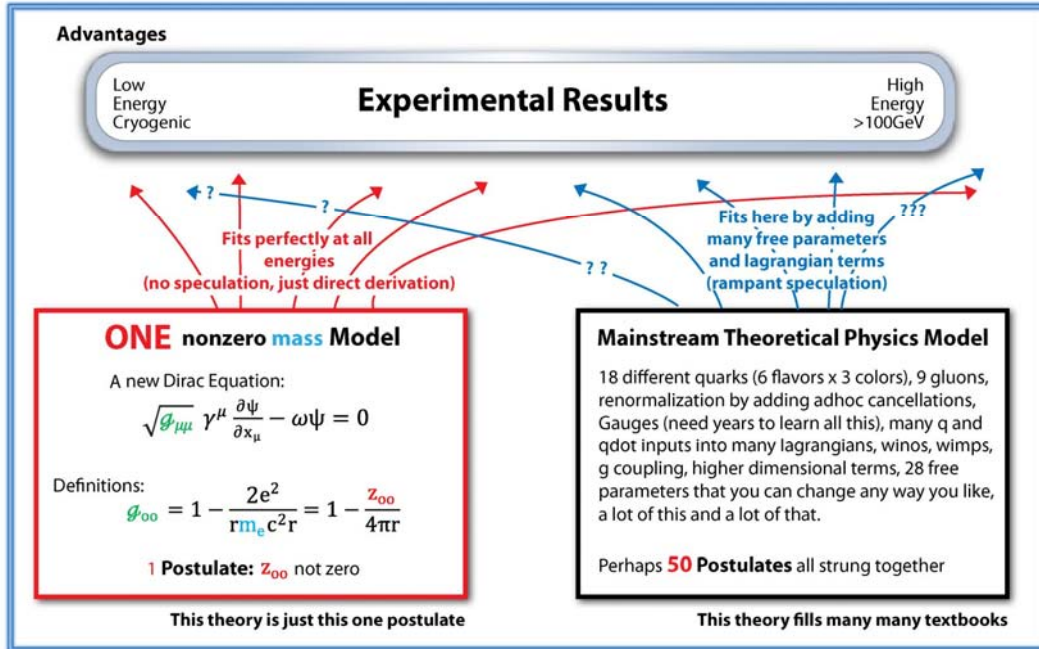


Figure 17-1 Comparison of new pde model with standard model and strings
17.7 Summary of Derivations From the New Pde

Even though there are no gauges in this theory the consequences of the SM standard gauges (such as SU(3)) can be derived from this theory as we have seen. In figure 17.1 below is a summary of these derivations. For example SU(3) comes from that new pde (equation 1.9) $2P_{3/2}$ at $r \approx r_H$. as we see in figure 17-2.

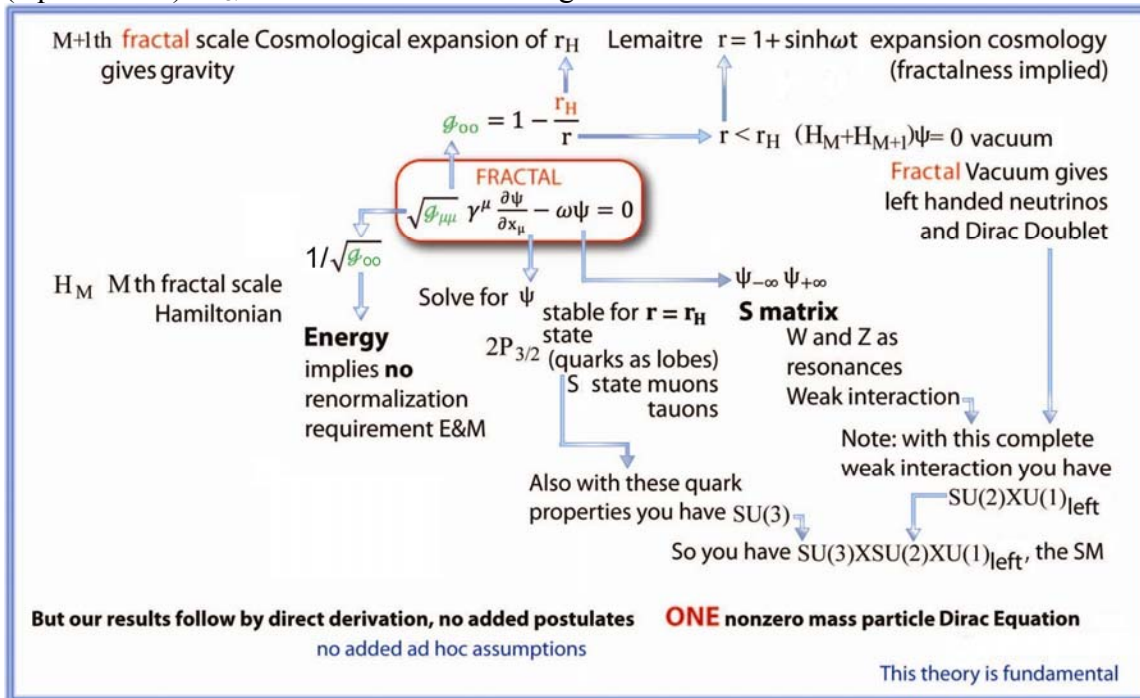


Figure 17-2 Summary of Derivations from equation 1.9

ⁱ Weinberg, Steve, *General Relativity and Cosmology*, P.257